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NEW DIFFUSION-MOBILITY-VISCOSITY-ACTIVATION ENERGY-FERMI ENERGY RELATIONS AND OPTICAL-ELECTRICAL-THERMOELECTRIC LAWS, INVESTIGATED IN N(P)-TYPE DEGENERATE “COMPENSATED” GATE(1-X) P(X)-CRYSTALLINE ALLOY, SUGGESTING A SAME VISCOSITY PHENOMENON OBSERVED IN BOTH LIQUIDS AND SOLIDS (XII)

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ABSTRACT

In degenerate $\mathbf{n^{+}(p^{+}) - p(n) - X(x) \equiv GaTe(1-x)P(x)}$ -crystalline alloy, $0 \le x \le 1$, various optical, electrical and thermoelectric laws and Stokes-Einstein-Sutherland-Reynolds-Van Cong relations, enhanced by: the optico-electrical phenomenon (O-EP) and the electro-optical phenomenon (E-OP), our static dielectric constant law given in Equations (1a, 1b), our accurate Fermi energy expression given in Eq. (11), and finally our conductivity model given in Eq. (18), are now investigated, by basing on the same physical model and the mathematical treatment method, as those used in our recent works.[1-5] One notes that, for x=0, this crystalline alloy is reduced to the n(p)-type degenerate **GaTe -crystal**. Some concluding remarks are given as follows. -By basing on our optical [electrical] conductivity models, $\sigma_{O[E]}$, given in Eq. (18), all the optical, electrical, thermoelectric coefficients have been determined, as those given in Equations (19a-19d, 20a-20d, 21-31). In particular, for the physical conditions, as those given in Eq. (15), one remarks that the optical

conductivity, σ_O, obtained from the O-EP, has a same form with that of the electrical conductivity, given from the E-OP, σ_E, as those determined in Eq. (20a), but $\sigma_O > \sigma_E$ since $m_r(x) < m_{c(v)}(x)$, $m_{c(v)}$ and m_r, being the unperturbed reduced effective electron (hole) mass in conduction (valence) bands and the relative carrier mass, respectively. Finally, the numerical results of such optical, electrical and thermoelectric coefficients, calculated by using Equations (18, 19a-19d, 20a-20d, 21-31), are reported in Tables 3-13, suggesting a same viscosity phenomenon observed in both liquids and solids.

KEYWORDS: Optical-and-electrical conductivity, Seebeck coefficient, Figure of merit, Van Cong relation, First Van-Cong coefficient, Second Van-Cong coefficient, Thomson coefficient, Peltier coefficient.

INTRODUCTION

In the $\mathbf{n^+(p^+) - X(x) \equiv GaTe(1-x)P(x)}$ -crystalline alloy, $0 \le x \le 1$, x being the concentration, the optical, electrical and thermoelectric coefficients, enhanced by : (i) the optico-electrical phenomenon (O-EP) and the electro-optical phenomenon (E-OP), (ii) our static dielectric constant law, $\varepsilon(r_{d(a)}, x)$, $r_{d(a)}$ being the donor (acceptor) d(a)-radius, given in Equations (1a, 1b), (iii) our accurate reduced Fermi energy, $\xi_{n(p)}$, given in Eq. (11), accurate with a precision of the order of 2.11×10^{-4}[9], affecting all the expressions of optical, electrical and thermoelectric coefficients, and (iv) our optical-and-electrical conductivity models, given in Eq. (18, 20a), are now investigated by basing on our physical model and Fermi-Dirac distribution function, as those given in our recent works.[1-5] It should be noted here that for x=0, the present obtained numerical results are reduced to those given in the n(p)-type degenerate **GaTe**-crystal.[1, 6-18]

Then, some important remarks can be reported as follows.

(1) As observed in Equations (3, 5, 6a, 6b), the critical impurity density $N_{CDn(CDp)}$, defined by the generalized Mott criterium in the metal-insulator transition (**MIT**), is just the density of electrons (holes), localized in the exponential conduction (valence)-band tail (**EBT**), $N^{EBT}_{CDn(CDp)}$, being obtained with a precision of the order of $\mathbf{2.91 \times 10^{-7}}$, as given in our recent work.[3] Therefore, the effective electron (hole)-density can be defined as: $N^* \equiv N - N_{CDn(CDp)} \simeq N - N^{EBT}_{CDn(CDp)}$, N being the total impurity density, as that observed in the compensated crystals.

(2) The ratio of the inverse effective screening length $k_{sn(sp)}$ to Fermi wave number $k_{Fn(kp)}$ at 0 K, $R_{sn(sp)}(N^*)$, defined in Eq. (7), is valid at any N^*.

(3) For given $[N, r_{d(a)}, x, T]$, the coefficients: $\sigma_{O[E]}(E)$, $\kappa_{O[E]}(E)$, $\varepsilon_{2\,O[E]}(E)$, and $\propto_{O[E]}(E)$, are determined in Equations (18, 19b-19d), as functions of the photon energy E, and then their numerical results are reported in Tables 3-8, being new ones.

(4) Finally, for particular physical conditions, as those given in Eq. (15), one observes that the optical conductivity σ_O has a same form with that of the electrical conductivity, σ_E, as those given in Eq. (20a), but $\sigma_O > \sigma_E$ since $m_r(x) < m_{c(v)}(x)$, $m_{c(v)}$ and m_r, being the unperturbed reduced effective electron (hole) mass in conduction (valence) bands and the relative carrier mass, respectively. Then, by basing on those $\sigma_{O[E]}$ -expressions, the thermoelectric laws, relations, and coefficients are determined in Equations (21-31), and their numerical results are reported in Tables 9 and 13, suggesting a same viscosity phenomenon observed in both liquids and solids.

In the following, various Sections are presented in order to investigate the optical, electrical and thermoelectric coefficients, given in the degenerate $n^+(p^+) - \mathbf{X}(\mathbf{x})$- crystalline alloy.

OUR STATIC DIELECTRIC CONSTANT LAW AND GENERALIZED MOTT CRITERIUM IN THE METAL-INSULATOR TRANSITION

First of all, in the degenerate $n^+(p^+) - \mathbf{X}(\mathbf{x})$- crystalline alloy, at T=0 K[1-5], we denote : the donor (acceptor) d(a)-radius by $r_{d(a)}$, the corresponding intrinsic one by: $\mathbf{r_{do(ao)}}=\mathbf{r_{Te(Ga)}}$, respectively, the effective averaged numbers of equivalent conduction (valence)-bands by: $g_{c(v)}$, the unperturbed reduced effective electron (hole) mass in conduction (valence) bands by $m_{c(v)}(x)/m_o$, m_o being the free electron mass, the relative carrier mass by: $m_r(x) \equiv \frac{m_c(x)\times m_v(x)}{m_c(x)+m_v(x)} < m_{c(v)}(x)$ for given x , the unperturbed static dielectric constant by: $\varepsilon_o(x)$, and the intrinsic band gap by: $E_{go}(x)$, as those given in **Table 1, reported in Appendix 1**.

Here, the effective carrier mass $m^*_{n(p)}(x)$ is equal to $m_{c(v)}(x)$. Therefore, we can define the effective donor (acceptor)-ionization energy in absolute values as:

$E_{do(ao)}(x) = \frac{13600\times[m_{c(v)}(x)/m_o]}{[\varepsilon_o(x)]^2}$ meV , and then, the isothermal bulk modulus, by : $B_{do(ao)}(x) \equiv \frac{E_{do(ao)}(x)}{\left(\frac{4\pi}{3}\right)\times\left(r_{do(ao)}\right)^3}$.

Our Static Dielectric Constant Law $\left[m^*_{n(p)}(x) \equiv m_{c(v)}(x)\right]$

Here, the changes in all the energy-band-structure parameters, expressed in terms of the effective dielectric constant $\varepsilon(r_{d(a)}, x)$, are developed as follows.

At $r_{d(a)} = r_{do(ao)}$, the needed boundary conditions are found to be, for the impurity-atom volumes : $V = (4\pi/3) \times \left(r_{d(a)}\right)^3$ and $V_{do(ao)} = (4\pi/3) \times \left(r_{do(ao)}\right)^3$, according to the pressures : p, $p_o = 0$, and to the deformation potential energies (or the strain energies) : α, $\alpha_o = 0$. Further, the two important equations, used to determine the α -variation, $\Delta\,\alpha \equiv \alpha - \alpha_o = \alpha$, are defined by : $\frac{dp}{dV} = -\frac{B_{do(ao)}(x)}{V}$ and $p = -\frac{d\alpha}{dV}$, giving rise to : $\frac{d}{dV}\left(\frac{d\alpha}{dV}\right) = \frac{B_{do(ao)}(x)}{V}$.

Then, by an integration, one gets :

$$\left[\Delta\alpha(r_{d(a)}, x)\right]_{n(p)} = B_{do(ao)}(x) \times (V - V_{do(ao)}) \times \ln\left(\frac{V}{V_{do(ao)}}\right) = E_{do(ao)}(x) \times \left[\left(\frac{r_{d(a)}}{r_{do(ao)}}\right)^3 - 1\right] \times \ln\left(\frac{r_{d(a)}}{r_{do(ao)}}\right)^3 \geq 0.$$

Furthermore, we also showed that, as $r_{d(a)} > r_{do(ao)}$ ($r_{d(a)} < r_{do(ao)}$), the compression (dilatation) gives rise to the increase (the decrease) in the energy gap $E_{gn(gp)}\left(r_{d(a)}, x\right)$, and the effective donor (acceptor)-ionization energy $E_{d(a)}\left(r_{d(a)}, x\right)$ in absolute values, obtained in the effective Bohr model, which is represented respectively by : $\pm\left[\Delta\alpha(r_{d(a)}, x)\right]_{n(p)}$,

$$E_{gn(gp)}(r_{d(a)}, x) - E_{go}(x) = E_{d(a)}(r_{d(a)}, x) - E_{do(ao)}(x) = E_{do(ao)}(x) \times \left[\left(\frac{\varepsilon_o(x)}{\varepsilon(r_{d(a)})}\right)^2 - 1\right] = + \left[\Delta\alpha(r_{d(a)}, x)\right]_{n(p)},$$

for $r_{d(a)} \geq r_{do(ao)}$, and for $r_{d(a)} \leq r_{do(ao)}$,

$$E_{gn(gp)}(r_{d(a)}, x) - E_{go}(x) = E_{d(a)}(r_{d(a)}, x) - E_{do(ao)}(x) = E_{do(ao)}(x) \times \left[\left(\frac{\varepsilon_o(x)}{\varepsilon(r_{d(a)})}\right)^2 - 1\right] = - \left[\Delta\alpha(r_{d(a)}, x)\right]_{n(p)}.$$

Therefore, one obtains the expressions for relative dielectric constant $\varepsilon(r_{d(a)}, x)$ and energy band gap $E_{gn(gp)}\left(r_{d(a)}, x\right)$, as :

(i)-for $r_{d(a)} \geq r_{do(ao)}$, since $\varepsilon(r_{d(a)}, x) = \frac{\varepsilon_o(x)}{\sqrt{1 + \left[\left(\frac{r_{d(a)}}{r_{do(ao)}}\right)^3 - 1\right] \times \ln\left(\frac{r_{d(a)}}{r_{do(ao)}}\right)^3}} \leq \varepsilon_o(x)$, being a **new** $\boldsymbol{\varepsilon(r_{d(a)}, x)}$**-law**,

$E_{gn(gp)}(r_{d(a)}, x) - E_{go}(x) = E_{d(a)}(r_{d(a)}, x) - E_{do(ao)}(x) = E_{do(ao)}(x) \times \left[\left(\frac{r_{d(a)}}{r_{do(ao)}}\right)^3 - 1\right] \times \ln\left(\frac{r_{d(a)}}{r_{do(ao)}}\right)^3 \geq 0,$ (1a)

according to the increase in both $E_{gn(gp)}(r_{d(a)}, x)$ and $E_{d(a)}(r_{d(a)}, x)$, with increasing $r_{d(a)}$ and for a given x, and

(ii)-for $r_{d(a)} \leq r_{do(ao)}$, since $\varepsilon(r_{d(a)}, x) = \frac{\varepsilon_o(x)}{\sqrt{1-\left[\left(\frac{r_{d(a)}}{r_{do(ao)}}\right)^3 - 1\right] \times \ln\left(\frac{r_{d(a)}}{r_{do(ao)}}\right)^3}} \geq \varepsilon_o(x)$, with a condition, given by: $\left[\left(\frac{r_{d(a)}}{r_{do(ao)}}\right)^3 - 1\right] \times \ln\left(\frac{r_{d(a)}}{r_{do(ao)}}\right)^3 < 1$, being a **new $\boldsymbol{\varepsilon}(\mathbf{r}_{\mathbf{d(a)}}, \mathbf{x})$-law**,

$E_{gn(gp)}(r_{d(a)}, x) - E_{go}(x) = E_{d(a)}(r_{d(a)}, x) - E_{do(ao)}(x) = -E_{do(ao)}(x) \times \left[\left(\frac{r_{d(a)}}{r_{do(ao)}}\right)^3 - 1\right] \times \ln\left(\frac{r_{d(a)}}{r_{do(ao)}}\right)^3 \leq 0,$ (1b)

corresponding to the decrease in both $E_{gn(gp)}(r_{d(a)}, x)$ and $E_{d(a)}(r_{d(a)}, x)$, with decreasing $r_{d(a)}$ and for a given x.

It should be noted that, in the following, all the optical, electrical and thermoelectric properties strongly depend on this **new $\boldsymbol{\varepsilon}(\mathbf{r}_{\mathbf{d(a)}}, \mathbf{x})$-law.**

Furthermore, the effective Bohr radius $a_{Bn(Bp)}(r_{d(a)}, x)$ is defined by:

$a_{Bn(Bp)}(r_{d(a)}, x) \equiv \frac{\varepsilon(r_{d(a)}, x) \times \hbar^2}{m^*_{n(p)}(x) \times m_o \times q^2} = 0.53 \times 10^{-8}\ \text{cm} \times \frac{\varepsilon(r_{d(a)}, x)}{m^*_{n(p)}(x)}$ (2)

where q=e, according to an electron charge equal to : -e.

Generalized Mott Criterium in the MIT $[m^*_{n(p)}(x) \equiv m_{c(v)}(x)]$

Now, it is interesting to remark that the critical total donor (acceptor)-density in the MIT at T=0 K, $N_{CDn(NDp)}(r_{d(a)}, x)$, was given by the Mott's criterium, with an empirical parameter, $M_{n(p)}$, as [3] :

$N_{CDn(CDp)}(r_{d(a)}, x)^{1/3} \times a_{Bn(Bp)}(r_{d(a)}, x) = M_{n(p)},\ M_{n(p)} = 0.25,$ (3)

depending thus on our **new $\boldsymbol{\varepsilon}(\mathbf{r}_{\mathbf{d(a)}}, \mathbf{x})$-law**.

This excellent one can be explained from the definition of the reduced effective Wigner-Seitz (**WS**) radius $r_{sn(sp),M}$, in the Mott's criterium, being characteristic of interactions, by :

$r_{sn(sp),M}\left(N = N_{CDn(CDp)}(r_{d(a)}, x), r_{d(a)}, x\right) \equiv \left(\frac{3}{4\pi N_{CDn(CDp)}(r_{d(a)},x)}\right)^{\frac{1}{3}} \times \frac{1}{a_{Bn(Bp)}(r_{d(a)},x)} = \mathbf{2.4813963}$, (4)

for any $(r_{d(a)}, x)$-values. Then, from Eq. (4), one also has :

$N_{CDn(CDp)}(r_{d(a)}, x)^{1/3} \times a_{Bn(Bp)}(r_{d(a)}, x) = \left(\frac{3}{4\pi}\right)^{\frac{1}{3}} \times \frac{1}{\mathbf{2.4813963}} = \mathbf{0.25} = \mathbf{M_{n(p)}}$ (5)

explaining thus the existence of the Mott's criterium.

Furthermore, by using $M_{n(p)} = 0.25$, according to the empirical Heisenberg parameter $\boldsymbol{\mathcal{H}_{n(p)}} = \mathbf{0.47137}$, as those given in our previous work[3], we have also showed that $N_{CDn(CDp)}$ is just **the density of electrons (holes) localized in the exponential conduction (valence)-band tail**, $\mathbf{N^{EBT}_{CDn(CDp)}}$, with a precision of the order of $\mathbf{2.91 \times 10^{-7}}$, respectively. [3]

So, $N_{CDn(NDp)}(r_{d(a)}, x) \cong N^{EBT}_{CDn(CDp)}(r_{d(a)}, x)$ (6a)

It shoud be noted that the values of $M_{n(p)}$ and $\mathcal{H}_{n(p)}$ could be chosen so that those of $N_{CDn(CDp)}$ and $N^{EBT}_{CDn(CDp)}$ are found to be in good agreement with their experimental results.

Therefore, the effective density of electrons (holes) given in parabolic conduction (valence) bands, N^*, can be defined, as that given in compensated materials:

$N^*(N, r_{d(a)}, x) \equiv N - N_{CDn(NDp)}(r_{d(a)}, x) \cong N - N^{EBT}_{CDn(CDp)}(r_{d(a)}, x) \geq 0.$ (6b)

In summary, as observed in our previous paper[3], for a given x and an increasing $r_{d(a)}$, $\varepsilon(r_{d(a)}, x)$ decreases, while $E_{gno(gpo)}\left(r_{d(a)}, x\right)$, $N_{CDn(NDp)}(r_{d(a)}, x)$ and $N^{EBT}_{CDn(CDp)}(r_{d(a)}, x)$ increase, affecting strongly all the optical, electrical, and thermoelectric coefficients, as those observed in following Sections.

PHYSICAL MODEL

In the degenerate $n^+(p^+) - \mathbf{X}(\mathbf{x})$ -crystalline alloy, the reduced effective Wigner-Seitz (**WS**) radius $r_{sn(sp)}$, characteristic of interactions, being given in Eq. (4), in which N is replaced by N^*, is now defined by:

$\gamma \times r_{sn(sp)}(N^*, r_{d(a)}, x) \equiv \frac{k_{Fn(Fp)}^{-1}}{a_{Bn(Bp)}} < 1$, $r_{sn(sp)}(N^*, r_{d(a)}, x) \equiv \left(\frac{3g_{c(v)}}{4\pi N^*}\right)^{1/3} \times \frac{1}{a_{Bn(Bp)}(r_{d(a)}, x)}$,

being proportional to $N^{*-1/3}$. Here, $\gamma = (4/9\pi)^{1/3}$, $k_{Fn(Fp)}(N^*) \equiv \left(\frac{3\pi^2 N^*}{g_{c(v)}}\right)^{\frac{1}{3}}$ is the Fermi wave.

Then, the ratio of the inverse effective screening length $k_{sn(sp)}$ to Fermi wave number $k_{Fn(kp)}$ is defined by:

$$R_{sn(sp)}(N^*) \equiv \frac{k_{sn(sp)}}{k_{Fn(Fp)}} = \frac{k_{Fn(Fp)}^{-1}}{k_{sn(sp)}^{-1}} = R_{snWS(spWS)} + \left[R_{snTF(spTF)} - R_{snWS(spWS)}\right] e^{-r_{sn(sp)}} < 1, \quad (7)$$

being valid at any N^*.

Here, these ratios, $R_{snTF(spTF)}$ and $R_{snWS(spWS)}$, can be determined as follows.

First, for $N \gg N_{CDn(NDp)}(r_{d(a)}, x)$, according to the **Thomas-Fermi (TF)-approximation**, the ratio $R_{snTF(spTF)}(N^*)$ is reduced to

$$R_{snTF(spTF)}(N^*) \equiv \frac{k_{snTF(spTF)}}{k_{Fn(Fp)}} = \frac{k_{Fn(Fp)}^{-1}}{k_{snTF(spTF)}^{-1}} = \sqrt{\frac{4\gamma r_{sn(sp)}}{\pi}} \ll 1, \quad (8)$$

being proportional to $N^{*-1/6}$.

Secondly, for $N \ll N_{CDn(NDp)}(r_{d(a)})$, according to the **Wigner-Seitz (WS)-approximation**, the ratio $R_{snWS(snWS)}$ is respectively reduced to

$$R_{sn(sp)WS}(N^*) \equiv \frac{k_{sn(sp)WS}}{k_{Fn}} = 0.5 \times \left(\frac{3}{2\pi} - \gamma \frac{d\left[r_{sn(sp)}^2 \times E_{CE}(N^*)\right]}{dr_{sn(sp)}}\right), \quad (9a)$$

where $E_{CE}(N^*)$ is the majority-carrier correlation energy (CE), being determined by:

$$E_{CE}(N^*) = \frac{-0.87553}{0.0908 + r_{sn(sp)}} + \frac{\frac{0.87553}{0.0908 + r_{sn(sp)}} + \left(\frac{2[1-\ln(2)]}{\pi^2}\right) \times \ln(r_{sn(sp)}) - 0.093288}{1 + 0.03847728 \times r_{sn(sp)}^{1.67378876}}.$$

Furthermore, in the highly degenerate case, the physical conditions are found to be given by:

$$\frac{k_{Fn(Fp)}^{-1}}{a_{Bn(Bp)}} < \frac{\sigma_{n(p)}}{E_{Fno(Fpo)}} \equiv \frac{1}{A_{n(p)}} < \frac{k_{Fn(Fp)}^{-1}}{k_{sn(sp)}^{-1}} \equiv R_{sn(sp)} < 1, \quad \sigma_{n(p)}(N^*) \equiv \frac{\sqrt{2\pi \times \left(\frac{N^*}{g_{c(v)}}\right)}}{\varepsilon(r_{d(a)})} \times q^2 k_{sn(sp)}^{-1/2} \quad (9b)$$

which gives: $A_{n(p)}(N^*) = \frac{E_{Fno(Fpo)}(N^*)}{\sigma_{n(p)}(N^*)}$, $E_{Fno(Fpo)}(N^*, r_{d(a)}, x) \equiv \frac{\hbar^2 \times k_{Fn(Fp)}^2(N^*)}{2 \times m_{n(p)}^*(x) \times m_o}$.

BAND GAP NARROWING (BGN)

First, the BGN is found to be given by:

$$\Delta E_{gn(gp)}\left(N^*, r_{d(a)}, x\right) \simeq$$

$$a_1 + \frac{\varepsilon_o(x)}{\varepsilon(r_{d(a)},x)} \times N_r^{\frac{1}{3}} + a_2 \times \frac{\varepsilon_o(x)}{\varepsilon(r_{d(a)},x)} \times N_r^{\frac{1}{3}} \times \left(2.503 \times \left[-E_{CE}\left(r_{sn(sp)}\right)\right] \times r_{sn(sp)}\right) + a_3 \times \left[\frac{\varepsilon_o(x)}{\varepsilon(r_{d(a)},x)}\right]^{\frac{5}{4}} \times \sqrt{\frac{m_{v(c)}}{m^*_{n(p)}(x)}} \times N_r^{\frac{1}{4}} + 2a_4 \times \left[\frac{\varepsilon_o(x)}{\varepsilon(r_{d(a)},x)}\right]^{\frac{1}{2}} \times N_r^{\frac{1}{2}} + 2a_5 \times \left[\frac{\varepsilon_o(x)}{\varepsilon(r_{d(a)},x)}\right]^{\frac{3}{2}} \times N_r^{\frac{1}{6}},\ N_r = \frac{N^*}{9.999\times 10^{17}\,cm^{-3}},\ (10a)$$

Here, for $\Delta E_{gn;N}(N^*, r_d, x)$, one has: $a_1 = 3.8 \times 10^{-3}(eV)$, $a_2 = 6.5 \times 10^{-4}(eV)$, $a_3 = 2.8 \times 10^{-3}(eV)$, $a_4 = 5.597 \times 10^{-3}(eV)$, and $a_5 = 8.1 \times 10^{-4}(eV)$, and for $\Delta E_{gp;N}(N^*, r_a, x)$, one has: $a_1 = 3.15 \times 10^{-3}(eV)$, $a_2 = 5.41 \times 10^{-4}(eV)$, $a_3 = 2.32 \times 10^{-3}(eV)$, $a_4 = 4.12 \times 10^{-3}(eV)$, and $a_5 = 9.8 \times 10^{-5}(eV)$.

Therefore, at T=0 K and $N^* = 0$, and for any x and $r_{d(a)}$, one gets: $\Delta E_{gn(gp)} = 0$, according to the metal-insulator transition (MIT).

Secondly, one has:

$$\Delta E_{gn(gp)}(T, x) = 10^{-4} T^2 \times \left[\frac{7.205\times x}{T+94} + \frac{5.405\times(1-x)}{T+204}\right] (eV). \quad (10b)$$

FERMI ENERGY AND FERMI-DIRAC DISTRIBUTION FUNCTION

Fermi Energy

Here, for a presentation simplicity, we change all the sign of various parameters, given in the degenerate $p^+ - \mathbf{X(x)}$-crystalline alloy, in order to obtain the same one, as given in the degenerate $n^+ - \mathbf{X(x)}$ - crystalline alloy, according to the reduced Fermi energy $E_{Fn(Fp)}$, $\xi_{n(p)}(N^*, r_{d(a)}, x, T) \equiv \frac{E_{Fn(Fp)}(N^*, r_{d(a)}, x, T)}{k_B T} > 0 (< 0)$, obtained respectively in the degenerate (non-degenerate) case.

For any $(N^*, r_{d(a)}, x, T)$, the reduced Fermi energy $\xi_{n(p)}(N^*, r_{d(a)}, x, T)$ or the Fermi energy $E_{Fn(Fp)}(N^*, r_{d(a)}, x, T)$, obtained in our previous paper[9], obtained with a precision of the order of $\mathbf{2.11 \times 10^{-4}}$, is found to be given by:

$$\xi_{n(p)}(u) \equiv \frac{E_{Fn(Fp)}(u)}{k_B T} = \frac{G(u) + Au^B F(u)}{1 + Au^B} \equiv \frac{V(u)}{W(u)},\ A = 0.0005372 \text{ and } B = 4.82842262,\ (11)$$

where u is the reduced electron density, $u(N^*, r_{d(a)}, x, T) \equiv \frac{N^*}{N_{c(v)}(T,x)}$, $N_{c(v)}(T,x) = 2g_{c(v)} \times \left(\frac{m^*_{n(p)}(x) \times m_o \times k_B T}{2\pi\hbar^2}\right)^{\frac{3}{2}}$ (cm^{-3}), $F(u) = au^{\frac{2}{3}}\left(1 + bu^{-\frac{4}{3}} + cu^{-\frac{8}{3}}\right)^{-\frac{2}{3}}$, $a = \left[3\sqrt{\pi}/4\right]^{2/3}$, $b = \frac{1}{8}\left(\frac{\pi}{a}\right)^2$, $c = \frac{62.3739855}{1920}\left(\frac{\pi}{a}\right)^4$, and $G(u) \simeq Ln(u) + 2^{-\frac{3}{2}} \times u \times e^{-du}$; $d = 2^{3/2}\left[\frac{1}{\sqrt{27}} - \frac{3}{16}\right] > 0$.

So, in the non-degenerate case ($u \ll 1$), one has: $E_{Fn(Fp)}(u) = k_B T \times G(u) \simeq k_B T \times Ln(u)$ as **$u \longrightarrow 0$, the limiting non-degenerate condition**, and in the very degenerate case ($u \gg 1$), one gets: $E_{Fn(Fp)}(u \gg 1) = k_B T \times F(u) = k_B T \times au^{\frac{2}{3}}\left(1 + bu^{-\frac{4}{3}} + cu^{-\frac{8}{3}}\right)^{-\frac{2}{3}} \simeq \frac{\hbar^2 \times k^2_{Fn(Fp)}(N^*)}{2 \times m^*_{n(p)}(x) \times m_o}$ as **$u \longrightarrow \infty$, the limiting degenerate condition**. In other words, $\xi_{n(p)} \equiv \frac{E_{Fn(Fp)}}{k_B T}$ is accurate, and it also verifies the correct limiting conditions.

In particular, as $T \to 0\,K$, since $u^{-1} \to 0$, Eq. (11) is reduced to: $E_{Fno(Fpo)}(N^*) \equiv \frac{\hbar^2 \times k^2_{Fn(Fp)}(N^*)}{2 \times m^*_{n(p)}(x) \times m_o}$, proportional to $(N^*)^{2/3}$, noting that, for a given N^*, $E_{Fno(Fpo)}\left(m^*_{n(p)}(x) = m_r(x)\right) > E_{Fno(Fpo)}\left(m^*_{n(p)}(x) = m_{c(v)}(x)\right)$ since $m_r(x) < m_{c(v)}(x)$ for given x. Further, at T=0 K and $N^* = 0$, being the physical conditions, given for the metal-insulator transition (MIT).

In the following, it should be noted that all the optical and electrical-and-thermoelectric properties strongly depend on such the accurate expression of $\xi_{n(p)}(N^*, r_{d(a)}, x, T)$.[9]

Fermi-Dirac Distribution Function (FDDF)

The Fermi-Dirac distribution function (FDDF) is given by: $f(E) \equiv (1 + e^{\gamma})^{-1}$, $\gamma \equiv (E - E_{Fn(Fp)})/(k_B T)$.

So, the average of E^p, calculated using the FDDF-method, as developed in our previous works [1, 6] is found to be given by:

$\langle E^p \rangle_{FDDF} \equiv G_p(E_{Fn(Fp)}) \times E^p_{Fn(Fp)} \equiv \int_{-\infty}^{\infty} E^p \times \left(-\frac{\partial f}{\partial E}\right) dE$, $\quad -\frac{\partial f}{\partial E} = \frac{1}{k_B T} \times \frac{e^{\gamma}}{(1+e^{\gamma})^2}$.

Further, one notes that, at 0 K, $-\frac{\partial f}{\partial E} = \delta\left(E - E_{Fno(Fpo)}\right)$, $\delta\left(E - E_{Fno(Fpo)}\right)$ being the Dirac delta (δ)-function. Therefore, $G_p\left(E_{Fno(Fpo)}\right) = 1$.

Then, at low T, by a variable change $\gamma \equiv (E - E_{Fn(Fp)})/(k_B T)$, one has:

$G_p(E_{Fn(Fp)}) \equiv 1 + E_{Fn(Fp)}^{-p} \times \int_{-\infty}^{\infty} \frac{e^\gamma}{(1+e^\gamma)^2} \times (k_B T\gamma + E_{Fn(Fp)})^p d\gamma = 1 + \sum_{\mu=1,2,...}^{p} C_p^\beta \times (k_B T)^\beta \times E_{Fn(Fp)}^{-\beta} \times I_\beta$, where $C_p^\beta \equiv p(p-1)...(p-\beta+1)/\beta!$ and the integral I_β is given by:

$I_\beta = \int_{-\infty}^{\infty} \frac{\gamma^\beta \times e^\gamma}{(1+e^\gamma)^2} d\gamma = \int_{-\infty}^{\infty} \frac{\gamma^\beta}{(e^{\gamma/2}+e^{-\gamma/2})^2} d\gamma$, vanishing for old values of β. Then, for even values of $\beta = 2n$, with n=1, 2, …, one obtains: .

$I_{2n} = 2\int_0^\infty \frac{\gamma^{2n} \times e^\gamma}{(1+e^\gamma)^2} d\gamma$.

Now, using an identity $(1+e^\gamma)^{-2} \equiv \sum_{s=1}^{\infty}(-1)^{s+1} s \times e^{\gamma(s-1)}$, a variable change: $s\gamma = -t$, the Gamma function: $\int_0^\infty t^{2n} e^{-t} dt \equiv \Gamma(2n+1) = (2n)!$, and also the definition of the Riemann's zeta function: $\zeta(2n) \equiv 2^{2n-1}\pi^{2n}|B_{2n}|/(2n)!$, B_{2n} being the Bernoulli numbers, one finally gets: $I_{2n} = (2^{2n}-2) \times \pi^{2n} \times |B_{2n}|$. So, from above Eq. of $\langle E^p \rangle_{FDDF}$, we get in the degenerate case the following ratio:

$G_p(E_{Fn(Fp)}) \equiv \frac{\langle E^p \rangle_{FDDF}}{E_{Fn(Fp)}^p} = 1 + \sum_{n=1}^{p} \frac{p(p-1)...(p-2n+1)}{(2n)!} \times (2^{2n}-2) \times |B_{2n}| \times y^{2n} \equiv G_{p\geq 1}(y)$,(12)

where $y \equiv \frac{\pi}{\xi_{n(p)}(N^*, r_{d(a)}, x, T)} = \frac{\pi k_B T}{E_{Fn(Fp)}(N^*, r_{d(a)}, x, T)}$, noting that $G_{p=1}(y \equiv \frac{\pi k_B T}{E_{Fn(Fp)}} = \frac{\pi}{\xi_{n(p)}}) = 1$, and as T→ 0 K, $G_{p>1}(y \to 0) \to 1$.

Then, some usual results of $G_{p\geq 1}(y)$ are given in the **Table 2, reported in Appendix 1**, being needed to determine all the following optical, electrical and thermoelectric properties.

OPTICAL-AND-ELECTRICAL PROPERTIES

Optico-Electrical Phenomenon (O-EP) and Electro-Optical Phenomenon (E-OP)

In the degenerate $n^+(p^+) - \mathbf{X}(\mathbf{x})$-crystalline alloy, one has:

(i) in the **E-OP**, the reduced band gap is defined by:

$E_{gn2(gp2)} \equiv E_{gn(gp)} - \Delta E_{gn(gp)}(N^*, r_{d(a)}, x) - \Delta E_{gn(gp)}(T, x)$ (13)

where the intrinsic band gap $E_{gn(gp)}$ is defined in Equations (1a, 1b), $\Delta E_{gn(gp)}(N^*, r_{d(a)}, x)$ and $\Delta E_{gn(gp)}(T, x)$ are respectively determined in Equations (10a, 10b), and

(ii) in the (**O-EP**), the photon energy is defined by: $E \equiv \hbar\omega$, and the optical band gap, as: $E_{gn1(gp1)} \equiv E_{gn2(gp2)} + E_{Fn(Fp)}$.

Therefore, for $E \geq E_{gn1(gp1)}(E_{gn2(gp2)})$, the effective photon energy E^* is found to be given by:

$E^* \equiv E - E_{gn1(gp1)}(E_{gn2(gp2)}) \geq 0.$ (14)

From above Equations, one notes that: $E^* \equiv [E - E_{gn1(gp1)}] = E_{Fn(Fp)}$, given in the O-EP, if $E = [E_{gn1(gp1)} + E_{Fn(Fp)}] \equiv E_{gn(gp)O}$ and $m^*_{n(p)}(x) = m_r(x)$, and $E^* \equiv E - E_{gn2(gp2)} = E_{Fn(Fp)}$, given in the E-OP, if $E = [E_{gn2(gp2)} + E_{Fn(Fp)}] \equiv E_{gn(gp)E}$ and $m^*_{n(p)}(x) = m_{c(v)}(x)$, noting that $E_{Fn(Fp)}(m_r(x)) > E_{Fn(Fp)}(m_{c(v)}(x))$, since $m_r(x) < m_{c(v)}(x)$, for a given x. (15)

Eq. (15) thus shows that, in both O-EP and E-OP, the Fermi energy-level penetrations into conduction (valence)-bands, observed in the $n^+(p^+) -$ type degenerate $n^+(p^+) - \mathbf{X}(\mathbf{x})$-crystalline alloy, $E_{Fn(Fp)}$, are well defined.

Optical Coefficients

The optical properties for any medium, defined in the O-EP and E-OP, respectively, according to: $[m^*_{n(p)} \equiv m_r(x)[m_{c(v)}(x)]]$, can be described by the complex refraction: $\mathbb{N}_{O[E]} \equiv n_{O[E]} - i\kappa_{O[E]}$, $n_{O[E]}$ and $\kappa_{O[E]}$ being the refraction index and the extinction coefficient, the complex dielectric function: $\mathcal{E}_{O[E]} = \varepsilon_{1\,O[E]} - i\varepsilon_{2\,O[E]}$, where $i^2 = -1$, and $\mathcal{E}_{O[E]} = \mathbb{N}_{O[E]}{}^2$. Further, if denoting the normal-incidence reflectance and the optical absorption by $R_{O[E]}$ and $\propto_{O[E]}$, and the effective joint parabolic conduction (parabolic valence)-band density of states by:

$JDOS_{n(p)\,O[E]}(E, N^*, r_{d(a)}, x, T) \equiv$

$\frac{1}{2\pi^2} \times \left(\frac{2m^*_{n(p)}(x)}{\hbar^2}\right)^{3/2} \times \sqrt{E_{Fno(Fpo)}(N^*)} \times \left[\frac{E - E_{gn1(gp1)}(E_{gn2(gp2)})}{E - [E_{gn1(gp1)}(E_{gn2(gp2)}) + E_{Fn(Fp)} - E_{Fno(Fpo)}]}\right]^2$, and

$F_{O[E]}(E) \equiv \frac{\hbar q^2 \times |v(E)|^2}{n_{O[E]}(E) \times cE \times \varepsilon(r_{d(a)}, x) \times \varepsilon_{free\ space}}$, one gets [2]:

$\propto_{O[E]}(E) = JDOS_{n(p)\,O[E]}(E) \times F_{O[E]}(E) = \frac{E \times \varepsilon_{2\,O[E]}(E)}{\hbar c\, n_{O[E]}(E)} = \frac{2E \times \kappa_{O[E]}(E)}{\hbar c} =$

$\frac{4\pi\sigma_{O[E]}(E)}{cn_{O[E]}(E) \times \varepsilon(r_{d(a)}, x) \times \varepsilon_{free\ space}}$,

$\varepsilon_{1\ O[E]}(E) \equiv {n_{O[E]}}^2 - {\kappa_{O[E]}}^2$, $\varepsilon_{2\ O[E]}(E) \equiv 2\kappa_{O[E]} n_{O[E]}$, and $R_{O[E]}(E) \equiv \frac{[n_{O[E]}-1]^2 + {\kappa_{O[E]}}^2}{[n_{O[E]}+1]^2 + {\kappa_{O[E]}}^2}$. (16a)

One notes here that, at the MIT-conditions : $N^* = 0$, both $E_{gn1(gp1)}(E_{gn2(gp2)}) = E_{gn(gp)}$, according to

$\left[\frac{E-E_{gn1(gp1)}(E_{gn2(gp2)})}{E-[E_{gn1(gp1)}(E_{gn2(gp2)})+E_{Fn(Fp)}-E_{Fno(Fpo)}]}\right]^2 = \frac{0}{0}$ for $E=E_{gn(gp)}$,

$\left[\frac{E-E_{gn1(gp1)}(E_{gn2(gp2)})}{E-[E_{gn1(gp1)}(E_{gn2(gp2)})+E_{Fn(Fp)}-E_{Fno(Fpo)}]}\right]^2 = 1$ for $E \gtrsim E_{gn(gp)}$, so that, in such the MIT,

$JDOS_{n(p)\ O[E]}(E, N^*, r_{d(a)}, x, T) \equiv \frac{1}{2\pi^2} \times \left(\frac{2m^*_{n(p)}(x)}{\hbar^2}\right)^{\frac{3}{2}} \times \sqrt{E_{Fno(Fpo)}(N^* = 0)} = 0$, for $E \gtrsim E_{gn(gp)}$, which is largely verified since $N_{CDn(NDp)}(r_{d(a)}, x) \cong N^{EBT}_{CDn(CDp)}(r_{d(a)}, x)$ or $E_{gn2(gp2)}(N_{CDn(NDp)}, T = 0K) \cong E_{gn2(gp2)}(N^{EBT}_{CDn(CDp)}, T = 0K) \cong E_{gn(gp)}$, as those given in Equations (6a, 6b). In other words, the critical photon energy can be defined by: $E \cong E_{gn(gp)}$.

Then, Eq. (6a) states that $N_{CDn(CDp)}$, given in parabolic conduction (parabolic valence)-band density of states, is just the density of electrons (holes) localized in the exponential conduction (valence)-band tail, $N^{EBT}_{CDn(CDp)}$, with a precision of the order of $\mathbf{2.91 \times 10^{-7}}$, respectively.[3] Therefore, for $E \cong E_{gn(gp)}$, the exponential conduction (valence)-band tail states can be approximated with a same precision to:

$$JDOS^{EBT}_{n(p)O[E]}(E, N^*, r_{d(a)}, x, T) \equiv \frac{1}{2\pi^2} \times \left(\frac{2m^*_{n(p)}(x)}{\hbar^2}\right)^{\frac{3}{2}} \times \sqrt{E_{Fno(Fpo)}(N^* = N_{CDn(NDp)})}. \quad (16b)$$

Here, $\varepsilon_{free\ space} = 8.854187817 \times 10^{-12} (\frac{C^2}{N \times m^2})$ is the permittivity of the free space, -q is the charge of the electron, $|v_{O[E]}(E)|$ is the matrix elements of the velocity operator between valence (conduction)-and-conduction (valence) bands, and our approximate expression for the refraction index $n_{O[E]}$ is found to be defined by:

$$n_{O[E]}(E, N^*, r_{d(a)}, x, T) = n_\infty(r_{d(a)}, x) + \sum_{i=1}^{4} \frac{X_i(E_{gn1(gp1)}) \times E + Y_i(E_{gn1(gp1)})}{E^2 - B_i E + C_i} \quad (17)$$

going to a constant as $E \to \infty$, since $n(E \to \infty, r_{d(a)}, x) \to n_\infty(r_{d(a)}, x) = \sqrt{\varepsilon(r_{d(a)}, x)} \times \frac{\omega_T}{\omega_L}$, given in the well-known Lyddane-Sachs-Teller relation, in which $\omega_T \simeq 5.1 \times 10^{13}\ s^{-1}$ and

$\omega_L \simeq 8.9755 \times 10^{13}\ s^{-1}$ are the transverse (longitudinal) optical phonon frequencies, giving rise to: $n_\infty(r_{d(a)}, x) \simeq \sqrt{\varepsilon(r_{d(a)}, x)} \times 0.568$.

Here, the other parameters are determined by: $X_i(E_{gn1(gp1)}) = \frac{A_i}{Q_i} \times \left[-\frac{B_i^2}{2} + E_{gn1(gp1)}B_i - E_{gn1(gp1)}^2 + C_i\right]$, $Y_i(E_{gn1(gp1)}) = \frac{A_i}{Q_i} \times \left[\frac{B_i \times (E_{gn1(gp1)}^2 + C_i)}{2} - 2E_{gn1(gp1)}C_i\right]$, $Q_i = \frac{\sqrt{4C_i - B_i^2}}{2}$, where, for i=(1, 2, 3, and 4),

$A_i = 4.7314 \times 10^{-4},\ 0.2313655, 0.1117995, 0.011632$, $B_i = 5.871, 6.154, 9.679,\ 13.232$, and $C_i = 8.619, 9.784, 23.803, 44.119$.

Now, the optical [electrical] conductivity $\sigma_{O[E]}$ can be defined and expressed in terms of the kinetic energy of the electron (hole), $E_k \equiv \frac{\hbar^2 \times k^2}{2 \times m_{n(p)}^*(x) \times m_o}$, k being the wave number, as:

$\sigma_{O[E]}(k) \equiv \frac{q^2 \times k}{\pi \times \hbar} \times \frac{k}{k_{sn(sp)}} \times [k \times a_{Bn(Bp)}] \times \left(\frac{E_k}{\eta_{n(p)}}\right)^{\frac{1}{2}} \left(\frac{1}{\Omega \times cm}\right)$, which is thus proportional to ${E_k}^2$, where $\frac{q^2}{\pi \times \hbar} = 7.7480735 \times 10^{-5}\ ohm^{-1}$.

Then, we obtain: $\langle E^2 \rangle_{FDDF} \equiv G_2(y = \frac{\pi k_B T}{E_{Fn(Fp)}}) \times E_{Fn(Fp)}^2$, and $G_2(y) = \left(1 + \frac{y^2}{3}\right) \equiv G_2(N, r_{d(a)}, x, T)$, with $y \equiv \frac{\pi}{\xi_{n(p)}}$, $\xi_{n(p)} = \xi_{n(p)}(N^*, r_{d(a)}, x, T)$ for a presentation simplicity.

Therefore, from above equations (16, 17), if denoting the function $H(N, r_{d(a)}, x, T)$ by:

$H(N^*, r_{d(a)}, x, T) =$

$\left[\frac{k_{Fn(Fp)}(N^*)}{R_{sn(sp)}(N^*)} \times [k_{Fn(Fp)}(N^*) \times a_{Bn(Bp)}(r_{d(a)}, x)] \times \sqrt{A_{n(p)}(N^*) = \frac{E_{Fno(Fpo)}(N^*)}{\sigma_{n(p)}(N^*)}}\right] \times G_2(N^*, r_{d(a)}, x, T)$, being proportional to $E_{Fno(Fpo)}^2$, where $\sigma_{n(p)}(N^*)$ is determined in Eq. Eq. (9b). Here, $R_{sn(sp)}(N^*) \equiv \frac{k_{sn(sp)}}{k_{Fn(Fp)}}$ is defined in Eq. (7). Then, our optical [electrical] conductivity models, defined in the O-EP and E-OP, respectively, for a simply representation, can thus be assumed to be as:

$\sigma_O(E, N^*, r_{d(a)}, x, T) =$

$\frac{q^2}{\pi \times \hbar} \times H(N^*, r_{d(a)}, x, T) \times \left[\frac{E - E_{gn1(gp1)}}{E - [E_{gn1(gp1)} + E_{Fn(Fp)} - E_{Fno(Fpo)}]}\right]^2 \left(\frac{1}{\Omega \times cm}\right)$, and

$$\sigma_E\left(E, N, r_{d(a)}, x, T\right) =$$

$$\frac{q^2}{\pi\times\hbar} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left[\frac{E-E_{gn2(gp2)}}{E-[(E_{gn2(gp2)}+E_{Fn(Fp)}-E_{Fno(Fpo)}]}\right]^2 \left(\frac{1}{\Omega\times cm}\right). \quad (18)$$

It should be noted here that:

(i) $\sigma_{O[E]}\left(E = E_{gn1(gp1)}[E_{gn2(gp2)}]\right) = 0$, and $\sigma_{O[E]}(E \to \infty) \to \text{Constant}$ for given $\left(N, r_{d(a)}, x, T\right)$ −physical conditions, and

(ii) as $T \to 0$ K and $N^* = 0$ [or $E_{Fno(Fpo)}(N^*) = 0$], according to: $H\left(N^*, r_{d(a)}, x, T\right) = 0$, and for a given E, $\left[E - E_{gn1(gp1)}\right] = \left[E - E_{gn(gp)}\right]$=Constant, then from Equations (16-18), $n_{O[E]}$ (E)= Constant, $\sigma_{O[E]}(E) = 0$, $\kappa_{O-EP[E-OP]}(E) = 0$, $\varepsilon_{1\,O[E]}(E) = (n_\infty)^2 = \text{Constant}$, $\varepsilon_2(E) = 0$, and $\propto_{O[E]}(E) = 0$.

This result (18) should be new, in comparison with that, obtained from an improved Forouhi-Bloomer parameterization, as given in our previous work.[2]

Using Equations (16-18), one obtains all the analytically results as:

$$\frac{|v(E)|^2}{E} =$$

$$\frac{8\pi^2\hbar}{(2m_r)^{\frac{3}{2}}\times\sqrt{\eta_{n(p)}}} \times \left[\frac{k_{Fn(Fp)}(N^*)}{R_{sn(sp)}(N^*)} \times \left[k_{Fn(Fp)}(N^*) \times a_{Bn(Bp)}\left(r_{d(a)}, x\right)\right]\right] \times G_2\left(N^*, r_{d(a)}, x, T\right), \quad (19a)$$

$$\kappa_O(E) = \frac{2q^2}{n(E)\times\varepsilon(r_{d(a)},x)\times\varepsilon_{\text{free space}}\times E} \times H\left(N^*, r_{d(a)}, x, T\right) \times$$

$$\left[\frac{E-E_{gn1(gp1)}}{E-[E_{gn1(gp1)}+E_{Fn(Fp)}-E_{Fno(Fpo)}]}\right]^2 \quad \text{and}$$

$$\kappa_E(E) = \frac{2q^2}{n(E)\times\varepsilon(r_{d(a)},x)\times\varepsilon_{\text{free space}}\times E} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left[\frac{E-E_{gn2(gp2)}}{E-[(E_{gn2(gp2)}+E_{Fn(Fp)}-E_{Fno(Fpo)}]}\right]^2 \quad ,$$

19b)

which gives: $\kappa_{O[E]}\left(E = E_{gn1(gp1)}[E_{gn2(gp2)}]\right) = 0$, and $\kappa_{O[E]}(E \to \infty) \to 0$, as those given in Ref. [2],

$$\varepsilon_{2\,O}(E) = \frac{4q^2}{\varepsilon(r_{d(a)},x)\times\varepsilon_{\text{free space}}\times E} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left[\frac{E-E_{gn1(gp1)}}{E-[E_{gn1(gp1)}+E_{Fn(Fp)}-E_{Fno(Fpo)}]}\right]^2 \quad \text{and}$$

$$\varepsilon_{2\,E}(E) = \frac{4q^2}{\varepsilon(r_{d(a)},x)\times\varepsilon_{\text{free space}}\times E} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left[\frac{E-E_{gn2(gp2)}}{E-[(E_{gn2(gp2)}+E_{Fn(Fp)}-E_{Fno(Fpo)}]}\right]^2, \quad (19c)$$

which gives: $\varepsilon_{2O-EP[2E-OP]}\left(E = E_{gn1(gp1)}[E_{gn2(gp2)}]\right) = 0$, and $\varepsilon_{2O-EP[2E-OP]}(E \to \infty) \to 0$, as those given in Ref. [2],

$$\propto_O (E) = \frac{4q^2}{\hbar c n(E) \times \varepsilon(r_{d(a)}, x) \times \varepsilon_{\text{free space}}} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left[\frac{E - E_{gn1(gp1)}}{E - [E_{gn1(gp1)} + E_{Fn(Fp)} - E_{Fno(Fpo)}]}\right]^2 \left(\frac{1}{cm}\right) \text{ and}$$

$$\propto_E (E) = \frac{4q^2}{\hbar c n(E) \times \varepsilon(r_{d(a)}, x) \times \varepsilon_{\text{free space}}} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left[\frac{E - E_{gn2(gp2)}}{E - [(E_{gn2(gp2)} + E_{Fn(Fp)} - E_{Fno(Fpo)}]}\right]^2 \left(\frac{1}{cm}\right), \quad (19d)$$

which gives: $\propto_{O[E]} \left(E = E_{gn1(gp1)}[E_{gn2(gp2)}]\right) = 0$, and $\propto_{O[E]} (E \to \infty) \to \text{Constant}$. Furthermore, from Equations (16, 17, 19b), we can also determine $\varepsilon_{1\ O[E]}(E)$ and $R_{O[E]}(E)$.

Now, from Equations (18, 19b, 19c, 19d), using Eq. (15) as $E \equiv E_{gn(gp)O[E]}$, one obtains respectively, as:

$$\sigma_O\left(N^*, r_{d(a)}, x, T\right) = \frac{q^2}{\pi \times \hbar} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left(\frac{E_{Fn(Fp)}}{E_{Fno(Fpo)}}\right)^2 \left(\frac{1}{\Omega \times cm}\right),$$

having the same form with that of $\sigma_E\left(N, r_{d(a)}, x, T\right)$ [1], as:

$$\sigma_E\left(N^*, r_{d(a)}, x, T\right) = \frac{q^2}{\pi \times \hbar} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left(\frac{E_{Fn(Fp)}}{E_{Fno(Fpo)}}\right)^2 \left(\frac{1}{\Omega \times cm}\right), \quad (20a)$$

noting here that for given $\left(N^*, r_{d(a)}, x, T\right)$ -physical conditions we obtain: $\sigma_O > \sigma_E$ since $m_r(x) < m_{c(v)}(x)$,

$$\kappa_O\left(N^*, r_{d(a)}, x, T\right) = \frac{2q^2}{n(E) \times \varepsilon(r_{d(a)}, x) \times \varepsilon_{\text{free space}} \times (E_{gn1(gp1)} + E_{Fn(Fp)})} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left(\frac{E_{Fn(Fp)}}{E_{Fno(Fpo)}}\right)^2 \text{ and}$$

$$\kappa_E\left(N^*, r_{d(a)}, x, T\right) = \frac{2q^2}{n(E) \times \varepsilon(r_{d(a)}, x) \times \varepsilon_{\text{free space}} \times (E_{gn2(gp2)} + E_{Fn(Fp)})} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left(\frac{E_{Fn(Fp)}}{E_{Fno(Fpo)}}\right)^2, \quad (20b)$$

$$\varepsilon_{2\ O}\left(N^*, r_{d(a)}, x, T\right) = \frac{4q^2}{\varepsilon(r_{d(a)}, x) \times \varepsilon_{\text{free space}} \times (E_{gn1(gp1)} + E_{Fn(Fp)})} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left(\frac{E_{Fn(Fp)}}{E_{Fno(Fpo)}}\right)^2$$

and

$$\varepsilon_{2\,E}\left(N^*, r_{d(a)}, x, T\right) = \frac{4q^2}{\varepsilon(r_{d(a)},x) \times \varepsilon_{\text{free space}} \times (E_{gn2(gp2)} + E_{Fn(Fp)})} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left(\frac{E_{Fn(Fp)}}{E_{Fno(Fpo)}}\right)^2,$$

(20c) $$\propto_O \left(N^*, r_{d(a)}, x, T\right) = \frac{4q^2}{\hbar cn(E) \times \varepsilon(r_{d(a)},x) \times \varepsilon_{\text{free space}}} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left(\frac{E_{Fn(Fp)}}{E_{Fno(Fpo)}}\right)^2 \left(\frac{1}{cm}\right)$$ and

$$\propto_E \left(N^*, r_{d(a)}, x, T\right) = \frac{4q^2}{\hbar cn(E) \times \varepsilon(r_{d(a)},x) \times \varepsilon_{\text{free space}}} \times H\left(N^*, r_{d(a)}, x, T\right) \times \left(\frac{E_{Fn(Fp)}}{E_{Fno(Fpo)}}\right)^2 \left(\frac{1}{cm}\right).$$

(20d)

Further, from Equations (16, 17, 20b), we can also determine $\varepsilon_{1\,O[E]}(E)$ and $R_{O[E]}(E)$.

Then, the numerical results of various O[E]-coefficients, $X_{O[E]}\left(E, N^*, r_{d(a)}, x, T\right)$, as functions of E, obtained from Equations (18, 19b-19d, 20a-20d) for given $\left(N^*, r_{d(a)}, x, T\right)$-physical conditions and $E \geq (\text{or} \leq)\ E_{gn1(gp1)}(E_{gn2(gp2)})$, giving raise to the metal-insulator transition (MIT) and the non-MIT (N-MIT), are reported and discussed as follows.

First of all, one notes that from Equations (3, 6a, 6b) the MIT occurs as T=0 K and $N^*(N, r_{d(a)}, x) \equiv N - N_{CDn(NDp)}(r_{d(a)}, x) \cong N - N^{EBT}_{CDn(CDp)}\left(r_{d(a)}, x\right) = 0$, according, for $E \geq E_{gn(gp)}$, to: $E_{Fno(Fpo)}(N^* = 0) \equiv \frac{\hbar^2 \times k^2_{Fn(Fp)}(N^*)}{2 \times m^*_{n(p)}(x) \times m_o} = 0$, and $\kappa^{MIT}_{O[E]}(E, N^* = 0) = 0$, $\varepsilon^{MIT}_{2\,O[E]}(E, N^* = 0) = 0$, $\sigma^{MIT}_{O[E]}(E, N^* = 0) = 0$ and $\propto^{MIT}_{O[E]}(E, N^* = 0)=0$, since, for example, $\sigma_{E[O]}(E, N^* = 0)$ is proportional to $E^2_{Fno(Fpo)}$, or to $(N^* = 0)^{\frac{4}{3}} = 0$. But, for such the same physical conditions: T=0 K, $N^* = 0$ and $E \geq E_{gn(gp)}$, we obtain other numerical results such as: $n^{N-MIT}_{O[E]}(N^* = 0, E) \neq 0$, $\varepsilon^{N-MIT}_{1\,O[E]}(N^* = 0, E) \neq 0$ and $R^{N-MIT}_{O[E]}(N^* = 0, E) \neq 0$, for $E \geq E_{gn(gp)}$, according to the non-MIT (N-MIT), as showed in **Tables 3, 4n and 4p, reported in Appendix 1**. These Tables also state that, at T=0 K and $N^* = 0$, and for $E \geq E_{gn(gp)}$, there is an **[O-EP]-[E-OP] transition at a given E**, characterized by: $n^{N-MIT}_O = n^{N-MIT}_E$, $\varepsilon^{N-MIT}_{1\,O} = \varepsilon^{N-MIT}_{1\,E}$ and $R^{N-MIT}_O = R^{N-MIT}_E$, since, in this case, $E_{gn1(gp1)} = E_{gn2(gp2)} = E_{gn(gp)}$.

Then, by using Eq. (16b), from Equations (18, 19b, 19c, 19d), for $E \cong E_{gn(gp)}$, one can determine the exponential conduction (valence)-band tail states, due to those coefficients: $\sigma^{EBT}_{O[E]}\left(E \cong E_{gn(gp)}, N^* = N_{CDn(NDp)}\right)$, $\kappa^{EBT}_{O[E]}\left(E \cong E_{gn(gp)}, N^* = N_{CDn(NDp)}\right)$, $\varepsilon^{EBT}_{2O[2E]}\left(E \cong\right.$

$E_{gn(gp)}, N^* = N_{CDn(NDp)})$ and $\propto_{O[E]}^{EBT}(E \cong E_{gn(gp)}, N^* = N_{CDn(NDp)})$, and then their numerical results are given in **Table 5, reported in Appendix 1.**

Further, the numerical results of $n_{O[E]}(E)$, $\kappa_{O[E]}(E)$, $\varepsilon_{2O[2E]}(E)$ and $\varepsilon_{1\,O[E]}(E)$, are obtained by using Equations (17, 19b, 19c and 16), expressed as functions of N for (E=3.2 eV and T=20 K)-conditions, and as functions of T for (E=3.2 eV and $N = 10^{20}\,cm^{-3}$)-conditions, as those given in **Tables 6n, 6p, 7n and 7p, being reported in Appendix 1**, respectively.

Finally, for T=20K and $N = 10^{20}\,cm^{-3}$, and for given x and r_d, the numerical results of $\sigma_{O[E]}$ (E), $\varepsilon_{2O[2E]}(E)$ and $\propto_{O[E]}$ (E), are obtained by using Equations (18, 19c, 19d), and given in **Tables 8n and 8p, being reported in Appendix 1.**

In the following, we will determine the electrical-and-thermoelectric laws, by basing on our $\sigma_{O[E]}$-models, given in Eq. (20a).

OPTICAL [ELECTRICAL]-AND-THERMOELECTRIC PROPERTIES $[m^*_{n(p)} \equiv m_r(x)[m_{c(v)}(x)]]$

Here, if denoting, for majority electrons (holes), the thermal conductivity by $\sigma_{Th.\,O[E]}(N^*, r_{d(a)}, x, T)$ in $\frac{W}{cm \times K}$, and the Lorenz number L by: $L = \frac{\pi^2}{3} \times \left(\frac{k_B}{q}\right)^2 = 2.4429637 \left(\frac{W \times ohm}{K^2}\right) = 2.4429637 \times 10^{-8}\,(V^2 \times K^{-2})$, then the well-known Wiedemann-Frank law states that the ratio, $\frac{\sigma_{Th.O[E]}}{\sigma_{O[E]]}}$, due to the O-EP [E-OP], is proportional to the temperature T(K), as:

$$\frac{\sigma_{Th\,.O[E]}(N^*, r_{d(a)}, x, T)}{\sigma_{O[E]}(N^*, r_{d(a)}, x, T)} = L \times T. \qquad (21)$$

Further, the resistivity is found to be given by: $\rho_{O[E]}(N^*, r_{d(a)}, x, T) \equiv 1/\sigma_{O[E]}(N^*, r_{d(a)}, x, T)$, noting again that $N^* \equiv N - N_{CDn(NDp)}(r_{d(a)}, x)$.

In Eq. (20a), one notes that at T= 0 K, $\sigma_{O[E]}(N^*, r_{d(a)}, x, T)$ is proportional to $E^2_{Fno(Fpo)}$, or to $(N^*)^{\frac{4}{3}}$. Thus, from Eq. (21), one has: $\sigma_{O[E]}(N^* = 0, r_{d(a)}, x, T = 0K) = 0$ and also $\sigma_{Th.\,O[E]}(N^* = 0, r_{d(a)}, x, T = 0K) = 0$ at $N^* = 0$, at which the MIT occurs.

New Optical [Electrical] Coefficients

The relaxation time $\tau_{O[E]}$ is related to $\sigma_{O[E]}$ by [1]:

$\tau_{O[E]}(N^*, r_{d(a)}, x, T) \equiv \sigma_{O[E]}(N^*, r_{d(a)}, x, T) \times \frac{m^*_{n(p)}(x)\times m_o}{q^2\times(N^*/g_{c(v)})}$. Therefore, the mobility $\mu_{O[E]}$ is given by:

$$\mu_{O[E]]}\left(N^*, r_{d(a)}, x, T\right) = \frac{q\times\tau_{O[E]}\left(N^*, r_{d(a)}, x, T\right)}{m^*_{n(p)}(x)\times m_o} = \frac{\sigma_{O[E]}\left(N^*, r_{d(a)}, x, T\right)}{q\times(N^*/g_{c(v)})}\ \left(\frac{cm^2}{V\times s}\right) \quad (22a)$$

Here, one notes that, as given in Eq. (18), both $\mu_{O[E]]}$ and $\sigma_{O[E]}$, expressed in terms of $G_2(y) = \left(1+\frac{y^2}{3}\right) \equiv G_2(N, r_{d(a)}, x, T)$, with $y \equiv \frac{\pi}{\xi_{n(p)}}$, $\xi_{n(p)} = \xi_{n(p)}\left(N^*, r_{d(a)}, x, T\right)$, thus **increase** with **increasing T ≥ 0**, for given $(N, r_{d(a)}, x)$, since $\xi_{n(p)}$ **decreases** with increasing T.

Then, from the idea of Stokes, Einstein, Sutherland and Reynolds, we can define our viscosity coefficient, $\mathbb{V}_{O[E]}(N^*, r_{d(a)}, x, T)$, and its reduced one, $R\mathbb{V}_{O[E]}(N^*, r_{d(a)}, x, T)$, by:

$$\frac{\mathbb{V}_{O[E]}(N^*, r_{d(a)}, x, T)}{q} \equiv \frac{1}{6\pi\times\ \mu_{O[E]]}(N^*, r_{d(a)}, x, T)\times R_{WS}(N^*, x)}\left(\frac{V}{cm}\times\frac{s}{cm^2}\right), \quad R\mathbb{V}_{O[E]}(N^*, r_{d(a)}, x, T) \equiv \frac{\mathbb{V}_{O[E]}(N^*, r_{d(a)}, x, T)}{\mathbb{V}_{O[E]}(N^*, r_{d(a)}, x, T=0K)}, \quad (22b)$$

where $R_{WS}(N^*, x) \equiv \left(\frac{3g_{c(v)}(x)}{4\pi N^*}\right)^{1/3}$ is the effective Wigner-Seitz radius. Therefore, as noted above, for given $(N, r_{d(a)}, x)$, $\mathbb{V}_{O[E]}$ thus **decreases** with **increasing T**, since $\mu_{O[E]]}$ increases with increasing T, giving raise to: $\mathbf{R\mathbb{V}_{O[E]} \leq 1}$; thus, $\mathbf{R\mathbb{V}_{O[E]} = 1}$ at **T=0K.**

Then, taking into account above remarks, it is interesting to define the activation energy, $AE_{O[E]}(N^*, r_{d(a)}, x, T)$, as [17] by:

$$AE_{O[E]}\left(N^*, r_{d(a)}, x, T\right) \equiv k_B T \times Ln\left(R\mathbb{V}_{O[E]}\left(N^*, r_{d(a)}, x, T\right)\right) \leq 0\ eV, \text{ for } T \geq 0, \quad (22c)$$

according to the reduced activation energy, $RAE_{O[E]}\left(N^*, r_{d(a)}, x, T\right)$, defined by:

$$RAE_{O[E]}\left(N^*, r_{d(a)}, x, T\right) \equiv \frac{AE_{O[E]}\left(N^*, r_{d(a)}, x, T\right)}{k_B T} \equiv Ln\left(R\mathbb{V}_{O[E]}\left(N^*, r_{d(a)}, x, T\right)\right) \leq 0.$$

Furthermore, the Hall factor is defined by

$r_{HO[HE]}(N^*, r_{d(a)}, x, T) \equiv \frac{\langle\tau_{O[E]}{}^2\rangle_{FDDF}}{\left[\langle\tau_{O[E]}\rangle_{FDDF}\right]^2} = \frac{G_4(y)}{[G_2(y)]^2}$, $y \equiv \frac{\pi}{\xi_{n(p)}(N^*, r_{d(a)}, x, T)} = \frac{\pi k_B T}{E_{Fn(Fp)}(N^*, r_{d(a)}, x, T)}$, and therefore, the Hall mobility yields:

$\mu_{HO[HE]}\left(N^*, r_{d(a)}, x, T\right) \equiv \mu_{O[E]}\left(N^*, r_{d(a)}, x, T\right) \times r_{HO[HE]}\left(N^*, r_{d(a)}, x, T\right) \left(\frac{cm^2}{V\times s}\right)$, (23)

noting that, at T=0K, since $r_{HE[HO]}(N^*, r_{d(a)}, x, T) = 1$, one therefore gets: $\mu_{HO[HE]}(N^*, r_{d(a)}, x, T) \equiv \mu_{O[E]}(N^*, r_{d(a)}, x, T)$.

Van-Cong (VC)-relation between the diffusion, the mobility and the viscosity

By taking into account Equations (22a, 22b), our relation is found to be defined by [1]:

$$\mathbb{R}_{E[O](VC)}\left(N^*, r_{d(a)}, x, T\right) \equiv \frac{D_{O[E]}\left(N^*, r_{d(a)}, x, T\right)}{\mu_{O[E]}\left(N^*, r_{d(a)}, x, T\right)} \equiv D_{O[E]}\left(N^*, r_{d(a)}, x, T\right) \times \frac{\mathbb{V}_{O[E]}\left(N^*, r_{d(a)}, x, T\right)}{q} \times 6\pi \times R_{WS}(N^*, x) \equiv \frac{N^*}{q} \times \frac{dE_{Fn(Fp)}}{dN^*} \equiv \frac{k_B\times T}{q} \times \left(u\frac{d\xi_{n(p)}(u)}{du}\right) = \sqrt{\frac{3\times L}{\pi^2}} \times T \times \left(u\frac{d\xi_{n(p)}(u)}{du}\right),$$

$$\frac{k_B}{q} = \sqrt{\frac{3\times L}{\pi^2}}, \quad (24)$$

where $D_{E[O]}\left(N^*, r_{d(a)}, x, T\right)$ is the diffusion coefficient, $\xi_{n(p)}(u)$ is defined in Eq. (11), the mobility $\mu_{O[E]}\left(N^*, r_{d(a)}, x, T\right)$ is determined in Eq. (22a), and finally the viscosity coefficient $\mathbb{V}_{O[E]}(N^*, r_{d(a)}, x, T)$ is defined in Eq. (22b).Then, by differentiating this function $\xi_{n(p)}(u) \equiv \frac{E_{Fn(Fp)}(u)}{k_B T} = \frac{G(u)+Au^B F(u)}{1+Au^B} \equiv \frac{V(u)}{W(u)}$, with respect to u, being defined in Eq. (11), one thus obtains $\frac{d\xi_{n(p)}(u)}{du}$.

Therefore, Eq. (24) can also be rewritten as: $\mathbb{R}_{E[O](VC)}(u) = \frac{k_B\times T}{q} \times u\frac{V'(u)\times W(u) - V(u)\times W'(u)}{W^2(u)}$

where $W'(u) = ABu^{B-1}$ and $V'(u) = u^{-1} + 2^{-\frac{3}{2}}e^{-du}(1-du) + \frac{2}{3}Au^{B-1}F(u)\left[\left(1+\frac{3B}{2}\right) + \frac{4}{3}\times\frac{bu^{-\frac{4}{3}}+2cu^{-\frac{8}{3}}}{1+bu^{-\frac{4}{3}}+cu^{-\frac{8}{3}}}\right]$. One remarks that: (i) as $u \to 0$, one has: $W^2 \simeq 1$ and $u[V' \times W - V \times W'] \simeq 1$, and therefore: $\mathbb{R}_{E[O](VC)}(u \to 0) \simeq \frac{k_B\times T}{q}$, being a well-known relation given by Stokes, Einstein, Sutherland and Reynolds, and (ii) as $u \to \infty$, one has: $W^2 \approx A^2u^{2B}$ and $u[V' \times W - V \times W'] \approx \frac{2}{3}au^{2/3}A^2u^{2B}$, and therefore, in this **highly degenerate case** and at T=0K, our relation (24) is reduced to: $\mathbb{R}_{E[O](VC)}\left(N^*, r_{d(a)}, x, T = 0K\right) \approx \frac{2}{3}E_{Fno(Fpo)}(N^*)/q$. In other words, **Eq. (24) verifies the correct limiting conditions**.

Furthermore, in the present degenerate case ($u \gg 1$), Eq. (24) can be rewritten as:

$$\mathbb{R}_{E[O](VC)}\left(N^*, r_{d(a)}, x, T=0K\right) \simeq \frac{2}{3} \times \frac{E_{Fno(Fpo)}(u)}{q} \times \left[1 + \frac{4}{3} \times \frac{\left(bu^{-\frac{4}{3}}+2cu^{-\frac{8}{3}}\right)}{\left(1+bu^{-\frac{4}{3}}+cu^{-\frac{8}{3}}\right)}\right],$$

where $a = \left[3\sqrt{\pi}/4\right]^{2/3}$, $b = \frac{1}{8}\left(\frac{\pi}{a}\right)^2$ and $c = \frac{62.3739855}{1920}\left(\frac{\pi}{a}\right)^4$.

Then, in **Tables 9n and 9p, reported in Appendix 1**, the numerical results of $\sigma_{O[E]}$, $\mu_{O[E]}$ and $D_{O[E]}$ for given x and T, are obtained by using Equations (20a, 22a and 24), suggesting that, for a given N, they decrease [decrease], with increasing $r_{d(a)}$.

Further, in **Tables 10n and 10p, reported in Appendix 1**, the numerical results of the viscosity coefficient $\mathbb{V}_{O[E]}(N^*, r_{d(a)}, x, T)$, expressed in $\left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$, are obtained by using Eq. (22b), suggesting that: (i) for given (x, T and N), they increase with increasing $r_{d(a)}$, (ii) for given (x, $r_{d(a)}$ and N) they decrease with increasing T, in good agreement with those, obtained in liquids by Ewell and Eyring[17] and complex fluids by Wenhao[18], and finally (iii) for given (x, T and $r_{d(a)}$) they increase with increasing N, in good agreement with those, obtained in complex fluids by Wenhao.[18]

Furthermore, in **Tables 11n and 11p, reported in Appendix 1**, the numerical results of the numerical results of reduced Fermi energy $\xi_{n(p)}(N^*, r_d, x, T)$, viscosity coefficient $\mathbb{V}_{O[E]}(N^*, r_d, x, T)$, and activation energy $AE_{O[E]}(N^*, r_d, x, T)$ are obtained by using Equations (11, 22b, 22c), respectively, suggesting that, for given (x, $r_{d(a)}$ and N), and with increasing T, both $\mathbb{V}_{O[E]}$ and $AE_{O[E]}$ decrease, since $\xi_{n(p)}$ decreases, in good agreement with those, obtained in liquids by Ewell and Eyring[17] and complex fluids by Wenhao.[18]

Thermoelectric Coefficients

Here, as noted above, $E_{Fn(Fp)}\left(m_r(x)\right) > E_{Fn(Fp)}\left(m_{c(v)}(x)\right)$ or $\xi_{n(p)}\left(m_r(x)\right) > \xi_{n(p)}\left(m_{c(v)}(x)\right)$ for a given T, since $m_r(x) < m_{c(v)}(x)$ for given x, corresponding to: $\sigma_O\left(m_r(x)\right) > \sigma_E\left(m_{c(v)}(x)\right)$.

Then, from Eq. (20a), obtained for $\sigma_{O[E]}(N^*, r_{d(a)}, x, T)$, the well-known Mott definition for the thermoelectric power or for the Seebeck coefficient, $S_{E[O]}$, is found to be given by:

$$S_{O[E]}(N^*, r_{d(a)}, x, T) \equiv \frac{-\pi^2}{3} \times \frac{k_B}{q>0} \times k_B T \times \frac{\partial \ln\sigma_{O[E]}}{\partial E}\Big]_{E=E_{Fn(Fp)}} = \frac{-\pi^2}{3} \times \frac{k_B}{q} \times \frac{\partial \ln\sigma_{O[E]}(\xi_{n(p)})}{\partial \xi_{n(p)}}.$$

Then, using Eq. (11), for the degenerate case, $\xi_{n(p)} \geq 0$, one gets, by putting

$$Y_{Sb\ O[E]}(N^*, r_{d(a)}, x, T) \equiv \left[1 - \frac{y^2}{3\times G_2\left(y=\frac{\pi}{\xi_{n(p)}}\right)}\right],$$

$$S_{O[E]}(N^*, r_{d(a)}, x, T) \equiv \frac{-\pi^2}{3} \times \frac{k_B}{q} \times \frac{2Y_{Sb\ O[E]}(N^*, r_{d(a)}, x, T)}{\xi_{n(p)}} = -\sqrt{\frac{3\times L}{\pi^2}} \times \frac{2\times\xi_{n(p)}}{\left(1+\frac{3\times\xi_{n(p)}{}^2}{\pi^2}\right)} = -2\sqrt{L} \times \frac{\sqrt{ZT_{O[E]Mott}}}{1+ ZT_{O[E]Mott}} \left(\frac{V}{K}\right) < 0, \quad ZT_{O[E]Mott} = \frac{\pi^2}{3\times\xi^2_{n(p)}}, \ (25)$$

according to:

$$\frac{\partial\, S_{O[E]}}{\partial\xi_{n(p)}} = \sqrt{\frac{3\times L}{\pi^2}} \times 2 \times \frac{\frac{3\times\xi_{n(p)}{}^2}{\pi^2} - 1}{\left(1+\frac{3\times\xi_{n(p)}{}^2}{\pi^2}\right)^2} = \sqrt{\frac{3\times L}{\pi^2}} \times 2 \times \frac{ZT_{O[E]Mott}\times[1- ZT_{O[E]Mott}]}{[1+ ZT_{O[E]Mott}]^2}.$$

Here, one notes that: (i) as $\xi_{n(p)} \to +\infty$ or $\xi_{n(p)} \to +0$, one has a same limiting value of $S_{O[E]}$: $S_{O[E]} \to -0$, (ii) at $\xi_{n(p)} = \sqrt{\frac{\pi^2}{3}} \simeq 1.8138$, since $\frac{\partial\, S_{O[E]}}{\partial\xi_{n(p)}} = 0$, one therefore gets: a minimum $\left(S_{O[E]}\right)_{min.} = -\sqrt{L} \simeq -1.563 \times 10^{-4} \left(\frac{V}{K}\right)$, and (iii) at $\xi_{n(p)} = 1$ one obtains: $S_{O[E]} \simeq -1.322 \times 10^{-4} \left(\frac{V}{K}\right)$.

Further, the figure of merit is found to be defined by:

$$ZT_{O[E]}(N^*, r_{d(a)}, x, T) \equiv \frac{S^2\times\sigma_{O[E]}\times T}{\kappa} = \frac{S^2}{L} = \frac{4\times ZT_{O[E]Mott}}{[1+ ZT_{O[E]Mott}]^2}.$$

(26)

Here, one notes that: (i) $\frac{\partial(ZT_{O[E]})}{\partial\xi_{n(p)}} = 2 \times \frac{S_{O[E]}}{L} \times \frac{\partial\, S_{O[E]}}{\partial\xi_{n(p)}}$, $S_{E[O]} < 0$, (ii) at $\xi_{n(p)} = \sqrt{\frac{\pi^2}{3}} \simeq 1.8138$, since $\frac{\partial(ZT_{O[E]})}{\partial\xi_{n(p)}} = 0$, one gets: a maximum $\left(ZT_{O[E]}\right)_{max.} = 1$, $ZT_{O[E]Mott} = 1$, and (iii) at $\xi_{n(p)} = 1$, one obtains: $ZT_{O[E]} \simeq 0.715$ and $ZT_{O[E]Mott} = \frac{\pi^2}{3} \simeq 3.290$.

Finally, the first Van-Cong coefficient can be defined by:

$$VC1_{O[E]}(N^*, r_{d(a)}, x, T) \equiv -N^* \times \frac{d\, S_{O[E]}}{dN^*} \left(\frac{V}{K}\right) = N^* \times \frac{\partial\, S_{O[E]}}{\partial\xi_{n(p)}} \times -\frac{\partial\xi_{n(p)}}{\partial N^*}, \qquad (27)$$

being equal to 0 for $\xi_{n(p)} = \sqrt{\frac{\pi^2}{3}}$,

and the second Van-Cong coefficient as:

$\mathrm{VC2}_{O[E]}(N^*, r_{d(a)}, x, T) \equiv T \times \mathrm{VC1}_{O[E]}$ (V), (28).

the Thomson coefficient, Ts , by:

$$\mathrm{Ts}_{O[E]}(N^*, r_{d(a)}, x, T) \equiv T \times \frac{d\, S_{O[E]}}{dT}\left(\frac{V}{K}\right) = T \times \frac{\partial\, S_{O[E]}}{\partial \xi_{n(p)}} \times \frac{\partial \xi_{n(p)}}{\partial T}, \qquad (29)$$

being equal to 0 for $\xi_{n(p)} = \sqrt{\frac{\pi^2}{3}}$,

and the Peltier coefficient, $\mathrm{Pt}_{E[O]}$, as:

$\mathrm{Pt}_{O[E]}(N^*, r_{d(a)}, x, T) \equiv T \times S_{O[E]}$ (V). (30)

Then, in **Tables 12n and 12p, reported in Appendix 1**, the numerical results of various thermoelectric coefficients such as: $\sigma_{Th.O[E]}$, $S_{O[E]}$, $\mathrm{VC1}_{O[E]}$, $\mathrm{VC2}_{O[E]}$, $\mathrm{Ts}_{O[E]}$, $\mathrm{Pt}_{O[E]}$ and $\mathrm{ZT}_{O[E]}$, for given x, $r_{d(a)}$, T=(3K and 80K) and N, are obtained by using Equations (21, 25, 27, 28, 29, 30 and 26), respectively.

In summary, in the O-EP [E-OP] and for given physical conditions: x, $r_{d(a)}$, N (or T), the same values of $\xi_{n(p)}$ decrease, according to the increasing T (or to the decreasing N), since $\mathrm{VC1}_{O[E]}(N, r_{d(a)}, x, T)$ and $\mathrm{Ts}_{O[E]}(N, r_{d(a)}, x, T)$ are expressed in terms of $\frac{-d\, S_{O[E]}}{dN^*}$ and $\frac{d\, S_{O[E]}}{dT}$, one has: $[\mathrm{VC1}_{O[E]}, \mathrm{Ts}_{O[E]}] < 0$ for $\xi_{n(p)} > \sqrt{\frac{\pi^2}{3}}$, $[\mathrm{VC1}_{O[E]}, \mathrm{Ts}_{O[E]}] = 0$ for $\xi_{n(p)} = \sqrt{\frac{\pi^2}{3}}$, and $[\mathrm{VC1}_{O[E]}, \mathrm{Ts}_{O[E]}] > 0$ for $\xi_{n(p)} < \sqrt{\frac{\pi^2}{3}}$, stating that for $\boldsymbol{\xi_{n(p)} = \sqrt{\frac{\pi^2}{3}} \simeq 1.8138}$: $S_{O[E]}$, determined in Eq. (25), thus presents **a same minimum** $S_{O[E]\,min.} = -\sqrt{L} \simeq -1.563 \times 10^{-4}\left(\frac{V}{K}\right)$, and $\mathrm{ZT}_{O[E]}$, determined in Eq. (26), therefore presents **a same maximum:** $\mathrm{ZT}_{O[E]\,max.} = 1$, and $(\mathrm{ZT})_{Mott} = 1$. Furthermore, for $\boldsymbol{\xi_{n(p)} = 1}$, the numerical results of $S_{O[E]}$, $\mathrm{ZT}_{O[E]}$, $\mathrm{ZT}_{O[E]Mott}$, $\mathrm{VC1}_{E[O]}$, and $\mathrm{Ts}_{O[E]}$, present the **same results**: $-1.322 \times 10^{-4}\frac{V}{K}$, 0.715, 3.290, $1.105 \times 10^{-4}\frac{V}{K}$, and $1.657 \times 10^{-4}\frac{V}{K}$, respectively, as those observed in [4, 5], and those given in **Table 13, reported in Appendix 1.**

It seems that these same obtained results could represent **a new law for the thermoelectric properties, obtained in the degenerate case ($\xi_{n(p)} \geq 0$)**.

Furthermore, it is interesting to remark that the $VC2_{O[E]}$-coefficient is related to our generalized Einstein relation (24) by:

$$\frac{k_B}{q} \times VC2_{O[E]}\left(N^*, r_{d(a)}, x, T\right) \equiv -\frac{\partial\, S_{O[E]}}{\partial \xi_{n(p)}} \times \frac{D_{O[E]}\left(N^*, r_{d(a)}, x, T\right)}{\mu_{O[E]}\left(N^*, r_{d(a)}, x, T\right)} \left(\frac{V^2}{K}\right), \quad \frac{k_B}{q} = \sqrt{\frac{3 \times L}{\pi^2}}, \quad (31)$$

according, in this work, with the use of our Eq. (25), to:

$$VC2_{O[E]}\left(N, r_{d(a)}, x, T\right) \equiv -\frac{D_{O[E]}\left(N^*, r_{d(a)}, x, T\right)}{\mu_{O[E]}\left(N^*, r_{d(a)}, x, T\right)} \times 2 \times \frac{ZT_{O[E]Mott} \times \left[1 - ZT_{O[E]Mott}\right]}{\left[1 + ZT_{O[E]Mott}\right]^2} \quad (V).$$

Of course, our relation (31) is reduced to: $\frac{D_{O[E]}}{\mu_{O[E]}}$, $VC1_{O[E]}$ and $VC2_{O[E]}$, being determined respectively by Equations (24, 27, 28). This may be a new result.

CONCLUDING REMARKS

In the $\mathbf{n^+(p^+) - X(x)}$-crystalline alloy, $0 \leq x \leq 1$, x being the concentration, the optical, electrical and thermoelectric coefficients, enhanced by : **(i)** the optico-electrical phenomenon (O-EP) and the electro-optical phenomenon (E-OP), **(ii)** our static dielectric constant law, $\varepsilon(r_{d(a)}, x)$, $r_{d(a)}$ being the donor (acceptor) d(a)-radius, given in Equations (1a, 1b), **(iii)** our accurate reduced Fermi energy, $\xi_{n(p)}$, given in Eq. (11), accurate with a precision of the order of $\mathbf{2.11} \times \mathbf{10^{-4}}$[9], affecting all the expressions of optical, electrical and thermoelectric coefficients, and **(iv)** our optical-and-electrical conductivity models, given in Eq. (18, 20a), are now investigated by basing on our physical model and Fermi-Dirac distribution function, as those given in our recent works.[1-5]

Some important concluding remarks can be given and discussed as follows.

(I)-First of all, one notes that from Equations (3, 6a, 6b) the MIT occurs as T=0 K and $N^*(N, r_{d(a)}, x) \equiv N - N_{CDn(NDp)}(r_{d(a)}, x\ ) \cong N - N^{EBT}_{CDn(CDp)}\left(r_{d(a)}, x\right) = 0$, according, for $E \geq E_{gn(gp)}$, to: $E_{Fno(Fpo)}(N^* = 0) \equiv \frac{\hbar^2 \times k^2_{Fn(Fp)}(N^*)}{2 \times m^*_{n(p)}(x) \times m_o} = 0$, and $\kappa^{MIT}_{O[E]}(E, N^* = 0) = 0$, $\varepsilon^{MIT}_{2\,O[E]}(E, N^* = 0) = 0$, $\sigma^{MIT}_{O[E]}(E, N^* = 0) = 0$ and $\propto^{MIT}_{O[E]}(E, N^* = 0)$=0, since, for example, $\sigma_{E[O]}(E, N^* = 0)$ is proportional to $E^2_{Fno(Fpo)}$, or to $(N^* = 0)^{\frac{4}{3}} = 0$. But, for such the same physical conditions: T=0 K, $N^* = 0$ and $E \geq E_{gn(gp)}$, we obtain other numerical results such as: $n^{N-MIT}_{O[E]}(N^* = 0, E) \neq 0$, $\varepsilon^{N-MIT}_{1\,O[E]}(N^* = 0, E) \neq 0$ and $R^{N-MIT}_{O[E]}(N^* = 0, E) \neq 0$, for

$E \geq E_{gn(gp)}$, according to the non-MIT (N-MIT), as showed in Tables 3, 4n and 4p, reported in Appendix 1. These Tables also state that, at T=0 K and $N^* = 0$, and for $E \geq E_{gn(gp)}$, there is an [O-EP]-[E-OP] transition at a given E, as: $n_O^{N-MIT} = n_E^{N-MIT}$, $\varepsilon_{1\,O}^{N-MIT} = \varepsilon_{1\,E}^{N-MIT}$ and $R_O^{N-MIT} = R_E^{N-MIT}$, since, in this case, $E_{gn1(gp1)} = E_{gn2(gp2)} = E_{gn(gp)}$.

Then, by using Eq. (16b), from Equations (18, 19b, 19c, 19d), for $E \cong E_{gn(gp)}$, one can determine the exponential conduction (valence)-band tail states, due to those coefficients: $\sigma_{O[E]}^{EBT}(E \cong E_{gn(gp)}, N^* = N_{CDn(NDp)})$, $\kappa_{O[E]}^{EBT}(E \cong E_{gn(gp)}, N^* = N_{CDn(NDp)})$, $\varepsilon_{2O[2E]}^{EBT}(E \cong E_{gn(gp)}, N^* = N_{CDn(NDp)})$ and $\propto_{O[E]}^{EBT}(E \cong E_{gn(gp)}, N^* = N_{CDn(NDp)})$, and then their numerical results are given in Table 5, reported in Appendix 1.

(II)-Further, the numerical results of $n_{O[E]}(E)$, $\kappa_{O[E]}(E)$, $\varepsilon_{2O[2E]}(E)$ and $\varepsilon_{1\,O[E]}(E)$, are obtained by using Equations (17, 19b, 19c and 16), expressed as functions of N for (E=3.2 eV and T=20 K)-conditions, and as functions of T for (E=3.2 eV and $N = 10^{20}\,cm^{-3}$)-conditions, as those given in Tables 6n, 6p, 7n and 7p, being reported in Appendix 1, respectively.

Finally, for T=20K and $N = 10^{20}\,cm^{-3}$, and for given x and r_d, the numerical results of $\sigma_{O[E]}$ (E), $\varepsilon_{2O[2E]}(E)$ and $\propto_{O[E]}$ (E), are obtained by using Equations (18, 19c, 19d), and given in Tables 8n and 8p, being reported in Appendix 1.

(III)- Then, in Tables 9n and 9p, reported in Appendix 1, the numerical results of $\sigma_{O[E]}$, $\mu_{O[E]}$ and $D_{O[E]}$ for given x and T, are obtained by using Equations (20a, 22a and 24), suggesting that, for a given N, they decrease [decrease], with increasing $r_{d(a)}$.

In Tables 10n and 10p, reported in Appendix 1, the numerical results of the viscosity coefficient $\mathbb{V}_{O[E]}(N^*, r_{d(a)}, x, T)$, expressed in $\left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$, are obtained by using Eq. (22b), suggesting that: (i) for given (x, T and N), they increase with increasing $r_{d(a)}$, (ii) for given (x, $r_{d(a)}$ and N) they decrease with increasing T, in good agreement with those, obtained in liquids by Ewell and Eyring[17] and complex fluids by Wenhao[18], and finally (iii) for given (x, T and $r_{d(a)}$) they increase with increasing N, in good agreement with those, obtained in complex fluids by Wenhao.[18]

Further, in Tables 11n and 11p, reported in Appendix 1, the numerical results of the numerical results of reduced Fermi energy $\xi_{n(p)}(N^*, r_d, x, T)$, viscosity coefficient $\mathbb{V}_{O[E]}(N^*, r_d, x, T)$, and activation energy $AE_{O[E]}(N^*, r_d, x, T)$ are obtained by using Equations (11, 22b, 22c), respectively, suggesting that, for given (x, $r_{d(a)}$ and N), $\mathbb{V}_{O[E]}$ and $AE_{O[E]}$ both decrease with increasing T, since $\xi_{n(p)}$ decreases with increasing T, in good agreement with those, obtained in liquids by Ewell and Eyring[17] and complex fluids by Wenhao[18], suggesting a same viscosity phenomenon observed in both liquids and solids.

Furthermore, in Tables 12n and 12p, reported in Appendix 1, the numerical results of various thermoelectric coefficients such as: $\sigma_{Th.O[E]}$, $S_{O[E]}$, $VC1_{O[E]}$, $VC2_{O[E]}$, $Ts_{O[E]}$, $Pt_{O[E]}$ and $ZT_{O[E]}$, for given x, $r_{d(a)}$, T=(3K and 80K) and N, are obtained by using Equations (21, 25, 27, 28, 29, 30 and 26), respectively.

(IV)-Finally, from Equations (20a, 21-30), for any given x, $r_{d(a)}$ and N (or T), with increasing T (or decreasing N), one obtains: (i) for $\xi_{n(p)} = \sqrt{\frac{\pi^2}{3}} \simeq 1.8138$, while the numerical results of $S_{O[E]}$ present a same minimum $S_{O[E]\,min.}\left(\simeq -1.563 \times 10^{-4}\,\frac{V}{K}\right)$, those of $ZT_{O[E]}$ show a same maximum $ZT_{ET[OT]\,max.} = 1$, (ii) for $\xi_{n(p)} = 1$, the numerical results of $S_{O[E]}$, $ZT_{O[E]}$, $ZT_{O[E]\,Mott}$, $VC1_{O[E]]}$, and $Ts_{O[E]}$, present the same results: $-1.322 \times 10^{-4}\,\frac{V}{K}$, 0.715, 3.290, $1.105 \times 10^{-4}\,\frac{V}{K}$, and $1.657 \times 10^{-4}\,\frac{V}{K}$, respectively, and finally (iii) for $\xi_{n(p)} = \sqrt{\frac{\pi^2}{3}} \simeq 1.8138$, $ZT_{O[E]\,Mott} = 1$, as those given in Table 13, reported in Appendix 1.

REFERENCES

1. Van Cong, H. Optical, Electrical and Thermoelectric Laws in n(p)-Type Degenerate “Compensated” GaAs(1-x)Sb(x)-Crystalline Alloy, Enhanced by: Optico-Electrical Phenomenon and Electro-Optical Phenomenon, and our Static Dielectric Constant Law, Accurate Fermi Energy and Conductivity Models (IX). WJERT, 2025; 11(12): 221-269.
2. Van Cong, H. Optical, Electrical and Thermoelectric Laws in n(p)-Type Degenerate “Compensated” GaAs(1-x)Te(x)-Crystalline Alloy, Enhanced by: Optico-Electrical Phenomenon and Electro-Optical Phenomenon, and our Static Dielectric Constant Law, Accurate Fermi Energy and Conductivity Models (VIII). WJERT, 2025; 11(12): 180-220.

3. Van Cong, H. New Critical Impurity Density in Metal-Insulator Transition, obtained in Various n(p)-Type Degenerate Crystalline Alloys, being just That of Carriers Localized in Exponential Band Tails (II). WJERT, 2024; 10(4): 65-96.
4. Van Cong H. Same maximum figure of merit ZT(=1), due to the effect of impurity size, obtained in the n(p)-type degenerate Ge -crystal ($\xi_{n(p)} \geq 1$), at same reduced Fermi energy $\xi_{n(p)} = \sqrt{\frac{\pi^2}{3}} \simeq 1.8138$, same minimum Seebeck coefficient $(S)_{min.} \left(\simeq -1.563 \times 10^{-4} \frac{V}{K}\right)$, same maximum $(ZT)_{max.} = 1$, and same $(ZT)_{Mott} \left(= \frac{\pi^2}{3\xi_{n(p)}^2} = 1\right)$, SCIREA Journal of Physics. 2023; 8(4): 407-430.
5. Van Cong, H. Same Maximal Figure of Merit ZT(=1), Due to the Effect of Impurity Size, Obtained in the n(p)-Type Degenerate GaAs-Crystal ($\xi_{n(p)} \geq 1$), at Same Reduced Fermi Energy $\xi_{n(p)}(= 1.8138)$ and Same Minimum Seebeck Coefficient $S\left(= -1.563 \times 10^{-4} \frac{V}{K}\right)$, at which Same $(ZT)_{Mott} \left(= \frac{\pi^2}{3\xi_{n(p)}^2} = 1\right)$. SCIREA Journal of Physics, 2023; 8(2): 133-157.
6. Van Cong, H. Effects of donor size and heavy doping on optical, electrical and thermoelectric properties of various degenerate donor-silicon systems at low temperatures. American Journal of Modern Physics, 2018; 7(4): 136-165.
7. Kim, H. S. et al. Characterization of Lorenz number with Seebeck coefficient measurement. APL Materials, 2015; 3(4): 041506.
8. Hyun, B. D. et al. Electrical-and-Thermoelectric Properties of $90\%Bi_2Te_3 - 5\%Sb_2Te_3 - 5\%Sb_2Se_3$ Single Crystals Doped with SbI_3. Scripta Materialia, 1998; 40(1): 49-56.
9. Van Cong, H. and Debiais, G. A simple accurate expression of the reduced Fermi energy for any reduced carrier density. J. Appl. Phys., 1993; 73: 1545-1546.
10. Van Cong, H. et al. Size effect on different impurity levels in semiconductors. Solid State Communications, 1984; 49: 697-699.
11. Van Cong, H. Diffusion coefficient in degenerate semiconductors. Phys. Stat. Sol. (b), 1984; 101: K27.
12. Van Cong, H. and Doan Khanh, B. Simple accurate general expression of the Fermi-Dirac integral $F_j(a)$ and for j> -1. Solid-State Electron., 1992; 35(7): 949-951.

13. Van Cong, H. New series representation of Fermi-Dirac integral $F_j(-\infty < a < \infty)$ for arbitrary j> -1, and its effect on $F_j(a \geq 0_+)$ for integer j≥ 0. Solid-State Electron., 1991; 34(5): 489-492.
14. Van Cong, H. and G. Mesnard. Thermoelectric effects of heavily doped semiconductors at low temperatures. Phys. Stat. Sol. (b), 1972; 50(1): 53-58.
15. Van Cong, H. Fermi energy and band-tail parameters in heavily doped semiconductors. Journal of Physics and Chemistry of Solids, 1975; 36(11): 1237-1240.
16. Van Cong, H. Quantum efficiency and radiative lifetime in degerate n-type GaAs. Journal of Physics and Chemistry of Solids, 1981; 36(11): 95-99.
17. Ewell, R. H and Eyring, H. Theory of the Viscosity of Liquids as a Function of Temperature and Pressure. Journal of Chemical Physics, 1937; 5: 726-736.
18. Wenhao, Z. Influence of Temperature and Concentration on Viscosity of Complex Fluids. Journal of Physics: Conference Series 1965, 2021; 1: 012064.

APPENDIX 1

Table 1: In the $X(x) \equiv GaTe_{1-x}P_x$-crystalline alloy, the different values of energy-band-structure parameters, for a given x, are given in the following [3].

In the $\mathbf{X(x)}$**-crystalline alloy**, in which $r_{do(ao)} = r_{Te(Ga)}$ =0.132 nm (0.126 nm), we have [3]: $g_{c(v)}(x) = 1 \times x + 1 \times (1-x) = 1$, $m_{c(v)}(x)/m_o = 0.13\ (0.5) \times x + 0.209\ (0.4) \times (1-x)$, $\varepsilon_o(x) = 11.1 \times x + 12.3 \times (1-x)$, $E_{go}(x) = 1.796 \times x + 1.796 \times (1-x)$.

Table 2: Expressions for $G_{p>1}(y \equiv \frac{\pi}{\xi_{n(p)}})$, due to the Fermi-Dirac distribution function, are used to determine the electrical-and-thermoelectric coefficients, suggesting that, for given (N^*, r_d, x), since $\xi_{n(p)}(N^*, r_d, x, T)$ decreases with increasing T, $G_{p>1}(T)$ increases.

$G_{3/2}(y)$	$G_2(y)$	$G_{5/2}(y)$	$G_3(y)$	$G_{7/2}(y)$	$G_4(y)$	$G_{9/2}(y)$
$\left(1+\frac{y^2}{8}+\frac{7y^4}{640}\right)$	$\left(1+\frac{y^2}{3}\right)$	$\left(1+\frac{5y^2}{8}-\frac{7y^4}{384}\right)$	$(1+y^2)$	$\left(1+\frac{35y^2}{24}+\frac{49y^4}{384}\right)$	$\left(1+2y^2+\frac{7y^4}{15}\right)$	$\left(1+\frac{21y^2}{8}+\frac{147y^4}{128}\right)$

Table 3. For T=0K and N=$N_{CDn(CDp)}(r_{d(a)}, x)$, and at $E = E_{gn(gp)}$, the numerical results of $n_{O[E]}^{N-MIT}$, $\varepsilon_{1\ O[E]}^{N-MIT}$ and $R_{O[E]}^{N-MIT}$ are obtained, using Equations (17, 16a), suggesting that they decrease (↘) with increasing (↗) $r_{d(a)}$ and $E_{gn(gp)}$, and further they are found to be the same, for given $r_{d(a)}$ and $E_{gn(gp)}$, since $E_{gn1(gp1)} = E_{gn2(gp2)} = E_{gn(gp)}$.

Donor		P	**Te**	Sb	Sn
r_d (nm) [4]	↗	0.110	**0.132**	0.136	0.140
At **x=0**,					
E_{gn}(meV)	↗	1791.7 [1791.7]	1796.0 [1796.0]	1796.2 [1796.2]	1796.6 [1796.6]
$n_{O[E]}^{N-MIT}$	↘	3.315 [3.315]	3.177 [3.177]	3.173 [3.173]	3.160 [3.160]
$\varepsilon_{1\ O[E]}^{N-MIT}$	↘	10.99 [10.99]	10.09 [10.09]	10.07 [10.07]	9.987 [9.987]
$R_{O[E]}^{N-MIT}$	↘	0.288 [0.288]	0.272 [0.272]	0.271 [0.271]	0.270 [0.270]
At **x=0.5**,					
E_{gn}(meV)	↗	1792.1 [1792.1]	1796.0 [1796.0]	1796.1 [1796.1]	1796.6 [1796.6]
$n_{O[E]}^{N-MIT}$	↘	3.262 [3.262]	3.128 [3.128]	3.124 [3.124]	3.111 [3.111]
$\varepsilon_{1\ O[E]}^{N-MIT}$	↘	10.64 [10.64]	9.785 [9.785]	9.759 [9.759]	9.681 [9.681]
$R_{O[E]}^{N-MIT}$	↘	0.282 [0.282]	0.266 [0.266]	0.265 [0.265]	0.264 [0.264]
At **x=1**,					
E_{gn}(meV)	↗	1792.1 [1792.1]	1796.0 [1796.0]	1796.1 [1796.1]	1796.5 [1796.5]

$n^{N-MIT}_{O[E]}$	↘	3.262 [3.262]	3.077 [3.077]	3.073 [3.073]	3.061 [3.061]
$\varepsilon^{N-MIT}_{1\ O[E]}$	↘	10.64 [10.64]	9.471 [9.471]	9.447 [9.447]	9.373 [9.373]
$R^{N-MIT}_{O[E]}$	↘	0.282 [0.282]	0.260 [0.260]	0.259 [0.259]	0.258 [0.258]

Acceptor		Ga	Mg	In	Cd
r_a (nm)	↗	**0.126**	0.140	0.144	0.148
At **x=0**,					
E_{gp}(meV)	↗	1796 [1796]	1800.2 [1800.2]	1803.1 [1803.1]	1806.8 [1806.8]
$n^{N-MIT}_{O[E]}$	↘	3.177 [3.177]	3.120 [3.120]	3.085 [3.085]	3.044 [3.044]
$\varepsilon^{N-MIT}_{1\ O[E]}$	↘	10.09 [10.09]	9.735 [9.735]	9.518 [9.518]	9.267 [9.267]
$R^{N-MIT}_{O[E]}$	↘	0.272 [0.272]	0.265 [0.265]	0.260 [0.260]	0.255 [0.255]
At **x=0.5**,					
E_{gp}(meV)	↗	1796 [1796]	1801.2 [1801.2]	1804.8 [1804.8]	1809.4 [1809.4]
$n^{N-MIT}_{O[E]}$	↘	3.128 [3.128]	3.071 [3.071]	3.037 [3.037]	2.996 [2.996]
$\varepsilon^{N-MIT}_{1\ O[E]}$	↘	9.785 [9.785]	9.434 [9.434]	9.223 [9.223]	8.979 [8.979]
$\varepsilon^{N-MIT}_{1\ O[E]}$	↘	0.266 [0.266]	0.259 [0.259]	0.255 [0.255]	0.249 [0.249]
At **x=1**,					
E_{gp}(meV)	↗	1796 [1796]	1802.5 [1802.5]	1806.9 [1806.9]	1812.5 [1812.5]
$n^{N-MIT}_{O[E]}$	↘	3.077 [3.077]	3.022 [3.022]	2.987 [2.987]	2.947 [2.947]
$\varepsilon^{N-MIT}_{1\ O[E]}$	↘	9.471 [9.471]	9.131 [9.131]	8.925 [8.925]	8.686 [8.686]
$\varepsilon^{N-MIT}_{1\ O[E]}$	↘	0.260 [0.260]	0.253 [0.253]	0.248 [0.248]	0.243 [0.243]

Table 4n. For T=0K and N=$N_{CDn}(r_d, x)$, and for given x and r_d, the numerical results of $n^{N-MIT}_{O[E]}$, $\varepsilon^{N-MIT}_{1\ O[E]}$ and $R^{N-MIT}_{O[E]}$ are obtained, using Equations (17, 16a), suggesting that, for a given E, they are found to be the same, since $E_{gn1} = E_{gn2} = E_{gn}$.

E in eV	$n^{N-MIT}_{O[E]}$	$\varepsilon^{N-MIT}_{1\ O[E]}$	$R^{N-MIT}_{O[E]}$
At x=0, and in the Te-X(x)-system, in which $E_{gn}(r_{Te}, x=0) = 1.796$ eV,			
1.796	**3.177 [3.177]**	**10.09 [10.09]**	**0.272 [0.272]**
2	3.320 [3.320]	11.02 [11.02]	0.288 [0.288]
2.5	3.849 [3.849]	14.80 [14.80]	0.345 [0.345]
3	4.034 [4.034]	16.27 [16.27]	0.363 [0.363]
3.5	3.502 [3.502]	12.27 [12.27]	0.310 [0.310]
4	3.634 [3.634]	13.20 [13.20]	0.323 [0.323]
4.5	3.947 [3.947]	15.58 [15.58]	0.355 [0.355]
5	2.475 [2.475]	6.127 [6.127]	0.180 [0.180]
5.5	1.403 [1.403]	1.968 [1.968]	0.028 [0.028]

6	1.484 [1.484]	2.203 [2.203]	0.038 [0.038]
…			
10^{22}	**1.992 [1.992]**	**3.968 [3.968]**	**0.110 [0.110]**
At x=0.5, and in the Te -X(x)-system, in which $E_{gn}(r_{Te}, x = 0.5) = 1.796$ eV,			
1.796	**3.128 [3.128]**	**9.785 [9.785]**	**0.266 [0.266]**
2	3.271 [3.271]	10.70 [10.70]	0.283 [0.283]
2.5	3.799 [3.799]	14.43 [14.43]	0.340 [0.340]
3	3.985 [3.985]	15.88 [15.88]	0.358 [0.358]
3.5	3.453 [3.453]	11.92 [11.92]	0.303 [0.303]
4	3.585 [3.585]	12.85 [12.85]	0.318 [0.318]
4.5	3.898 [3.898]	15.19 [15.19]	0.350 [0.350]
5	2.426 [2.426]	5.886 [5.886]	0.173 [0.173]
5.5	1.354 [1.354]	1.832 [1.832]	0.022 [0.022]
6	1.435 [1.435]	2.059 [2.059]	0.032 [0.032]
…			
10^{22}	**1.943 [1.943]**	**3.775 [3.775]**	**0.103 [0.103]**
At x=1, and in the Te -X(x)-system, in which $E_{gn}(r_{Te}, x = 1) = 1.796$ eV,			
1.7960	**3.077 [3.077]**	**9.471 [9.471]**	**0.260 [0.260]**
2	3.220 [3.220]	10.37 [10.37]	0.277 [0.277]
2.5	3.748 [3.748]	14.05 [14.05]	0.335 [0.335]
3	3.934 [3.934]	15.48 [15.48]	0.354 [0.354]
3.5	3.403 [3.403]	11.58 [11.58]	0.298 [0.298]
4	3.534 [3.534]	12.49 [12.49]	0.312 [0.312]
4.5	3.847 [3.847]	14.80 [14.80]	0.345 [0.345]
5	2.376 [2.376]	5.644 [5.644]	0.166 [0.166]
5.5	1.303 [1.303]	1.698 [1.698]	0.017 [0.017]
6	1.384 [1.384]	1.917 [1.917]	0.026 [0.026]
…			
10^{22}	**1.892 [1.892]**	**3.581 [3.581]**	**0.095 [0.095]**

E in eV	$n^{MIT}_{O[E]}$	$\varepsilon^{MIT}_{1\ O[E]}$	$R^{MIT}_{O[E]}$
At x=0, and in the Sb-X(x)-system, in which $E_{gn}(r_{Sb}, x = 0) = 1.7961$ eV,			
1.7961	**3.173 [3.173]**	**10.07 [10.07]**	**0.271 [0.271]**
2	3.316 [3.316]	10.99 [10.99]	0.288 [0.288]
2.5	3.843 [3.843]	14.77 [14.77]	0.345 [0.345]
3	4.030 [4.030]	16.24 [16.24]	0.363 [0.363]
3.5	3.498 [3.498]	12.24 [12.24]	0.308 [0.308]
4	3.630 [3.630]	13.17 [13.17]	0.323 [0.323]
4.5	3.943 [3.943]	15.55 [15.55]	0.354 [0.354]
5	2.471 [2.471]	6.107 [6.107]	0.180 [0.180]
5.5	1.399 [1.399]	1.957 [1.957]	0.028 [0.028]
6	1.480 [1.480]	2.191 [2.191]	0.037 [0.037]
…			
10^{22}	**1.988 [1.988]**	**3.952 [3.952]**	**0.109 [0.109]**

At x=0.5, and in the Sb-X(x)-system, in which $E_{gn}(r_{Sb}, x = 0.5) = 1.7961$ eV,			
1.7961	**3.124 [3.124]**	**9.759 [9.759]**	**0.265 [0.265]**
2	3.267 [3.267]	10.67 [10.67]	0.282 [0.282]
2.5	3.794 [3.794]	14.40 [14.40]	0.340 [0.340]
3	3.980 [3.980]	15.84 [15.84]	0.358 [0.358]
3.5	3.449 [3.449]	11.90 [11.90]	0.303 [0.303]
4	3.580 [3.580]	12.82 [12.82]	0.317 [0.317]
4.5	3.894 [3.894]	15.16 [15.16]	0.350 [0.350]
5	2.422 [2.422]	5.867 [5.867]	0.173 [0.173]
5.5	1.350 [1.350]	1.822 [1.822]	0.022 [0.022]
6	1.431 [1.431]	2.048 [2.048]	0.031 [0.031]
…			
10^{22}	**1.940 [1.940]**	**3.759 [3.759]**	**0.102 [0.102]**
At x=1, and in the Sb-X(x)-system, in which $E_{gn}(r_{Sb}, x = 1) = 1.7961$ eV,			
1.7961	**3.073 [3.073]**	**9.447 [9.447]**	**0.259 [0.259]**
2	3.216 [3.216]	10.34 [10.34]	0.276 [0.276]
2.5	3.744 [3.744]	14.02 [14.02]	0.334 [0.334]
3	3.930 [3.930]	15.45 [15.45]	0.353 [0.353]
3.5	3.399 [3.399]	11.55 [11.55]	0.297 [0.297]
4	3.530 [3.530]	12.46 [12.46]	0.312 [0.312]
4.5	3.843 [3.843]	14.77 [14.77]	0.345 [0.345]
5	2.372 [2.372]	5.625 [5.625]	0.165 [0.165]
5.5	1.299 [1.299]	1.688 [1.688]	0.017 [0.017]
6	1.381 [1.381]	1.096 [1.096]	0.025 [0.025]
…			
10^{22}	**1.888 [1.888]**	**3.566 [3.566]**	**0.095 [0.095]**

Table 4p. **For T=0K and N=$N_{CDp}(r_a, x)$, and for given x and r_d, the numerical results of $n^{N-MIT}_{O[E]}$, $\varepsilon^{N-MIT}_{1\,O[E]}$ and $R^{N-MIT}_{O[E]}$ are obtained, using Equations (17, 16a), suggesting that, for a given E, they are found to be the same, since $E_{gp1} = E_{gp2} = E_{gp}$.**

E in eV	$n^{N-MIT}_{O[E]}$	$\varepsilon^{N-MIT}_{1\,O[E]}$	$R^{N-MIT}_{O[E]}$
At x=0, and in the Mg-X(x)-system, in which $E_{gp}(r_{Mg}, x = 0) = 1.8002$ eV,			
1.8002	**3.120 [3.120]**	**9.735 [9.735]**	**0.265 [0.265]**
2	3.259 [3.259]	10.62 [10.62]	0.281 [0.281]
2.5	3.785 [3.785]	14.33 [14.33]	0.339 [0.339]
3	3.973 [3.973]	15.79 [15.79]	0.357 [0.357]
3.5	3.446 [3.446]	11.88 [11.88]	0.303 [0.303]
4	3.577 [3.577]	12.80 [12.80]	0.317 [0.317]
4.5	3.890 [3.890]	15.13 [15.13]	0.349 [0.349]
5	2.423 [2.423]	5.869 [5.869]	0.173 [0.173]
5.5	1.352 [1.352]	1.829 [1.829]	0.022 [0.022]

6	1.433 [1.433]	2.053 [2.053]	0.032 [0.032]
…			
10^{22}	**1.937 [1.937]**	**3.754 [3.754]**	**0.102 [0.102]**
At x=0.5, and in the Mg-X(x)-system, in which $E_{gp}(r_{Mg}, x = 0.5) = 1.8012$ eV,			
1.8012	**3.071 [3.071]**	**9.434 [9.434]**	**0.259 [0.259]**
2.5	3.735 [3.735]	13.95 [13.95]	0.334 [0.334]
3	3.924 [3.924]	15.40 [15.40]	0.353 [0.353]
3.5	3.398 [3.398]	11.55 [11.55]	0.297 [0.297]
4	3.529 [3.529]	12.46 [12.46]	0.312 [0.312]
4.5	3.842 [3.842]	14.76 [14.76]	0.344 [0.344]
5	2.375 [2.375]	5.642 [5.642]	0.166 [0.166]
5.5	1.305 [1.305]	1.704 [1.704]	0.017 [0.017]
6	1.386 [1.386]	1.921 [1.921]	0.026 [0.026]
…			
10^{22}	**1.890 [1.890]**	**3.571 [3.571]**	**0.095 [0.095]**
At x=1, and in the Mg-X(x)-system, in which $E_{gp}(r_{Mg}, x = 1) = 1.8025$ eV,			
1.8025	**3.022 [3.022]**	**9.131 [9.131]**	**0.253 [0.253]**
2.5	3.684 [3.684]	13.57 [13.57]	0.328 [0.328]
3	3.873 [3.873]	15.00 [15.00]	0.348 [0.348]
3.5	3.348 [3.348]	11.21 [11.21]	0.292 [0.292]
4	3.480 [3.480]	12.11 [12.11]	0.306 [0.306]
4.5	3.792 [3.792]	14.38 [14.38]	0.339 [0.339]
5	2.327 [2.327]	5.414 [5.414]	0.159 [0.159]
5.5	1.257 [1.257]	1.581 [1.581]	0.013 [0.013]
6	1.338 [1.338]	1.790 [1.790]	0.021 [0.021]
…			
10^{22}	**1.840 [1.840]**	**3.388 [3.388]**	**0.087 [0.087]**

E in eV	$n_{O-EP[E-OP]}$	$\varepsilon_{1O-EP[E-OP]}$	$R_{O-EP[E-OP]}$
At x=0, and in the In-X(x)-system, in which $E_{gp}(r_{In}, x = 0) = 1.8031$ eV,			
1.8031	**3.085 [3.085]**	**9.518 [9.518]**	**0.260 [0.260]**
2.5	3.747 [3.747]	14.04 [14.04]	0.335 [0.335]
3	3.936 [3.936]	15.49 [15.49]	0.354 [0.354]
3.5	3.412 [3.412]	11.64 [11.64]	0.299 [0.299]
4	3.543 [3.543]	12.55 [12.55]	0.313 [0.313]
4.5	3.855 [3.855]	14.86 [14.86]	0.346 [0.346]
5	2.391 [2.391]	5.716 [5.716]	0.168 [0.168]
5.5	1.322 [1.322]	1.747 [1.747]	0.019 [0.019]
6	1.402 [1.402]	1.966 [1.966]	0.028 [0.028]
…			
10^{22}	**1.904 [1.904]**	**3.626 [3.626]**	**0.097 [0.097]**

At x=0.5, and in the In-X(x)-system, in which $E_{gp}(r_{In}, x = 0.5) = 1.8048$ eV,

1.8048	**3.037 [3.037]**	**9.223 [9.223]**	**0.255 [0.255]**
2.5	3.696 [3.696]	13.66 [13.66]	0.330 [0.330]
3	3.887 [3.887]	15.11 [15.11]	0.349 [0.349]
3.5	3.364 [3.364]	11.32 [11.32]	0.293 [0.293]
4	3.495 [3.495]	12.22 [12.22]	0.308 [0.308]
4.5	3.807 [3.807]	14.50 [14.50]	0.341 [0.341]
5	2.344 [2.344]	5.497 [5.497]	0.162 [0.162]
5.5	1.276 [1.276]	1.629 [1.629]	0.015 [0.015]
6	1.356 [1.356]	1.840 [1.840]	0.023 [0.023]
…			
$\mathbf{10^{22}}$	**1.857 [1.857]**	**3.449 [3.449]**	**0.090 [0.090]**

At x=1, and in the In-X(x)-system, in which $E_{gno}(r_{In}, x = 1) = 0.8069$ eV,

1.8069	**2.987 [2.987]**	**8.925 [8.925]**	**0.248 [0.248]**
2.5	3.644 [3.644]	13.28 [13.28]	0.324 [0.324]
3	3.835 [3.835]	14.71 [14.71]	0.344 [0.344]
3.5	3.315 [3.315]	10.99 [10.99]	0.288 [0.288]
4	3.446 [3.446]	11.88 [11.88]	0.303 [0.303]
4.5	3.758 [3.758]	14.12 [14.12]	0.336 [0.336]
5	2.297 [2.297]	5.277 [5.277]	0.155 [0.155]
5.5	1.230 [1.230]	1.513 [1.513]	0.011 [0.011]
6	1.310 [1.310]	1.716 [1.716]	0.018 [0.018]
…			
$\mathbf{10^{22}}$	**1.809 [1.809]**	**3.273 [3.273]**	**0.083 [0.083]**

Table 5. **For T=0K, $E \cong E_{gn(gp)}$ and $N^* = N_{CDn(NDp)}$, and from Eq. (16b), the numerical results of $\sigma_{O[E]}^{EBT}$, $\kappa_{O[E]}^{EBT}$, $\varepsilon_{2O[2E]}^{EBT}$ and $\propto_{O[E]}^{EBT}$ are obtained, using Equations (18, 19b, 19c, 19d), suggesting that they increase (↗) with increasing (↗) $r_{d(a)}$.**

Donor		P	Te	Sb	Sn
r_d (nm) [4]	↗	0.110	**0.132**	0.136	0.140
At **x=0**,					
$\sigma_{O[E]}^{EBT}\left(\frac{10^2}{\Omega\times cm}\right)$	↗	4.079 [2.364]	4.645 [2.695]	4.664 [2.707]	4.722 [2.741]
$\kappa_{O[E]}^{EBT} \times 10^3$	↗	2.283 [1.317]	3.087 [1.780]	3.116 [1.797]	3.207 [1.849]
$\varepsilon_{2O[2E]}^{EBT} \times 10^2$	↗	1.517 [0.879]	1.964 [1.140]	1.980 [1.149]	2.030 [1.178]
$\propto_{O[E]}^{EBT}\left(\frac{10^2}{cm}\right)$	↗	4.146 [2.392]	5.619 [3.239]	5.673 [3.271]	5.840 [3.367]

At **x=0.5**,

$\sigma_{O[E]}^{EBT}\left(\frac{10^2}{\Omega\times cm}\right)\nearrow$	3.048 [2.016]	3.473 [2.298]	3.488 [2.308]	3.532 [2.337]
$\kappa_{O[E]}^{EBT}\times 10^3\nearrow$	1.820 [1.199]	2.461 [1.620]	2.484 [1.635]	2.557 [1.683]
$\varepsilon_{2O[2E]}^{EBT}\times 10^2\nearrow$	1.191 [0.788]	1.544 [1.022]	1.557 [1.030]	1.596 [1.056]
$\propto_{O[E]}^{EBT}\left(\frac{10^2}{cm}\right)\nearrow$	3.306 [2.178]	4.479 [2.948]	4.523 [2.977]	4.656 [3.064]
At **x=1**,				
$\sigma_{O[E]}^{EBT}\left(\frac{10^2}{\Omega\times cm}\right)\nearrow$	2.196 [1.630]	2.503 [1.858]	2.514 [1.866]	2.546 [1.889]
$\kappa_{O[E]}^{EBT}\times 10^3\nearrow$	1.403 [1.038]	1.897 [1.402]	1.915 [1.415]	1.971 [1.457]
$\varepsilon_{2O[2E]}^{EBT}\times 10^2\nearrow$	0.904 [0.671]	1.173 [0.871]	1.183 [0.878]	1.213 [0.900]
$\propto_{O[E]}^{EBT}\left(\frac{10^2}{cm}\right)\nearrow$	2.550 [1.887]	3.453 [2.552]	3.487 [2.577]	3.589 [2.652]

Acceptor	**Ga**	Mg	In	Cd
r_a (nm) $\nearrow$	**0.126**	0.140	0.144	0.148
At **x=0**,				
$\sigma_{O[E]}^{EBT}\left(\frac{10^3}{\Omega\times cm}\right)\nearrow$	2.229 [0.516]	2.357 [0.545]	2.440 [0.565]	2.542 [0.588]
$\kappa_{O[E]}^{EBT}\times 10^2\nearrow$	1.512 [0.341]	1.722 [0.388]	1.866 [0.419]	2.054 [0.460]
$\varepsilon_{2O[2E]}^{EBT}\times 10^1\nearrow$	0.943 [0.218]	1.051 [0.243]	1.124 [0.260]	1.218 [0.282]
$\propto_{O[E]}^{EBT}\left(\frac{10^3}{cm}\right)\nearrow$	2.933 [0.663]	3.167 [0.713]	3.323 [0.746]	3.518 [0.788]
At **x=0.5**,				
$\sigma_{O[E]}^{EBT}\left(\frac{10^3}{\Omega\times cm}\right)\nearrow$	3.740 [0.610]	3.953 [0.645]	4.092 [0.668]	4.263 [0.696]
$\kappa_{O[E]}^{EBT}\times 10^2\nearrow$	2.746 [0.431]	3.132 [0.489]	3.399 [0.529]	3.746 [0.580]
$\varepsilon_{2O[2E]}^{EBT}\times 10^1\nearrow$	1.663 [0.271]	1.852 [0.302]	1.981 [0.323]	2.145 [0.350]
$\propto_{O[E]}^{EBT}\left(\frac{10^3}{cm}\right)\nearrow$	5.329 [0.837]	5.765 [0.901]	6.057 [0.943]	6.424 [0.995]
At **x=1**,				
$\sigma_{O[E]}^{EBT}\left(\frac{10^3}{\Omega\times cm}\right)\nearrow$	6.897 [0.715]	7.291 [0.755]	7.548 [0.782]	7.863 [0.815]
$\kappa_{O[E]}^{EBT}\times 10^2\nearrow$	5.566 [0.541]	6.365 [0.614]	6.921 [0.663]	7.648 [0.727]
$\varepsilon_{2O[2E]}^{EBT}\times 10^1\nearrow$	3.232 [0.335]	3.599 [0.373]	3.847 [0.399]	4.162 [0.431]
$\propto_{O[E]}^{EBT}\left(\frac{10^3}{cm}\right)\nearrow$	10.80 [1.050]	11.72 [1.130]	12.35 [1.183]	13.14 [1.249]

Table 6n. **In the X(x)-system, at E=3.2 eV and T=20 K, for given r_d and x, the numerical results of $n_{O[E]}(E)$, $\kappa_{O[E]}(E)$, $\varepsilon_{1O[E]}(E)$ and $\varepsilon_{2O[E]}(E)$, are obtained, as functions of N, by using Equations (17, 19b, 19c and 16), respectively, noting that, with increasing N, $\eta_{o[E]}$ increases [increases], and $E_{gn1\ O[E]}$ increases [decreases], respectively.**

N (10^{18} cm^{-3}) ↗	15	26	60	100
At x=0				
For $\mathbf{r_d} = \mathbf{r_{Te}}$,				
$\xi_{nO[E]} \gg 1$	91.55 [60.12]	133.4 [87.63]	234.8 [154.2]	330.8 [217.3]
$E_{gn1\ O[E]}$ in eV	1.87 [1.71]	1.92 [1.69]	2.04 [1.64]	2.17 [1.60]
$n_{O[E]}$	3.80 [3.95]	3.75 [3.98]	3.62 [4.02]	3.49 [4.06]
$\kappa_{O[E]}$	0.02 [0.01]	0.04 [0.02]	0.08 [0.04]	0.14 [0.06]
$\varepsilon_{1O[E]}$	14.43 [15.63]	14.06 [15.81]	13.11 [16.18]	12.17 [16.50]
$\varepsilon_{2O[E]}$	0.17 [0.09]	0.29 [0.14]	0.62 [0.30]	0.98 [0.48]
For $\mathbf{r_d} = \mathbf{r_{Sb}}$,				
$\xi_{nO[E]} \gg 1$	91.52 [60.10]	133.4 [87.62]	234.8 [154.2]	330.8 [217.3]
$E_{gn1\ O[E]}$ in eV	1.87 [1.71]	1.92 [1.69]	2.04 [1.64]	2.17 [1.60]
$n_{O[E]}$	3.79 [3.95]	3.74 [3.97]	3.62 [4.02]	3.49 [4.06]
$\kappa_{O[E]}$	0.02 [0.01]	0.04 [0.02]	0.08 [0.04]	0.14 [0.06]
$\varepsilon_{1O[E]}$	14.40 [15.59]	14.03 [15.77]	13.08 [16.15]	12.14 [16.50]
$\varepsilon_{2O[E]}$	0.17 [0.09]	0.29 [0.14]	0.61 [0.30]	0.98 [0.48]
At x=0.5				
For $\mathbf{r_d} = \mathbf{r_{Te}}$,				
$\xi_{nO[E]} \gg 1$	103.0 [74.80]	149.5 [108.6]	262.3 [190.6]	369.3 [268.3]
$E_{gn1\ O[E]}$ in eV	1.87 [1.69]	1.92 [1.66]	2.05 [1.60]	2.18 [1.55]
$n_{O[E]}$	3.75 [3.93]	3.70 [3.95]	3.57 [4.01]	3.43 [4.06]
$\kappa_{O[E]}$	0.03 [0.01]	0.04 [0.02]	0.10 [0.05]	0.17 [0.08]
$\varepsilon_{1O[E]}$	14.08 [15.41]	13.70 [15.64]	12.71 [16.10]	11.72 [16.49]
$\varepsilon_{2O[E]}$	0.20 [0.12]	0.33 [0.20]	0.72 [0.42]	1.15 [0.66]
For $\mathbf{r_d} = \mathbf{r_{Sb}}$,				
$\xi_{nO[E]} \gg 1$	103.0 [74.79]	149.5 [108.6]	262.3 [190.5]	369.3 [268.2]
$E_{gn1\ O[E]}$ in eV	1.87 [1.69]	1.92 [1.66]	2.05 [1.60]	2.18 [1.55]
$n_{O[E]}$	3.75 [3.92]	3.70 [3.95]	3.56 [4.01]	3.42 [4.06]
$\kappa_{O[E]}$	0.03 [0.01]	0.04 [0.02]	0.10 [0.05]	0.17 [0.08]
$\varepsilon_{1O[E]}$	14.05 [15.38]	13.67 [15.60]	12.67 [16.07]	11.68 [16.45]

$\varepsilon_{2O[E]}$	0.20 [0.12]	0.33 [0.20]	0.72 [0.41]	1.15 [0.66]
At x=1				
For $\mathbf{r_d = r_{Te}}$,				
$\xi_{nO[E]} \gg 1$	123.7 [98.21]	179.1 [142.2]	313.6 [248.9]	441.1 [350.1]
$E_{gn1\ O[E]}$ in eV	1.86 [1.65]	1.92 [1.61]	2.07 [1.53]	2.22 [1.46]
$n_{O[E]}$	3.70 [3.91]	3.65 [3.95]	3.50 [4.03]	3.34 [4.09]
$\kappa_{O[E]}$	0.04 [0.02]	0.06 [0.04]	0.13 [0.08]	0.23 [0.12]
$\varepsilon_{1O[E]}$	13.72 [15.31]	13.31 [15.61]	12.22 [16.24]	11.11 [16.76]
$\varepsilon_{2O[E]}$	0.27 [0.18]	0.44 [0.30]	0.94 [0.63]	1.51 [1.01]
For $\mathbf{r_d = r_{Sb}}$,				
$\xi_{nO[E]} \gg 1$	123.7 [98.20]	179.1 [142.2]	313.6 [248.9]	441.1 [350.1]
$E_{gn1\ O[E]}$ in eV	1.86 [1.65]	1.92 [1.61]	2.07 [1.53]	2.21 [1.46]
$n_{O[E]}$	3.70 [3.91]	3.64 [3.95]	3.49 [4.02]	3.33 [4.09]
$\kappa_{O[E]}$	0.03 [0.02]	0.06 [0.04]	0.13 [0.08]	0.22 [0.12]
$\varepsilon_{1O[E]}$	13.69 [15.27]	13.28 [15.57]	12.18 [16.20]	11.07 [16.72]
$\varepsilon_{2O[E]}$	0.27 [0.18]	0.44 [0.30]	0.94 [0.63]	1.51 [1.00]
$N\ (10^{18}\ cm^{-3})$ ↗	15	26	60	100

Table 6p. In the X(x)-system, at E=3.2 eV and T=20 K, for given $\mathbf{r_a}$ and x, the numerical results of $\mathbf{n_{O[E]}(E)}$, $\boldsymbol{\kappa_{O[E]}(E)}$, $\boldsymbol{\varepsilon_{1O[E]}(E)}$ and $\boldsymbol{\varepsilon_{2O[E]}(E)}$, are obtained, as functions of N, by using Equations (17, 19b, 19c and 16), respectively, noting that, with increasing N, $\boldsymbol{\eta_{o[E]}}$ increases [increases], and $\mathbf{E_{gp1\ O[E]}}$ increases [decreases], respectively.

$N\ (10^{18}\ cm^{-3})$ ↗	15	26	60	100
At x=0				
For $\mathbf{r_a = r_{Mg}}$,				
$\xi_{pO[E]} \gg 1$	74.97 [25.70]	120.0 [41.16]	224.8 [77.14]	322.4 [110.6]
$E_{gp1\ O[E]}$ in eV	1.91 [1.78]	1.98 [1.77]	2.14 [1.76]	2.30 [1.74]
$n_{O[E]}$	3.71 [3.83]	3.64 [3.84]	3.46 [3.86]	3.30 [3.87]
$\kappa_{O[E]}$	0.02 [0.004]	0.03 [0.006]	0.08 [0.013]	0.14 [0.02]
$\varepsilon_{1O[E]}$	13.74 [14.7]	13.22 [14.8]	11.99 [14.9]	10.85 [15.0]
$\varepsilon_{2O[E]}$	0.13 [0.03]	0.24 [0.05]	0.56 [0.10]	0.91 [0.16]
For $\mathbf{r_a = r_{In}}$,				
$\xi_{pO[E]} \gg 1$	72.79 [24.95]	118.3 [40.57]	223.6 [76.71]	321.4 [110.3]

$E_{gp1\ O[E]}$ in eV	1.91 [1.78]	1.98 [1.77]	2.14 [1.76]	2.30 [1.75]
$n_{O[E]}$	3.67 [3.80]	3.60 [3.80]	3.43 [3.82]	3.26 [3.83]
$\kappa_{O[E]}$	0.02 [0.003]	0.03 [0.006]	0.08 [0.013]	0.13 [0.02]
$\varepsilon_{1O[E]}$	13.50 [14.4]	12.97 [14.5]	11.75 [14.6]	10.62 [14.7]
$\varepsilon_{2O[E]}$	0.12 [0.03]	0.23 [0.05]	0.54 [0.10]	0.88 [0.15]
At x=0.5				
For $\mathbf{r_a} = \mathbf{r_{Mg}}$,				
$\xi_{pO[E]} \gg 1$	68.39 [18.67]	122.1 [33.38]	242.2 [66.26]	352.5 [96.43]
$E_{gp1\ O[E]}$ in eV	1.90 [1.78]	1.99 [1.78]	2.18 [1.77]	2.36 [1.76]
$n_{O[E]}$	3.66 [3.78]	3.58 [3.79]	3.37 [3.80]	3.18 [3.81]
$\kappa_{O[E]}$	0.01 [0.002]	0.03 [0.005]	0.09 [0.011]	0.16 [0.02]
$\varepsilon_{1O[E]}$	13.42 [14.2]	12.79 [14.3]	11.37 [14.4]	10.08 [14.5]
$\varepsilon_{2O[E]}$	0.11 [0.02]	0.25 [0.04]	0.62 [0.08]	1.03 [0.13]
For $\mathbf{r_a} = \mathbf{r_{In}}$,				
$\xi_{pO[E]} \gg 1$	63.90 [17.44]	118.7 [32.47]	239.9 [65.61]	350.5 [95.90]
$E_{gp1\ O[E]}$ in eV	1.90 [1.79]	1.99 [1.78]	2.18 [1.77]	2.36 [1.76]
$n_{O[E]}$	3.63 [3.74]	3.55 [3.75]	3.34 [3.76]	3.14 [3.77]
$\kappa_{O[E]}$	0.01 [0.002]	0.03 [0.005]	0.09 [0.010]	0.16 [0.02]
$\varepsilon_{1O[E]}$	13.21 [14.0]	12.57 [14.1]	11.15 [14.2]	09.87 [14.3]
$\varepsilon_{2O[E]}$	0.10 [0.02]	0.23 [0.03]	0.60 [0.08]	1.00 [0.12]
At x=1				
For $\mathbf{r_a} = \mathbf{r_{Mg}}$,				
$\xi_{pO[E]} \gg 1$	48.74 [9.978]	122.8 [25.31]	273.3 [56.37]	407.6 [84.10]
$E_{gp1\ O[E]}$ in eV	1.87 [1.79]	1.99 [1.78]	2.24 [1.77]	2.47 [1.76]
$n_{O[E]}$	3.64 [3.72]	3.52 [3.73]	3.26 [3.74]	3.01 [3.75]
$\kappa_{O[E]}$	0.01 [0.001]	0.04 [0.003]	0.11 [0.009]	0.21 [0.01]
$\varepsilon_{1O[E]}$	13.27 [13.9]	12.39 [13.92]	10.61 [14.0]	09.04 [14.1]
$\varepsilon_{2O[E]}$	0.08 [0.009]	0.25 [0.02]	0.75 [0.06]	1.30 [0.10]
For $\mathbf{r_a} = \mathbf{r_{In}}$,				
$\xi_{pO[E]} \gg 1$	37.05 [7.538]	115.8 [23.87]	268.6 [55.41]	403.8 [83.32]
$E_{gp1\ O[E]}$ in eV	1.86 [1.80]	1.99 [1.79]	2.24 [1.78]	2.47 [1.77]
$n_{O[E]}$	3.62 [3.69]	3.49 [3.70]	3.23 [3.70]	2.98 [3.71]
$\kappa_{O[E]}$	0.007 [0.001]	0.03 [0.003]	0.11 [0.008]	0.21 [0.01]
$\varepsilon_{1O[E]}$	13.14 [13.5]	12.21 [13.6]	10.42 [13.7]	08.86 [13.8]
$\varepsilon_{2O[E]}$	0.05 [0.006]	0.23 [0.02]	0.71 [0.06]	1.25 [0.10]

$N\ (10^{18}\ cm^{-3})$ ↗	15	26	60	100

Table 7n. In the X(x)-system, at E=3.2 eV and $N = 10^{20} cm^{-3}$, for given r_d and x, the numerical results of $n_{O[E]}(E)$, $\kappa_{O[E]}(E)$, $\varepsilon_{1O[E]}(E)$ and $\varepsilon_{2O[E]}(E)$, are obtained, as functions of T, by using Equations (17, 19b, 19c and 16), respectively, noting that $\eta_{o[E]}$ and $E_{gn1\ O[E]}$ both decrease with increasing T, respectively.

T ↗	20 K	50 K	100 K	300 K
At x=0				
For $r_d = r_{Te}$,				
$\xi_{nO[E]} \gg 1$	330.8 [217.2]	132.3 [86.90]	66.15 [43.44]	22.02 [14.43]
$E_{gn1\ O[E]}$ in eV	2.17 [1.600]	2.16 [1.596]	2.15 [1.584]	2.07 [1.505]
$n_{O[E]}$	3.49 [4.062]	3.50 [4.066]	3.51 [4.078]	3.59 [4.153]
$\kappa_{O[E]}$	0.1411 [0.0589]	0.1409 [0.0588]	0.1404 [0.0587]	0.1378 [0.0584]
$\varepsilon_{1O[E]}$	12.17 [16.50]	12.20 [16.53]	12.30 [16.63]	12.89 [17.24]
$\varepsilon_{2O[E]}$	0.9852 [0.4787]	0.9853 [0.4788]	0.9857 [0.4794]	0.990 [0.485]
For $r_d = r_{Sb}$,				
$\xi_{nO[E]} \gg 1$	330.8 [217.3]	132.3 [86.90]	66.15 [43.43]	22.01 [14.43]
$E_{gn1\ O[E]}$ in eV	2.17 [1.601]	2.16 [1.597]	2.15 [1.584]	2.07 [1.506]
$n_{O[E]}$	3.48 [4.057]	3.49 [4.061]	3.50 [4.073]	3.59 [4.148]
$\kappa_{O[E]}$	0.1408 [0.0588]	0.1406 [0.0587]	0.1402 [0.0586]	0.1375 [0.0583]
$\varepsilon_{1O[E]}$	12.14 [16.46]	12.17 [16.49]	12.26 [16.59]	12.86 [17.20]
$\varepsilon_{2O[E]}$	0.9821 [0.4773]	0.9822 [0.4774]	0.9826 [0.4780]	0.987 [0.484]
At x=0.5				
For $r_d = r_{Te}$,				
$\xi_{nO[E]} \gg 1$	369.3 [268.3]	147.7 [107.3]	73.85 [53.64]	24.59 [17.84]
$E_{gn1\ O[E]}$ in eV	2.182 [1.549]	2.17 [1.541]	2.16 [1.523]	2.05 [1.420]
$n_{O[E]}$	3.42 [4.062]	3.43 [4.069]	3.45 [4.086]	3.56 [4.182]
$\kappa_{O[E]}$	0.1678 [0.0813]	0.1675 [0.0812]	0.1666 [0.0809]	0.1620 [0.0797]
$\varepsilon_{1O[E]}$	11.72 [16.49]	11.77 [16.55]	11.91 [16.69]	12.67 [17.49]
$\varepsilon_{2O[E]}$	1.1503 [0.6606]	1.1504 [0.6608]	1.1508 [0.6612]	1.1548 [0.6665]
For $r_d = r_{Sb}$,				
$\xi_{nO[E]} \gg 1$	369.3 [268.2]	147.7 [107.3]	73.85 [53.64]	24.59 [17.84]
$E_{gn1\ O[E]}$ in eV	2.18 [1.549]	2.17 [1.542]	2.16 [1.524]	2.05 [1.421]
$n_{O[E]}$	3.42 [4.057]	3.43 [4.064]	3.45 [4.081]	3.56 [4.177]
$\kappa_{O[E]}$	0.1675 [0.0812]	0.1672 [0.0810]	0.1663 [0.0808]	0.1617 [0.0795]

$\varepsilon_{1O[E]}$	11.68 [16.45]	11.73 [16.51]	11.87 [16.65]	12.63 [17.44]
$\varepsilon_{2O[E]}$	1.1466 [0.6586]	1.1467 [0.6588]	1.1471 [0.6593]	1.1511 [0.6645]
At x=1				
For $\mathbf{r_d} = \mathbf{r_{Te}}$,				
$\xi_{nO[E]} \gg 1$	441.2 [350.1]	176.5 [140.0]	88.22 [70.01]	29.38 [23.31]
$E_{gn1\ O[E]}$ in eV	2.22 [1.459]	2.21 [1.449]	2.18 [1.424]	2.05 [1.297]
$n_{O[E]}$	3.34 [4.096]	3.35 [4.105]	3.38 [4.128]	3.51 [4.244]
$\kappa_{O[E]}$	0.2263 [0.123]	0.2256 [0.122]	0.2239 [0.121]	0.2158 [0.119]
$\varepsilon_{1O[E]}$	11.11 [16.76]	11.18 [16.83]	11.36 [17.02]	12.28 [17.99]
$\varepsilon_{2O[E]}$	1.51169 [1.0068]	1.51178 [1.0069]	1.51209 [1.0074]	1.51558 [1.0119]
For $\mathbf{r_d} = \mathbf{r_{Sb}}$,				
$\xi_{nO[E]} \gg 1$	441.1 [350.1]	176.4 [140.0]	88.22 [70.01]	29.38 [23.31]
$E_{gn1\ O[E]}$ in eV	2.22 [1.460]	2.21 [1.450]	2.18 [1.425]	2.05 [1.298]
$n_{O[E]}$	3.33 [4.091]	3.34 [4.100]	3.37 [4.123]	3.51 [4.239]
$\kappa_{O[E]}$	0.2259 [0.123]	0.2252 [0.122]	0.2235 [0.121]	0.2154 [0.119]
$\varepsilon_{1O[E]}$	11.07 [16.72]	11.14 [16.79]	11.32 [16.98]	12.25 [17.95]
$\varepsilon_{2O[E]}$	1.50678 [1.0037]	1.50687 [1.0038]	1.50718 [1.0042]	1.51066 [1.0088]
T ↗	20 K	50 K	100 K	300 K

Table 7p. **In the X(x)-system, at E=3.2 eV and $\mathbf{N = 10^{20} cm^{-3}}$, for given $\mathbf{r_a}$ and x, the numerical results of $\mathbf{n_{O[E]}(E)}$, $\boldsymbol{\kappa_{O[E]}(E)}$, $\boldsymbol{\varepsilon_{1O[E]}(E)}$ and $\boldsymbol{\varepsilon_{2O[E]}(E)}$, are obtained, as functions of T, by using Equations (17, 19b, 19c and 16), respectively, noting that $\boldsymbol{\eta_{o[E]}}$ and $\mathbf{E_{gp1\ O[E]}}$ both decrease with increasing T, respectively.**

T ↗	20 K	50 K	100 K	300 K
At x=0				
For $\mathbf{r_a} = \mathbf{r_{Mg}}$,				
$\xi_{pO[E]} \gg 1$	322.4 [110.6]	129.0 [44.24]	64.47 [22.09]	21.46 [7.262]
$E_{gp1\ O[E]}$ in eV	2.30 [1.745]	2.29 [1.741]	2.28 [1.728]	2.20 [1.650]
$n_{O[E]}$	3.29 [3.868]	3.30 [3.872]	3.31 [3.884]	3.40 [3.960]
$\kappa_{O[E]}$	0.1384 [0.02054]	0.1382 [0.02055]	0.1377 [0.02058]	0.135 [0.021]
$\varepsilon_{1O[E]}$	10.85 [14.96]	10.88 [14.99]	10.97 [15.09]	11.54 [15.68]
$\varepsilon_{2O[E]}$	0.9128 [0.158]	0.9129 [0.1591]	0.9133 [0.1599]	0.9175 [0.168]
For $\mathbf{r_a} = \mathbf{r_{In}}$,				
$\xi_{pO[E]} \gg 1$	321.4 [110.3]	128.5 [44.10]	64.27 [22.02]	21.39 [7.237]

$E_{gp1\ O[E]}$ in eV	2.30 [1.749]	2.29 [1.745]	2.28 [1.733]	2.20 [1.654]
$n_{O[E]}$	3.26 [3.830]	3.27 [3.835]	3.28 [3.847]	3.36 [3.923]
$\kappa_{O[E]}$	0.1358 [0.02031]	0.1356 [0.02032]	0.1351 [0.02035]	0.132 [0.0209]
$\varepsilon_{1O[E]}$	10.618 [14.67]	10.65 [14.70]	10.74 [14.80]	11.31 [15.39]
$\varepsilon_{2O[E]}$	0.8859 [0.1556]	0.8860 [0.1558]	0.8864 [0.1565]	0.890 [0.164]
At x=0.5				
For $\mathbf{r_a} = \mathbf{r_{Mg}}$,				
$\xi_{pO[E]} \gg 1$	352.5 [96.43]	141.0 [38.55]	70.48 [19.24]	23.46 [6.296]
$E_{gp1\ O[E]}$ in eV	2.36 [1.76]	2.35 [1.75]	2.34 [1.73]	2.23 [1.63]
$n_{O[E]}$	3.18 [3.808]	3.19 [3.815]	3.21 [3.833]	3.32 [3.933]
$\kappa_{O[E]}$	0.1627 [0.016724]	0.1623 [0.016726]	0.1614 [0.016747]	0.156 [0.0174]
$\varepsilon_{1O[E]}$	10.08 [14.50]	10.13 [14.55]	10.26 [14.69]	10.99 [15.47]
$\varepsilon_{2O[E]}$	1.0349 [0.1274]	1.0350 [0.1276]	1.0353 [0.1284]	1.039 [0.137]
For $\mathbf{r_a} = \mathbf{r_{In}}$,				
$\xi_{pO[E]} \gg 1$	350.5 [95.90]	140.2 [38.34]	70.10 [19.14]	23.33 [6.259]
$E_{gp1\ O[E]}$ in eV	2.36 [1.76]	2.35 [1.75]	2.34 [1.74]	2.23 [1.63]
$n_{O[E]}$	3.14 [3.771]	3.15 [3.778]	3.17 [3.796]	3.29 [3.896]
$\kappa_{O[E]}$	0.1591 [0.01652]	0.1587 [0.01652]	0.1577 [0.01654]	0.153 [0.0172]
$\varepsilon_{1O[E]}$	9.871 [14.22]	9.921 [14.27]	10.05 [14.41]	10.77 [15.18]
$\varepsilon_{2O[E]}$	1.0009 [0.1246]	1.0010 [0.1248]	1.0014 [0.1256]	1.005 [0.134]
At x=1				
For $\mathbf{r_a} = \mathbf{r_{Mg}}$,				
$\xi_{pO[E]} \gg 1$	407.6 [84.10]	163.0 [33.62]	81.52 [16.77]	27.14 [5.452]
$E_{gp1\ O[E]}$ in eV	2.46 [1.77]	2.45 [1.76]	2.43 [1.73]	2.30 [1.60]
$n_{O[E]}$	3.01 [3.749]	3.02 [3.759]	3.05 [3.783]	3.19 [3.906]
$\kappa_{O[E]}$	0.2154 [0.01377]	0.2146 [0.01377]	0.2127 [0.01379]	0.204 [0.0145]
$\varepsilon_{1O[E]}$	09.04 [14.06]	9.11 [14.13]	9.27 [14.31]	10.16 [15.26]
$\varepsilon_{2O[E]}$	1.2988 [0.1033]	1.2989 [0.1035]	1.2992 [0.1044]	1.302 [0.1135]
For $\mathbf{r_a} = \mathbf{r_{In}}$,				
$\xi_{pO[E]} \gg 1$	403.8 [83.32]	161.5 [33.31]	80.76 [16.62]	26.89 [5.398]
$E_{gp1\ O[E]}$ in eV	2.46 [1.77]	2.45 [1.76]	2.43 [1.74]	2.30 [1.61]
$n_{O[E]}$	2.98 [3.712]	2.99 [3.722]	3.02 [3.746]	3.16 [3.870]
$\kappa_{O[E]}$	0.2092 [0.01357]	0.2085 [0.01357]	0.2066 [0.01359]	0.198 [0.0143]
$\varepsilon_{1O[E]}$	8.86 [13.78]	8.93 [13.85]	9.10 [14.04]	9.97 [14.97]
$\varepsilon_{2O[E]}$	1.2490 [0.1008]	1.2491 [0.1010]	1.2493 [0.1018]	1.252 [0.1109]

T	↗	20 K	50 K	100 K	300 K

Table 8n. **For T=20K and $N = 10^{20} cm^{-3}$, and for given x and r_d, the numerical results of $\sigma_{O[E]}$ (E), $\varepsilon_{2O[2E]}(E)$ and $\propto_{O[E]}$ (E) , are obtained by using Equations (18, 19c, 19d), noting that, as given in Eq. (15), $E_{gnE} \equiv E_{gn2} + E_{Fn}$ and $E_{gnO} \equiv E_{gn1} + E_{Fn}$.**

E in eV	$\sigma_{O[E]}$ $\left(\frac{10^5}{\Omega\times cm}\right)$	$\varepsilon_{2O[2E]}$	$\propto_{O[E]}$ $\left(\frac{10^5}{cm}\right)$
At x=0 and $r_d = r_{Te}$,			
1.6004 $= E_{gn2}$	**… [0]**	**… [0]**	**… [0]**
1.9749 $= E_{gnE}$	**… [0.2016]**	**… [0.775]**	**… [0.216]**
2.1674 $= E_{gn1}$	**0** [0.2016]	**0** [0.706]	**0** [0.205]
2.7376 $=E_{gnO}$	**0.4150** [0.2017]	**1.1514** [0.559]	**0.4628** [0.168]
3.5	0.4151 [0.2017]	0.901 [0.438]	0.480 [0.217]
4	0.4151 [0.2017]	0.788 [0.383]	0.463 [0.209]
4.5	0.4152 [0.2017]	0.701 [0.340]	0.430 [0.191]
5	0.4152 [0.2017]	0.631 [0.306]	0.614 [0.326]
5.5	0.4152 [0.2017]	0.573 [0.278]	0.930 [0.645]
6	0.4152 [0.2017]	0.525 [0.255]	0.913 [0.589]
…			
10^{22}	**0.4152 [0.2017]**	**0 [0]**	**0.8022 [0.3897]**
At x=0 and $r_d = r_{Sb}$,			
1.6012 $= E_{gn2}$	**… [0]**	**… [0]**	**… [0]**
1.9757 $= E_{gnE}$	**… [0.2002]**	**… [0.7727]**	**… [0.2161]**
2.1681 $= E_{gn1}$	**0** [0.2002]	**0** [0.7043]	**0** [0.2047]
2.7383 $=E_{gnO}$	**0.4120** [0.2002]	**1.1474** [0.5577]	**0.4620** [0.1683]
4	0.4121 [0.2002]	0.786 [0.3818]	0.462 [0.2090]
4.5	0.4121 [0.2003]	0.698 [0.3394]	0.429 [0.1911]
5	0.4121 [0.2003]	0.629 [0.3055]	0.613 [0.3258]
5.5	0.4121 [0.2003]	0.571 [0.2777]	0.929 [0.6446]
6	0.4121 [0.2003]	0.524 [0.2546]	0.912 [0.5891]
…			
10^{22}	**0.4122 [0.2003]**	**0 [0]**	**0.8013 [0.3894]**
At x=0.5 and $r_d = r_{Te}$,			
1.5486 $= E_{gn2}$	**… [0]**	**… [0]**	**… [0]**
2.0110 $= E_{gnE}$	**… [0.2647]**	**… [1.0508]**	**… [0.2933]**
2.1820 $= E_{gn1}$	**0** [0.2647]	**0** [0.9686]	**0** [0.2787]
2.8186 $=E_{gnO}$	**0.4609** [0.2647]	**1.3057** [0.7500]	**0.5381** [0.2234]
4	0.4611 [0.2648]	0.920 [0.5285]	0.549 [0.2915]
4.5	0.4611 [0.2648]	0.818 [0.4698]	0.510 [0.2658]
5	0.4611 [0.2648]	0.736 [0.4228]	0.730 [0.4654]
5.5	0.4611 [0.2648]	0.669 [0.3844]	1.111 [0.9754]
6	0.4611 [0.2648]	0.614 [0.3524]	1.091 [0.8776]

…			
10^{22}	**0.46115 [0.2648]**	**0 [0]**	**0.96039 [0.5515]**
At x=0.5 and $r_d = r_{Sb}$,			
$\mathbf{1.5496} = E_{gn2}$	**… [0]**	**… [0]**	**… [0]**
$\mathbf{2.0119} = E_{gnE}$	**… [0.2628]**	**… [1.0472]**	**… [0.2927]**
$\mathbf{2.1829} = E_{gn1}$	**0** [0.2628]	**0** [0.9653]	**0** [0.2782]
$\mathbf{2.8195} = E_{gnO}$	**0.4575** [0.2628]	**1.3011** [0.7475]	**0.5371** [0.2229]
4	0.4577 [0.2629]	0.917 [0.5269]	0.548 [0.2909]
4.5	0.4577 [0.2629]	0.815 [0.4684]	0.509 [0.2653]
5	0.4577 [0.2629]	0.734 [0.4216]	0.729 [0.4648]
5.5	0.4577 [0.2629]	0.667 [0.3832]	1.110 [0.9751]
6	0.4577 [0.2629]	0.612 [0.3513]	1.090 [0.8773]
…			
10^{22}	**0.45776 [0.2629]**	**0 [0]**	**0.95932 [0.5510]**
At x=1 and and $r_d = r_{Te}$,			
$\mathbf{1.4590} = E_{gn2}$	**… [0]**	**… [0]**	**… [0]**
$\mathbf{2.0625} = E_{gnE}$	**… [0.3827]**	**… [1.5616]**	**… [0.4289]**
$\mathbf{2.2162} = E_{gn1}$	**0** [0.3827]	**0** [1.4535]	**0** [0.4081]
$\mathbf{2.9766} = E_{gnO}$	**0.5747** [0.3828]	**1.6250** [1.0824]	**0.7334** [0.3675]
4	0.5749 [0.3828]	1.209 [0.8055]	0.737 [0.4470]
4.5	0.5749 [0.3828]	1.075 [0.7160]	0.684 [0.4056]
5	0.5749 [0.3828]	0.968 [0.6444]	0.975 [0.7428]
5.5	0.5749 [0.3829]	0.880 [0.5858]	1.483 [1.7252]
6	0.5749 [0.3829]	0.806 [0.5370]	1.460 [1.5045]
…			
10^{22}	**0.5750 [0.3829]**	**0 [0]**	**1.29585 [0.8629]**
At x=1 and $r_d = r_{Sb}$,			
$\mathbf{1.4602} = E_{gn2}$	**… [0]**	**… [0]**	**… [0]**
$\mathbf{2.0637} = E_{gnE}$	**… [0.3799]**	**… [1.5559]**	**… [0.4281]**
$\mathbf{2.2174} = E_{gn1}$	**0** [0.3799]	**0** [1.4481]	**0** [0.4072]
$\mathbf{2.9778} = E_{gnO}$	**0.5705** [0.3800]	**1.6190** [1.0786]	**0.7322** [0.3670]
4	0.5706 [0.3800]	1.206 [0.8030]	0.736 [0.4462]
4.5	0.5706 [0.3801]	1.072 [0.7138]	0.682 [0.4048]
5	0.5706 [0.3801]	0.964 [0.6424]	0.974 [0.7416]
5.5	0.5707 [0.3801]	0.877 [0.5840]	1.481 [1.7245]
6	0.5707 [0.3801]	0.804 [0.5354]	1.458 [1.5037]
…			
10^{22}	**0.5707 [0.3801]**	**0 [0]**	**1.29435 [0.8620]**

Table 8p. For T=20K and $N = 10^{20} cm^{-3}$, and for given x and r_d, the numerical results of $\sigma_{O[E]}(E)$, $\varepsilon_{2O[2E]}(E)$ and $\propto_{O[E]}(E)$, are obtained by using Equations (18, 19c, 19d), noting that, as given in Eq. (15), $E_{gpE} \equiv E_{gp2} + E_{Fp}$ and $E_{gpO} \equiv E_{gp1} + E_{Fp}$.

E in eV	$\sigma_{O[E]}\ (\frac{10^5}{\Omega\times cm})$	$\varepsilon_{2O[2E]}$	$\propto_{O[E]}\ (\frac{10^5}{cm})$
At x=0 and $r_a = r_{Mg}$,			
1.7452 $= E_{gp2}$	**... [0]**	**... [0]**	**... [0]**
1.9359 $= E_{gpE}$	**... [0.0633]**	**... [0.2627]**	**... [0.0791]**
2.2984 $= E_{gp1}$	**0** [0.0633]	**0** [0.2213]	**0** [0.0715]
2.8542 $=E_{gpO}$	**0.3638** [0.0633]	**1.0234** [0.1782]	**0.4532** [0.0592]
4.5	0.3638 [0.0633]	0.649 [0.1130]	0.417 [0.0663]
5	0.3638 [0.0633]	0.584 [0.1017]	0.580 [0.1084]
5.5	0.3638 [0.0633]	0.531 [0.0925]	0.851 [0.2001]
6	0.3638 [0.0633]	0.487 [0.0848]	0.842 [0.1872]
…			
10^{22}	**0.3638 [0.0633]**	**0 [0]**	**0.7704 [0.1341]**
At x=0 and $r_a = r_{In}$,			
1.7494 $= E_{gp2}$	**... [0]**	**... [0]**	**... [0]**
1.9395 $= E_{gpE}$	**... [0.0599]**	**... [0.2567]**	**... [0.0756]**
2.3008 $= E_{gp1}$	**0** [0.0599]	**0** [0.2164]	**0** [0.0684]
2.8548 $=E_{gpO}$	**0.3411** [0.0599]	**0.9930** [0.1744]	**0.4297** [0.0565]
4.5	0.3411 [0.0599]	0.630 [0.1107]	0.395 [0.0633]
5	0.3411 [0.0599]	0.567 [0.0996]	0.550 [0.1039]
5.5	0.3411 [0.0599]	0.515 [0.0905]	0.812 [0.1936]
6	0.3411 [0.0599]	0.472 [0.0830]	0.804 [0.1810]
…			
10^{22}	**0.3411 [0.0599]**	**0 [0]**	**0.7349 [0.1291]**
At x=0.5 and $r_a = r_{Mg}$,			
1.7574 $= E_{gp2}$	**... [0]**	**... [0]**	**... [0]**
1.9236 $= E_{gpE}$	**... [0.0483]**	**... [0.2118]**	**... [0.0648]**
2.3625 $= E_{gp1}$	**0** [0.0483]	**0** [0.1725]	**0** [0.0572]
2.9701 $=E_{gpO}$	**0.3923** [0.0483]	**1.1150** [0.1372]	**0.5393** [0.0520]
4.5	0.3923 [0.0483]	0.736 [0.0906]	0.486 [0.0539]
5	0.3923 [0.0483]	0.662 [0.0815]	0.667 [0.0884]
5.5	0.3923 [0.0483]	0.602 [0.0741]	0.968 [0.1649]
6	0.3923 [0.0483]	0.552 [0.0679]	0.962 [0.1542]
…			
10^{22}	**0.3923 [0.0483]**	**0 [0]**	**0.8955 [0.1102]**
At x=0.5 and $r_a = r_{In}$,			
1.7621 $= E_{gp2}$	**... [0]**	**... [0]**	**... [0]**
1.9274 $= E_{gpE}$	**... [0.0456]**	**... [0.2068]**	**... [0.0619]**

$\mathbf{2.3638} = E_{gp1}$	**0** [0.0456]	**0** [0.1687]	**0** [0.0547]
2.9680 $=E_{gpO}$	**0.3666** [0.0456]	**1.0792** [0.1343]	**0.5095** [0.0496]
4.5	0.3666 [0.0456]	0.712 [0.0886]	0.458 [0.0514]
5	0.3666 [0.0456]	0.641 [0.0797]	0.631 [0.0846]
5.5	0.3666 [0.0456]	0.582 [0.0725]	0.921 [0.1593]
6	0.3666 [0.0456]	0.534 [0.0664]	0.915 [0.1489]
…			
$\mathbf{10^{22}}$	**0.3666 [0.0456]**	**0 [0]**	**0.8514 [0.1060]**
At x=1 and $r_a = r_{Mg}$,			
$\mathbf{1.7670} = E_{gp2}$	**… [0]**	**… [0]**	**… [0]**
$\mathbf{1.9119} = E_{gpE}$	**… [0.0371]**	**… [0.1728]**	**… [0.0537]**
$\mathbf{2.4673} = E_{gp1}$	**0** [0.0371]	**0** [0.1339]	**0** [0.0455]
3.1699 $=E_{gpO}$	**0.4672** [0.0371]	**1.3112** [0.1043]	**0.7046** [0.0445]
4.5	0.4672 [0.0371]	0.924 [0.0734]	0.633 [0.0443]
5	0.4672 [0.0371]	0.831 [0.0661]	0.850 [0.0731]
5.5	0.4672 [0.0371]	0.756 [0.0601]	1.206 [0.1381]
6	0.4672 [0.0371]	0.693 [0.0551]	1.204 [0.1290]
…			
$\mathbf{10^{22}}$	**0.4672 [0.0371]**	**0 [0]**	**1.1540 [0.0918]**
At x=1 and $r_a = r_{In}$,			
$\mathbf{1.7723} = E_{gp2}$	**… [0]**	**… [0]**	**… [0]**
$\mathbf{1.9160} = E_{gpE}$	**… [0.0350]**	**… [0.1683]**	**… [0.0512]**
$\mathbf{2.4661} = E_{gp1}$	**0** [0.0350]	**0** [0.1308]	**0** [0.0434]
3.1621 $=E_{gpO}$	**0.4340** [0.0350]	**1.2640** [0.1020]	**0.6612** [0.0422]
4.5	0.4340 [0.0350]	0.888 [0.0717]	0.593 [0.0421]
5	0.4340 [0.0350]	0.799 [0.0645]	0.799 [0.0698]
5.5	0.4340 [0.0350]	0.727 [0.0586]	1.142 [0.1330]
6	0.4340 [0.0350]	0.666 [0.0537]	1.140 [0.1242]
…			
$\mathbf{10^{22}}$	**0.4340 [0.0350]**	**0 [0]**	**1.0907 [0.0880]**

Table 9n: For given x, r_d, and T=(4.2 K and 77 K), the numerical results of $\sigma_{O[E]}$, $\mu_{O[E]}$ and $D_{O[E]}$, expressed respectively in $\left(\frac{10^4}{ohm\times cm}, \frac{10^3\times cm^2}{V\times s}, \frac{10^3\times cm^2}{s}\right)$, and as functions of N, are obtained by using Equations (20a, 22a and 24), suggesting that, for a given N, they decrease [decrease], with increasing r_d. In particular, for given (x, r_d and N), those of $\mu_{O[E]}(T)$ increase [increase] with increasing T, due to the decreasing reduced Fermi energy $\xi_{nO[E]}$.

Donor	**Te**	Sb
r_d (nm) ↗	**0.132**	0.136
For x=0 and at T=4.2 K		
N (10^{19} cm^{-3})		
3	1.378 [0.692], 2.918 [1.466], 0.493 [0.163]	1.368 [0.688], 2.897 [1.456], 0.489 [0.161]

7	2.989 [1.465], 2.685 [1.316], 0.803 [0.259]	2.968 [1.455], 2.666 [1.307], 0.798 [0.257]
10	4.150 [2.016], 2.604 [1.265], 0.990 [0.316]	4.120 [2.002], 2.585 [1.256], 0.982 [0.313]
For x=0 and at T=77 K		
N (10^{19} cm^{-3})		
3	76.82 [3.563], 162.6 [7.542], 27.45 [0.835]	76.29 [3.540], 161.5 [7.495], 27.27 [0.830]
7	55.10 [3.397], 49.49 [3.051], 14.81 [0.599]	54.71 [3.373], 49.14 [3.030], 14.70 [0.595]
10	48.97 [3.662], 30.72 [2.298], 11.68 [0.573]	48.62 [3.637], 30.50 [2.282], 11.59 [0.569]
For x=0.5 and at T=4.2 K		
N (10^{19} cm^{-3})		
3	1.531 [0.900], 3.220 [1.892], 0.609 [0.260]	1.520 [0.894], 3.198 [1.880], 0.605 [0.258]
7	3.319 [1.918], 2.973 [1.718], 0.994 [0.417]	3.295 [1.904], 2.952 [1.706], 0.987 [0.414]
10	4.609 [2.647], 2.886 [1.657], 1.225 [0.511]	4.575 [2.628], 2.865 [1.645], 1.216 [0.507]
For x=0.5 and at T=77 K		
N (10^{19} cm^{-3})		
3	84.58 [4.592], 177.8 [9.657], 33.63 [1.326]	83.99 [4.561], 176.6 [9.593], 33.40 [1.317]
7	60.96 [4.435], 54.60 [3.972], 18.25 [0.964]	60.52 [4.404], 54.21 [3.945], 18.12 [0.957]
10	54.26 [4.801], 33.97 [3.006], 14.41 [0.926]	53.86 [4.767], 33.72 [2.985], 14.31 [0.920]
For x=1 and at T=4.2 K		
N (10^{19} cm^{-3})		
3	1.901 [1.284], 3.978 [2.687], 0.901 [0.483]	1.888 [1.275], 3.950 [2.669], 0.895 [0.480]
7	4.133 [2.763], 3.694 [2.469], 1.476 [0.783]	4.103 [2.743], 3.667 [2.452], 1.465 [0.777]
10	5.747 [3.827], 3.593 [2.393], 1.821 [0.963]	5.705 [3.799], 3.567 [2.375], 1.808 [0.956]
For x=1 and at T=77 K		
N (10^{19} cm^{-3})		
3	104.3 [6.513], 218.2 [13.63], 49.41 [2.448]	103.5 [6.467], 216.6 [13.53], 49.06 [2.431]
7	75.69 [6.377], 67.65 [5.700], 27.02 [1.806]	75.14 [6.331], 67.16 [5.659], 26.82 [1.794]
10	67.52 [6.935], 42.21 [4.336], 21.40 [1.744]	67.02 [6.885], 41.90 [4.304], 21.24 [1.731]

Table 9p: For given x, r_a, and T=(4.2 K and 300 K), the numerical results of $\sigma_{O[E]}$, $\mu_{O[E]}$ and $D_{O[E]}$, expressed respectively in $\left(\frac{10^4}{ohm \times cm}, \frac{10^4 \times cm^2}{V \times s}, \frac{10^3 \times cm^2}{s}\right)$, and as functions of N, are obtained by using Equations (20a, 22a and 24), suggesting that, for a given N, they decrease [decrease] with increasing r_a. In particular, for given (x, r_a and N), those of $\mu_{O[E]}(T)$ increase [increase] with increasing T, due to the decreasing reduced Fermi energy $\xi_{pO[E]}$.

Acceptor	Mg	In
r_a (nm) ↗	0.140	0.144
For x=0 and at T=4.2 K		
N (10^{19} cm^{-3})		
3	1.113 [0.218], 0.270 [0.053], 0.417 [0.028]	1.034 [0.205], 0.255 [0.051], 0.389 [0.026]
5	1.861 [0.345], 0.254 [0.047], 0.575 [0.037]	1.739 [0.326], 0.240 [0.045], 0.539 [0.035]
10	3.638 [0.633], 0.237 [0.041], 0.879 [0.052]	3.411 [0.599], 0.223 [0.039], 0.825 [0.050]
For x=0 and at T=300 K		
N (10^{19} cm^{-3})		
3	1.194 [0.286], 0.290 [0.069], 0.439 [0.052]	1.111 [0.268], 0.274 [0.066], 0.411 [0.049]
5	1.924 [0.439], 0.262 [0.060], 0.590 [0.047]	1.798 [0.416], 0.248 [0.057], 0.553 [0.045]
10	3.683 [0.701], 0.240 [0.046], 0.887 [0.057]	3.453 [0.664], 0.226 [0.043], 0.833 [0.054]
For x=0.5 and at T=4.2 K		
N (10^{19} cm^{-3})		
3	1.108 [0.159], 0.301 [0.043], 0.481 [0.019]	1.015 [0.148], 0.286 [0.042], 0.445 [0.018]
5	1.944 [0.260], 0.282 [0.038], 0.684 [0.025]	1.803 [0.244], 0.267 [0.036], 0.638 [0.024]
10	3.923 [0.483], 0.263 [0.032], 1.067 [0.036]	3.666 [0.456], 0.248 [0.031], 1.000 [0.034]
For x=0.5 and at T=300 K		
N (10^{19} cm^{-3})		
3	1.184 [0.207], 0.322 [0.056], 0.505 [0.034]	1.088 [0.193], 0.306 [0.054], 0.469 [0.032]
5	2.001 [0.343], 0.291 [0.050], 0.699 [0.039]	1.857 [0.323], 0.275 [0.048], 0.653 [0.037]
10	3.964 [0.551], 0.266 [0.037], 1.076 [0.040]	3.704 [0.522], 0.251 [0.035], 1.008 [0.038]
For x=1 and at T=4.2 K		
N (10^{19} cm^{-3})		
3	1.135 [0.111], 0.379 [0.037], 0.629 [0.013]	1.008 [0.100], 0.361 [0.036], 0.572 [0.012]
5	2.187 [0.193], 0.353 [0.031], 0.951 [0.017]	2.001 [0.180], 0.334 [0.030], 0.880 [0.016]
10	4.671 [0.371], 0.329 [0.026], 1.540 [0.025]	4.339 [0.350], 0.310 [0.025], 1.437 [0.024]
For x=1 and at T=300 K		
N (10^{19} cm^{-3})		
3	1.207 [0.155], 0.403 [0.052], 0.659 [0.024]	1.078 [0.144], 0.386 [0.051], 0.601 [0.023]
5	2.239 [0.255], 0.361 [0.041], 0.968 [0.032]	2.051 [0.237], 0.342 [0.039], 0.896 [0.030]
10	4.708 [0.440], 0.331 [0.031], 1.550 [0.029]	4.374 [0.416], 0.312 [0.030], 1.446 [0.027]

Table 10n: The numerical results of the viscosity coefficient $\mathbb{V}_{O[E]}(N^{*}, r_d, x, T)$, expressed in $\left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$, are obtained by using Eq. (22b), suggesting that: (i) for given (x, T and N), they increase with increasing r_d, (ii) for given (x, r_d and N) they decrease with increasing T, being due to the decreasing reduced Fermi energy $\xi_{nO[E]}$, in good agreement with those, obtained in liquids by Ewell and Eyring [17] and complex fluids by Wenhao [18], and finally (iii) for given (x, T and r_d) they increase with increasing N, in good agreement with those, obtained in complex fluids by Wenhao [18].

Donor	P	**Te**	Sb	Sn
r_d (nm) [4] ↗	0.110	**0.132**	0.136	0.140
For x=0 and at T=4.2 K				
N (10^{19} cm^{-3})				
3	72.69 [146.5]	90.54 [180.2]	91.17 [181.4]	93.10 [185.0]
7	104.3 [214.9]	130.9 [267.2]	131.9 [269.0]	134.7 [274.6]
10	120.9 [251.2]	152.2 [313.3]	153.3 [315.5]	156.7 [322.2]
For x=0 and at T=77 K				
N (10^{19} cm^{-3})				
3	1.314 [28.64]	1.625 [35.03]	1.635 [35.25]	1.668 [35.92]
7	5.677 [92.86]	7.103 [115.2]	7.154 [116.0]	7.308 [118.4]
10	10.27 [138.4]	12.90 [172.4]	12.99 [173.6]	13.27 [177.3]
For x=0.5 and at T=4.2 K				
N (10^{19} cm^{-3})				
3	65.85 [113.0]	82.22 [139.9]	82.80 [140.8]	84.57 [200.9]
7	94.15 [164.1]	118.3 [204.9]	119.2 [206.3]	121.8 [298.6]
10	109.0 [191.1]	137.4 [239.3]	138.4 [241.0]	141.5 [350.6]
For x=0.5 and at T=77 K				
N (10^{19} cm^{-3})				
3	1.198 [22.23]	1.489 [27.42]	1.499 [27.60]	1.530 [39.29]
7	5.137 [71.02]	6.445 [88.59]	6.491 [89.21]	6.633 [129.0]
10	9.277 [105.4]	11.67 [131.9]	11.76 [132.8]	12.02 [193.2]
For x=1 and at T=4.2 K				
N (10^{19} cm^{-3})				
3	53.21 [79.18]	66.67 [98.69]	67.15 [99.37]	111.1 [218.7]
7	75.65 [113.6]	95.32 [142.6]	96.02 [143.6]	161.2 [325.8]
10	87.46 [131.8]	110.4 [165.8]	111.2 [167.0]	187.6 [383.0]
For x=1 and at T=77 K				
N (10^{19} cm^{-3})				
3	0.972 [15.65]	1.216 [19.46]	1.224 [19.60]	2.024 [43.02]
7	4.135 [49.27]	5.205 [61.78]	5.243 [62.23]	8.798 [141.0]

10	7.449 [72.77]	9.398 [91.50]	9.467 [92.16]	15.97 [211.2]

Table 10p: The numerical results of the viscosity coefficient $\mathbb{V}_{O[E]}(N^{*}, r_a, x, T)$, expressed in $\left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$, are obtained by using Eq. (22b), suggesting that: (i) for given (x, T and N), they increase with increasing r_d, (ii) for given (x, r_a and N) they decrease with increasing T, being due to the decreasing reduced Fermi energy $\xi_{pO[E]}$, in good agreement with those, obtained in liquids by Ewell and Eyring [17] and complex fluids by Wenhao [18], and finally (iii) for given (x, T and r_a) they increase with increasing N, in good agreement with those, obtained in complex fluids by Wenhao [18].

Acceptor	**Ga**	Mg	In	Cd
r_a (nm) ↗	**0.126**	0.140	0.144	0.148
For x=0 and at T=4.2 K				
N (10^{19} cm^{-3})				
3	86.20 [448.1]	93.53 [477.2]	98.24 [495.2]	103.9 [516.4]
5	110.2 [603.7]	120.4 [648.9]	127.1 [677.8]	135.4 [712.8]
10	150.1 [874.4]	164.9 [947.6]	174.8 [995.3]	187.1 [1054]
For x=0 and at T=300 K				
N (10^{19} cm^{-3})				
3	80.53 [341.0]	87.18 [364.0]	91.41 [378.3]	96.49 [395.3]
5	106.7 [475.7]	116.5 [509.8]	122.9 [531.4]	130.8 [557.3]
10	148.3 [789.9]	162.9 [855.3]	172.6 [897.8]	184.8 [950.1]
For x=0.5 and at T=4.2 K				
N (10^{19} cm^{-3})				
3	74.94 [534.8]	80.59 [561.1]	84.06 [576.0]	88.06 [591.7]
5	97.42 [744.7]	106.0 [793.2]	111.5 [823.3]	118.3 [858.6]
10	134.1 [1109]	147.0 [1195]	155.6 [1251]	166.3 [1318]
For x=0.5 and at T=300 K				
N (10^{19} cm^{-3})				
3	70.43 [412.3]	75.44 [431.7]	78.46 [442.0]	81.82 [451.7]
5	94.73 [564.6]	103.0 [600.1]	108.3 [621.9]	114.7 [647.5]
10	132.7 [972.6]	145.5 [1046]	154.0 [1093]	164.6 [1150]
For x=1 and at T=4.2 K				
N (10^{19} cm^{-3})				
3	56.58 [600.5]	59.77 [613.3]	61.43 [616.7]	62.90 [614.4]
5	75.82 [881.3]	81.92 [925.9]	85.72 [951.5]	90.19 [978.9]
10	106.0 [1364]	116.0 [1459]	122.5 [1519]	130.6 [1591]

For x=1 and at T=300 K

$N\ (10^{19}\ cm^{-3})$				
3	53.59 [440.5]	56.22 [437.4]	57.46 [428.4]	58.29 [408.3]
5	74.16 [665.8]	80.02 [701.8]	83.65 [723.2]	87.88 [747.0]
10	105.2 [1154]	115.1 [1230]	121.5 [1277]	129.6 [1332]

Table 11n: For given x, r_d, T=(3K, 80K, 300K) and N, the numerical results of reduced Fermi energy $\xi_{nO[E]}(N^*, r_d, x, T)$, viscosity coefficient $\mathbb{V}_{O[E]}(N^*, r_d, x, T)$, and activation energy $AE_{O[E]}(N^*, r_d, x, T)$ are obtained by using Equations (11, 22b, 22c), respectively, suggesting that for given (x, r_d and N) $\mathbb{V}_{O[E]}$ and $AE_{O[E]}$ decrease with increasing T, due to the decreasing reduced Fermi energy $\xi_{nO[E]}$ with increasing T, in good agreement with those, obtained in liquids by Ewell and Eyring [17] and complex fluids by Wenhao [18].

Donor		Te	Sb	Sn
For **x=0** and **N=9.4754×10^{18} cm^{-3}**,				
$\xi_{nO[E](T=3K)}$	↘	443.181 [291.090]	442.967 [290.950]	442.306 [290.515]
$\xi_{nO[E](T=80K)}$	↘	16.694 [11.031]	16.686 [11.026]	16.661 [11.009]
$\xi_{nO[E](T=300K)}$	↘	4.6738 [2.783]	4.6716 [2.781]	4.6648 [2.775]
$\mathbb{V}_{O[E]}{}_{O[E]\,(3K)} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$	↗	52.786 [100.66]	53.121 [101.25]	54.137 [103.02]
$\mathbb{V}_{O[E]}{}_{O[E]\,(80K)} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$	↗	0.1641 [4.0081]	0.1650 [4.0275]	0.1676 [4.0860]
$\mathbb{V}_{O[E]}{}_{O[E]\,(300K)} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$	↗	0.000500 [0.0224]	0.000503 [0.0226]	0.000511 [0.0229]
$-AE_{O[E]}{}_{O[E]\,(3K)}$ (meV$\times 10^{-6}$)	↗	3.443 [7.708]	3.446 [7.716]	3.457 [7.739]
$-AE_{O[E]}{}_{O[E]\,(80K)}$ (meV)	↗	39.8001 [22.2224]	39.8075 [22.2289]	39.8282 [22.2491]
$-AE_{O[E]}{}_{O[E]\,(300K)}$ (meV)	↗	299.034 [217.369]	299.059 [217.382]	299.138 [217.421]
For **x=0.5** and **N=7.95325×10^{18} cm^{-3}**, one has:				
$\xi_{nO[E](T=3K)}$	↘	444.066 [322.568]	443.910 [322.455]	443.428 [322.104]
$\xi_{nO[E](T=80K)}$	↘	16.727 [12.200]	16.721 [12.195]	16.702 [12.182]
$\xi_{nO[E](T=300K)}$	↘	4.6829 [3.231]	4.6813 [3.229]	4.6763 [3.224]
$\mathbb{V}_{O[E]}{}_{O[E]\,(3K)} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$	↗	44.588 [73.224]	44.875 [73.666]	45.748 [75.011]
$\mathbb{V}_{O[E]}{}_{O[E]\,(80K)} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$	↗	0.1120 [2.3827]	0.1127 [2.3954]	0.1146 [2.4340]
$\mathbb{V}_{O[E]}{}_{O[E]\,(300K)} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$	↗	0.000341 [0.0120]	0.000343 [0.0121]	0.000349 [0.0123]
$-AE_{O[E]}{}_{O[E]\,(3K)}$ (meV$\times 10^{-6}$)	↗	3.477 [6.354]	3.479 [6.358]	3.487 [6.372]
$-AE_{O[E]}{}_{O[E]\,(80K)}$ (meV)	↗	41.2689 [23.614]	41.2738 [23.618]	41.2888 [23.633]
$-AE_{O[E]}{}_{O[E]\,(300K)}$ (meV)	↗	304.553 [225.29]	304.572 [225.30]	304.629 [225.33]
For **x=1** and **N=6.01616×10^{18} cm^{-3}**, one has:				
$\xi_{nO[E](T=3K)}$	↘	443.635 [352.092]	443.527 [352.007]	443.195 [351.743]

$\xi_{nO[E](T=80K)}$ $\searrow$	16.711 [13.298]	16.707 [13.295]	16.694 [13.285]
$\xi_{nO[E](T=300K)}$ $\searrow$	4.6785 [3.630]	4.6774 [3.629]	4.6739 [3.626]
$\mathbb{V}_{O[E]_{O[E]\,(3K)}} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$ $\nearrow$	32.280 [46.442]	32.492 [46.733]	33.136 [47.621]
$\mathbb{V}_{O[E]_{O[E]\,(80K)}} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$ $\nearrow$	0.0569 [1.0727]	0.0572 [1.0789]	0.0583 [1.0977]
$\mathbb{V}_{O[E]_{O[E]\,(300K)}} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$ $\nearrow$	0.000173 [0.00504]	0.000174 [0.00507]	0.000177 [0.000516]
$-AE_{O[E]_{O[E]\,(3K)}}$ $(meV \times 10^{-6})$ $\nearrow$	3.586 [5.4637]	3.588 [5.4664]	3.593 [5.4746]
$-AE_{O[E]_{O[E]\,(80K)}}$ (meV) $\nearrow$	43.7144 [25.976]	43.7177 [25.980]	43.7281 [25.990]
$-AE_{O[E]_{O[E]\,(300K)}}$ (meV) $\nearrow$	313.744 [235.975]	313.757 [235.984]	313.796 [236.013]

Table 11p: For given x, r_a, T=(3K, 80K, 300K) and N, the numerical results of reduced Fermi energy $\xi_{pO[E]}(N^*, r_a, x, T)$, viscosity coefficient $\mathbb{V}_{O[E]}(N^*, r_a, x, T)$, and activation energy $AE_{O[E]}(N^*, r_a, x, T)$ are obtained by using Equations (11, 22b, 22c), respectively, suggesting that for given (x, r_a and N) $\mathbb{V}_{O[E]}$ and $AE_{O[E]}$ decrease with increasing T, due to the decreasing reduced Fermi energy $\xi_{pO[E]}$ with increasing T, in good agreement with those, obtained in liquids by Ewell and Eyring [17] and complex fluids by Wenhao [18].

Acceptor	Ga	Mg	In
For x=0 and **N=6.85437×10^{19} cm^{-3}** one has:			
$\xi_{pO[E](T=3K)}$ $\searrow$	1659.6 [569.556]	1648.4 [565.722]	1640.5 [562.989]
$\xi_{pO[E](T=80K)}$ $\searrow$	62.255 [21.416]	61.836 [21.273]	61.538 [21.171]
$\xi_{nO[E](T=300K)}$ $\searrow$	16.671 [5.918]	16.560 [5.881]	16.481 [5.854]
$\mathbb{V}_{O[E]_{O[E]\,(3K)}} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$ $\nearrow$	127.18 [717.43]	139.37 [774.40]	147.42 [811.21]
$\mathbb{V}_{O[E]_{O[E]\,(80K)}} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$ $\nearrow$	126.99 [708.48]	139.16 [764.60]	147.20 [800.85]
$\mathbb{V}_{O[E]_{O[E]\,(300K)}} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$ $\nearrow$	124.5689 [607.49]	136.4682 [654.40]	144.3236 [684.51]
$-AE_{O[E]_{O[E]\,(3K)}}$ $(meV \times 10^{-6})$ $\nearrow$	0.540 [4.588]	0.548 [4.650]	0.553 [4.696]
$-AE_{O[E]_{O[E]\,(80K)}}$ (meV) $\nearrow$	0.01024 [0.087]	0.01038 [0.088]	0.01048 [0.089]
$-AE_{O[E]_{O[E]\,(300K)}}$ (meV) $\nearrow$	0.53731 [4.301]	0.5446 [4.353]	0.5498 [4.391]
For x=0.5 and **N=6.11283×10^{19} cm^{-3}** one has:			
$\xi_{pO[E](T=3K)}$ $\searrow$	1659.6 [454.084]	1637.8 [448.121]	1622.2 [443.859]
$\xi_{pO[E](T=80K)}$ $\searrow$	62.255 [17.101]	61.438 [16.878]	60.854 [16.719]
$\xi_{nO[E](T=300K)}$ $\searrow$	16.671 [4.785]	16.454 [4.724]	16.299 [4.681]
$\mathbb{V}_{O[E]_{O[E]\,(3K)}} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$ $\nearrow$	107.22 [839.88]	117.01 [898.30]	123.40 [935.09]
$\mathbb{V}_{O[E]_{O[E]\,(80K)}} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$ $\nearrow$	107.06 [823.48]	116.8 [880.30]	123.21 [915.99]
$\mathbb{V}_{O[E]_{O[E]\,(300K)}} \left(\frac{eV}{cm} \times \frac{s}{cm^2}\right)$ $\nearrow$	105.0126 [661.27]	114.5428 [704.36]	120.7455 [731.07]
$-AE_{O[E]_{O[E]\,(3K)}}$ $(meV \times 10^{-6})$ $\nearrow$	0.540 [7.218]	0.555 [7.412]	0.565 [7.555]
$-AE_{O[E]_{O[E]\,(80K)}}$ (meV) $\nearrow$	0.01024 [0.136]	0.01052 [0.140]	0.01072 [0.142]

$-AE_{O[E]_{O[E]\,(300K)}}$ (meV)	↗	0.53731 [6.182]	0.5516 [6.288]	0.5622 [6.364]
For x=1 and **N=5.1903×10^{19} cm^{-3}**, one has:				
$\xi_{pO[E](T=3K)}$	↘	1659.6 [342.462]	1613.8 [333.010]	1580.8 [326.210]
$\xi_{pO[E](T=80K)}$	↘	62.255 [12.940]	60.538 [12.588]	59.302 [12.335]
$\xi_{nO[E](T=300K)}$	↘	16.671 [3.503]	16.215 [3.375]	15.887 [3.281]
$\mathbb{V}_{O[E]_{O[E]\,(3K)}}\left(\frac{eV}{cm}\times\frac{s}{cm^2}\right)$	↗	77.306 [904.03]	83.608 [951.14]	87.563 [978.46]
$\mathbb{V}_{O[E]_{O[E]\,(80K)}}\left(\frac{eV}{cm}\times\frac{s}{cm^2}\right)$	↗	77.192 [873.35]	83.476 [917.06]	87.420 [941.96]
$\mathbb{V}_{O[E]_{O[E]\,(300K)}}\left(\frac{eV}{cm}\times\frac{s}{cm^2}\right)$	↗	75.71628 [681.32]	81.79069 [718.40]	85.5819 [740.63]
$-AE_{O[E]_{O[E]\,(3K)}}$ (meV$\times 10^{-6}$)	↗	0.540 [12.69]	0.571 [13.42]	0.595 [13.99]
$-AE_{O[E]_{O[E]\,(80K)}}$ (meV)	↗	0.01024 [0.238]	0.01083 [0.252]	0.01129 [0.262]
$-AE_{O[E]_{O[E]\,(300K)}}$ (meV)	↗	0.53731 [7.313]	0.5680 [7.256]	0.5918 [7.200]

Table 12n: For given x, r_d, T=(3K and 80K) and N, the numerical results of various thermoelectric coefficients: $\sigma_{Th.O[E]}$, $S_{O[E]}$, $VC1_{O[E]}$, $VC2_{O[E]}$, $Ts_{O[E]}$, $Pt_{O[E]}$ and $ZT_{O[E]}$, are obtained by using Equations (21, 25, 27, 28, 29, 30 and 26), respectively. Further, their variations with increasing r_d are represented by the arrows: ↗ (increase), and ↘ (decrease).

Donor		**Te**	Sb	Sn
For **x=0** and **N=9.4754×10^{18} cm^{-3}**,				
$\xi_{nO[E](T=3K)}$	↘	443.181 [291.090]	442.967 [290.950]	442.306 [290.515]
$\xi_{nO[E](T=80K)}$	↘	16.694 [11.031]	16.686 [11.026]	16.661 [11.009]
$\sigma_{Th.O[E]\,(3K)}\left(\frac{10^{-4}\times W}{cm\times K}\right)$	↘	3.540 [1.857]	3.515 [1.844]	3.438 [1.807]
$\sigma_{Th.O[E]\,(80K)}\left(\frac{10^{-1}\times W}{cm\times K}\right)$	↘	30.36 [1.243]	30.17 [1.236]	29.61 [1.215]
$-S_{O[E]\,(3K)}\left(\frac{10^{-6}\times V}{K}\right)$	↗	1.279 [1.948]	1.280 [1.949]	1.282 [1.952]
$-S_{O[E]\,(80K)}\left(\frac{10^{-5}\times V}{K}\right)$	↗	3.356 [5.004]	3.358 [5.007]	3.363 [5.014]
$-VC1_{O[E]\,(3K)}\left(\frac{10^{-7}\times V}{K}\right)$	↗	8.528 [12.98]	8.533 [12.99]	8.545 [13.01]
$-VC1_{O[E]\,(80K)}\left(\frac{10^{-5}\times V}{K}\right)$	↗	2.165 [3.093]	2.167 [3.095]	2.170 [3.098]
$-VC2_{O[E]\,(3K)}\left(\frac{10^{-6}\times V}{K}\right)$	↗	2.558 [3.895]	2.560 [3.897]	2.564 [3.903]
$-VC2_{O[E]\,(80K)}\left(\frac{10^{-3}\times V}{K}\right)$	↗	1.732 [2.475]	1.733 [2.476]	1.736 [2.479]
$-Ts_{O[E]\,(3K)}\left(\frac{10^{-6}\times V}{K}\right)$	↗	1.279 [1.947]	1.280 [1.948]	1.282 [1.951]
$-Ts_{O[E]\,(80K)}\left(\frac{10^{-5}\times V}{K}\right)$	↗	3.248 [4.640]	3.250 [4.642]	3.254 [4.648]
$-Pt_{O[E]\,(3K)}(10^{-6}\times V)$	↗	3.838 [5.843]	3.840 [5.846]	3.846 [5.855]
$-Pt_{O[E]\,(80K)}(10^{-3}\times V)$	↗	2.685 [4.004]	2.687 [4.005]	2.690 [4.011]
$ZT_{O[E]\,(3K)}\ (10^{-5})$	↗	6.699 [15.53]	6.706 [15.54]	6.726 [15.59]

$ZT_{O[E]\,(80K)}(10^{-2})$	↗	4.612 [10.25]	4.617 [10.26]	4.630 [10.29]

For **x=0.5** and **N=$\mathbf{7.95325 \times 10^{18}\ cm^{-3}}$**, one has:

$\xi_{nO[E](T=3K)}$	↘	444.066 [322.568]	443.910 [322.455]	443.428 [322.104]
$\xi_{nO[E](T=80K)}$	↘	16.727 [12.200]	16.721 [12.195]	16.702 [12.182]
$\sigma_{Th.O[E]]\,(3K)}\ (\frac{10^{-4}\times W}{cm\times K})$	↘	3.385 [2.061]	3.361 [2.048]	3.290 [2.006]
$\sigma_{Th.O[E]]\,(80K)}\ (\frac{10^{-1}\times W}{cm\times K})$	↘	35.92 [1.689]	35.69 [1.679]	35.01 [1.649]
$-S_{O[E]\,(3K)}\ (\frac{10^{-6}\times V}{K})$	↗	1.276 [1.757]	1.277 [1.758]	1.279 [1.760]
$-S_{O[E]\,(80K)}\ (\frac{10^{-5}\times V}{K})$	↗	3.350 [4.547]	3.351 [4.548]	3.355 [4.553]
$-VC1_{O[E]\,(3K)}\ (\frac{10^{-7}\times V}{K})$	↗	8.511 [11.71]	8.514 [11.72]	8.524 [11.73]
$-VC1_{O[E]\,(80K)}\ \left(\frac{10^{-5}\times V}{K}\right)$	↗	2.162 [2.850]	2.163 [2.851]	2.165 [2.854]
$-VC2_{O[E]\,(3K)}\ (\frac{10^{-6}\times V}{K})$	↗	2.553 [3.515]	2.554 [3.516]	2.557 [3.520]
$-VC2_{O[E]\,(80K)}\ \left(\frac{10^{-3}\times V}{K}\right)$	↗	1.729 [2.280]	1.730 [2.281]	1.732 [2.283]
$-Ts_{O[E]\,(3K)}\ (\frac{10^{-6}\times V}{K})$	↗	1.276 [1.757]	1.277 [1.758]	1.278 [1.760]
$-Ts_{O[E]\,(80K)}\ \left(\frac{10^{-5}\times V}{K}\right)$	↗	3.243 [4.275]	3.244 [4.276]	3.247 [4.280]
$-Pt_{O[E]\,(3K)}(10^{-6}\times V)$	↗	3.830 [5.273]	3.831 [5.275]	3.836 [5.280]
$-Pt_{O[E]\,(80K)}(10^{-3}\times V)$	↗	2.680 [3.637]	2.681 [3.638]	2.684 [3.642]
$ZT_{O[E]\,(3K)}\ (10^{-5})$	↗	6.673 [12.64]	6.677 [12.65]	6.692 [12.68]
$ZT_{O[E]\,(80K)}(10^{-2})$	↗	4.594 [8.463]	4.597 [8.468]	4.607 [8.486]

For **x=1** and **N=$\mathbf{6.01616 \times 10^{18}\ cm^{-3}}$**, one has:

$\xi_{nO[E](T=3K)}$	↘	443.635 [352.092]	443.527 [352.007]	443.195 [351.743]
$\xi_{nO[E](T=80K)}$	↘	16.711 [13.298]	16.707 [13.295]	16.694 [13.285]
$\sigma_{Th.O[E]]\,(3K)}\ (\frac{10^{-4}\times W}{cm\times K})$	↘	3.277 [2.278]	3.254 [2.263]	3.186 [2.217]
$\sigma_{Th.O[E]]\,(80K)}\ (\frac{10^{-1}\times W}{cm\times K})$	↘	49.58 [2.630]	49.26 [2.613]	48.31 [2.565]
$-S_{O[E]\,(3K)}\ (\frac{10^{-6}\times V}{K})$	↗	1.277 [1.610]	1.278 [1.611]	1.279 [1.612]
$-S_{O[E]\,(80K)}\ (\frac{10^{-5}\times V}{K})$	↗	3.353 [4.186]	3.354 [4.187]	3.357 [4.190]
$-VC1_{O[E]\,(3K)}\ (\frac{10^{-7}\times V}{K})$	↗	8.520 [10.73]	8.522 [10.74]	8.528 [10.75]
$-VC1_{O[E]\,(80K)}\ \left(\frac{10^{-5}\times V}{K}\right)$	↗	2.163 [2.649]	2.164 [2.650]	2.166 [2.652]
$-VC2_{O[E]\,(3K)}\ (\frac{10^{-6}\times V}{K})$	↗	2.556 [3.220]	2.557 [3.221]	2.558 [3.223]
$-VC2_{O[E]\,(80K)}\ \left(\frac{10^{-3}\times V}{K}\right)$	↗	1.730 [2.119]	1.731 [2.120]	1.733 [2.122]
$-Ts_{O[E]\,(3K)}\ (\frac{10^{-6}\times V}{K})$	↗	1.277 [1.610]	1.278 [1.611]	1.279 [1.612]

$-Ts_{O[E]\,(80K)}\left(\frac{10^{-5}\times V}{K}\right)$	↗	3.245 [3.975]	3.246 [3.976]	3.249 [3.978]
$-Pt_{O[E]\,(3K)}(10^{-6}\times V)$	↗	3.834 [4.831]	3.835 [4.432]	3.838 [4.836]
$-Pt_{O[E]\,(80K)}(10^{-3}\times V)$	↗	2.682 [3.348]	2.683 [3.349]	2.685 [3.352]
$ZT_{O[E]\,(3K)}\ (10^{-5})$	↗	6.686 [10.61]	6.689 [10.62]	6.699 [10.63]
$ZT_{O[E]\,(80K)}(10^{-2})$	↗	4.603 [7.172]	4.605 [7.175]	4.612 [7.186]

Table 12p: For given x, r_a, T=(3K and 80K) and N, the numerical results of various thermoelectric coefficients: $\sigma_{Th.O[E]}$, $S_{O[E]}$, $VC1_{O[E]}$, $VC2_{O[E]}$, $Ts_{O[E]}$, $Pt_{O[E]}$ and $ZT_{O[E]}$, are obtained by using Equations (21, 25, 27, 28, 29, 30 and 26), respectively. Further, their variations with increasing r_a are represented by the arrows: ↗ (increase), and ↘ (decrease).

Acceptor		Ga	Mg	In
For x=0 and **N=6.85437×10^{19} cm^{-3}** one has:				
$\xi_{pO[E](T=3K)}$	↘	1659.6 [569.556]	1648.4 [565.722]	1640.5 [562.989]
$\xi_{pO[E](T=80K)}$	↘	62.255 [21.416]	61.836 [21.273]	61.538 [21.171]
$\sigma_{Th.O[E]]\,(3K)}\left(\frac{10^{-3}\times W}{cm\times K}\right)$	↘	2.061 [0.365]	1.855 [0.334]	1.737 [0.316]
$\sigma_{Th.O[E]]\,(80K)}\left(\frac{10^{-2}\times W}{cm\times K}\right)$	↘	5.503 [0.986]	4.955 [0.901]	4.639 [0.853]
$-S_{O[E]\,(3K)}\left(\frac{10^{-7}\times V}{K}\right)$	↗	3.416 [9.955]	3.439 [10.02]	3.456 [10.07]
$-S_{O[E]\,(80K)}\left(\frac{10^{-6}\times V}{K}\right)$	↗	9.100 [26.28]	9.161 [26.46]	9.206 [26.59]
$-VC1_{O[E]\,(3K)}\left(\frac{10^{-7}\times V}{K}\right)$	↗	2.277 [6.636]	2.293 [6.681]	2.304 [6.714]
$-VC1_{O[E]\,(80K)}\left(\frac{10^{-6}\times V}{K}\right)$	↗	6.052 [17.18]	6.093 [17.29]	6.122 [17.37]
$-VC2_{O[E]\,(3K)}\left(\frac{10^{-7}\times V}{K}\right)$	↗	6.833 [19.91]	6.879 [20.04]	6.912 [20.14]
$-VC2_{O[E]\,(80K)}\left(\frac{10^{-4}\times V}{K}\right)$	↗	4.842 [13.74]	4.874 [13.83]	4.898 [13.89]
$-Ts_{O[E]\,(3K)}\left(\frac{10^{-7}\times V}{K}\right)$	↗	3.416 [9.954]	3.439 [10.02]	3.456 [10.07]
$-Ts_{O[E]\,(80K)}\left(\frac{10^{-6}\times V}{K}\right)$	↗	9.078 [25.77]	9.139 [25.93]	9.184 [26.05]
$-Pt_{O[E]\,(3K)}(10^{-6}\times V)$	↗	1.025 [2.98]	1.032 [3.00]	1.037 [3.02]
$-Pt_{O[E]\,(80K)}(10^{-4}\times V)$	↗	7.280 [21.03]	7.329 [21.17]	7.364 [21.27]
$ZT_{O[E]\,(3K)}\ (10^{-6})$	↗	4.778 [40.56]	4.843 [41.12]	4.890 [41.52]
$ZT_{O[E]\,(80K)}(10^{-3})$	↗	3.389 [28.28]	3.435 [28.66]	3.469 [28.93]
For x=0.5 and **N=6.11283×10^{19} cm^{-3}** one has:				
$\xi_{pO[E](T=3K)}$	↘	1659.6 [454.084]	1637.8 [448.121]	1622.2 [443.859]
$\xi_{pO[E](T=80K)}$	↘	62.255 [17.101]	61.438 [16.878]	60.854 [16.719]
$\sigma_{Th.O[E]]\,(3K)}\left(\frac{10^{-3}\times W}{cm\times K}\right)$	↘	1.966 [0.251]	1.755 [0.228]	1.632 [0.215]
$\sigma_{Th.O[E]]\,(80K)}\left(\frac{10^{-2}\times W}{cm\times K}\right)$	↘	5.252 [0.683]	4.687 [0.622]	4.360 [0.586]

$-S_{O[E]\,(3K)}\ (\frac{10^{-7}\times V}{K})$	↗	3.416 [12.48]	3.462 [12.65]	3.495 [12.77]
$-S_{O[E]\,(80K)}\ (\frac{10^{-6}\times V}{K})$	↗	9.100 [32.78]	9.220 [33.21]	9.309 [33.52]
$-VC1_{O[E]\,(3K)}\ (\frac{10^{-7}\times V}{K})$	↗	2.277 [8.324]	2.308 [8.434]	2.330 [8.515]
$-VC1_{O[E]\,(80K)}\ \left(\frac{10^{-6}\times V}{K}\right)$	↗	6.052 [21.19]	6.132 [21.44]	6.191 [21.63]
$-VC2_{O[E]\,(3K)}\ (\frac{10^{-7}\times V}{K})$	↗	6.833 [24.97]	6.924 [25.30]	6.990 [25.55]
$-VC2_{O[E]\,(80K)}\ \left(\frac{10^{-4}\times V}{K}\right)$	↗	4.842 [16.95]	4.906 [17.15]	4.953 [17.30]
$-Ts_{O[E]\,(3K)}\ (\frac{10^{-7}\times V}{K})$	↗	3.416 [12.48]	3.462 [12.65]	3.495 [12.77]
$-Ts_{O[E]\,(80K)}\ \left(\frac{10^{-6}\times V}{K}\right)$	↗	9.078 [31.78]	9.198 [32.16]	9.286 [32.44]
$-Pt_{O[E]\,(3K)}(10^{-6}\times V)$	↗	1.025 [3.746]	1.038 [3.795]	1.048 [3.832]
$-Pt_{O[E]\,(80K)}(10^{-4}\times V)$	↗	7.280 [26.23]	7.376 [26.56]	7.447 [26.81]
$ZT_{O[E]\,(3K)}\ (10^{-6})$	↗	4.778 [63.82]	4.906 [65.53]	5.000 [66.76]
$ZT_{O[E]\,(80K)}(10^{-3})$	↗	3.389 [44.00]	3.480 [45.14]	3.547 [45.99]

For x=1 and **N=$\mathbf{5.1903\times10^{19}\ cm^{-3}}$**, one has:

$\xi_{pO[E](T=3K)}$	↘	1659.6 [342.462]	1613.8 [333.010]	1580.8 [326.210]
$\xi_{pO[E](T=80K)}$	↘	62.255 [12.940]	60.538 [12.588]	59.302 [12.335]
$\sigma_{Th\,.O[E]]\,(3K)}\ (\frac{10^{-3}\times W}{cm\times K})$	↘	1.915 [0.164]	1.674 [0.147]	1.534 [0.137]
$\sigma_{Th\,.O[E]]\,(80K)}\ (\frac{10^{-2}\times W}{cm\times K})$	↘	5.114 [0.452]	4.472 [0.407]	4.098 [0.380]
$-S_{O[E]\,(3K)}\ (\frac{10^{-7}\times V}{K})$	↗	3.416 [16.55]	3.513 [17.02]	3.587 [17.38]
$-S_{O[E]\,(80K)}\ (\frac{10^{-6}\times V}{K})$	↗	9.100 [42.97]	9.357 [44.12]	9.552 [44.99]
$-VC1_{O[E]\,(3K)}\ (\frac{10^{-7}\times V}{K})$	↗	2.277 [11.03]	2.342 [11.35]	2.391 [11.59]
$-VC1_{O[E]\,(80K)}\ \left(\frac{10^{-6}\times V}{K}\right)$	↗	6.052 [27.12]	6.223 [27.76]	6.352 [28.24]
$-VC2_{O[E]\,(3K)}\ (\frac{10^{-7}\times V}{K})$	↗	6.833 [33.11]	7.027 [34.04]	7.173 [34.76]
$-VC2_{O[E]\,(80K)}\ \left(\frac{10^{-4}\times V}{K}\right)$	↗	4.842 [21.70]	4.978 [22.21]	5.081 [22.59]
$-Ts_{O[E]\,(3K)}\ (\frac{10^{-7}\times V}{K})$	↗	3.416 [16.55]	3.513 [17.02]	3.587 [17.38]
$-Ts_{O[E]\,(80K)}\ \left(\frac{10^{-6}\times V}{K}\right)$	↗	9.078 [40.68]	9.334 [41.65]	9.527 [42.36]
$-Pt_{O[E]\,(3K)}(10^{-6}\times V)$	↗	1.025 [4.967]	1.054 [5.108]	1.076 [5.214]
$-Pt_{O[E]\,(80K)}(10^{-4}\times V)$	↗	7.280 [34.38]	7.486 [35.30]	7.642 [35.99]
$ZT_{O[E]\,(3K)}\ (10^{-6})$	↗	4.778 [112.2]	5.053 [118.6]	5.266 [123.6]
$ZT_{O[E]\,(80K)}(10^{-3})$	↗	3.389 [75.59]	3.584 [79.70]	3.735 [82.86]

Table 13: Here, in the O-EP [E-OP] and for given physical conditions: x, $r_{d(a)}$, N (or T), the same values of $\xi_{n(p)}$ decrease, according to the increasing T (or to the decreasing N), and other thermoelectric coefficients are in variations, as indicated by the arrows as: (increase: ↗, decrease: ↘). One notes here that (i) for $\xi_{n(p)O[E]} \simeq 1.8138$, while the numerical results of $S_{O[E]}$ present a same minimum $S_{O[E]\,min.} \left(\simeq -1.563 \times 10^{-4}\frac{V}{K}\right)$, those of $ZT_{O[E]}$ show a same maximum $ZT_{O[E]\,max.} = 1$, (ii) for $\xi_p = 1$, those of $S_{O[E]}$, $ZT_{O[E]}$, $ZT_{O[E]Mott}$, $VC1_{E[O]}$, and $Ts_{O[E]}$ present the same results: $-1.322 \times 10^{-4}\frac{V}{K}$, 0.715, 3.290, $1.105 \times 10^{-4}\frac{V}{K}$, and $1.657 \times 10^{-4}\frac{V}{K}$, respectively, and (iii) for $\xi_{n(p)O[E]} \simeq 1.8138$, $(ZT)_{O[E]\,Mott} = 1$.

$\xi_{n(p)O[E]}$ ↘	1.880 [1.880]		**1.8138 [1.8138]**		1.750 [1.750]		**1 [1]**		0.998 [0.998]
$S_{O[E]}\left(10^{-4}\frac{V}{K}\right)$	−1.562 [−1.562]	↘	**−1.563** [**−1.563**]	↗	−1.562 [−1.562]	↗	**−1.322** [**−1.322**]	↗	−1.320 [−1.320]
$ZT_{O[E]}$	0.999 [0.999]	↗	**1** [1]	↘	0.999 [0.999]	↘	**0.715 [0.715]**	↘	0.713 [0.713]
$(ZT)_{O[E]\,Mott}$ ↗	0.931 [0.931]		**1** [1]		1.074 [1.074]		**3.290 [3.290]**		3.306 [3.306]
$VC1_{E[O]}\left(10^{-4}\frac{V}{K}\right)$	−0.061[−0.061]	↗	0 [0]	↗	0.063 [0.063]	↗	**1.105 [1.105]**	↗	1.109 [1.109]
$Ts_{O[E]}\left(10^{-4}\frac{V}{K}\right)$	−0.092 [−0.092]	↗	0 [0]	↗	0.094 [0.094]	↗	**1.657 [1.657]**	↗	1.663 [1.663]